

# Experimental Studies of Color Propagation

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# Outline

- Introduction
- Direct measurement of quark energy loss
- Extracting characteristic times: semi-inclusive DIS
  - HERMES data - comparison of our model results to Lund string model
  - JLab data - strong evidence for time dilation, comparison to string model HERMES
  - Connections to QCD factorization, to much higher energies, and hadronization in vacuum
- Extrapolation to 12 GeV and EIC kinematics

# Aims

## Quark-Hadron Transition

Discover new fundamental features of hadronization

- Characteristic time distributions
- Mechanisms of color neutralization

## Quark-Nucleus Interaction

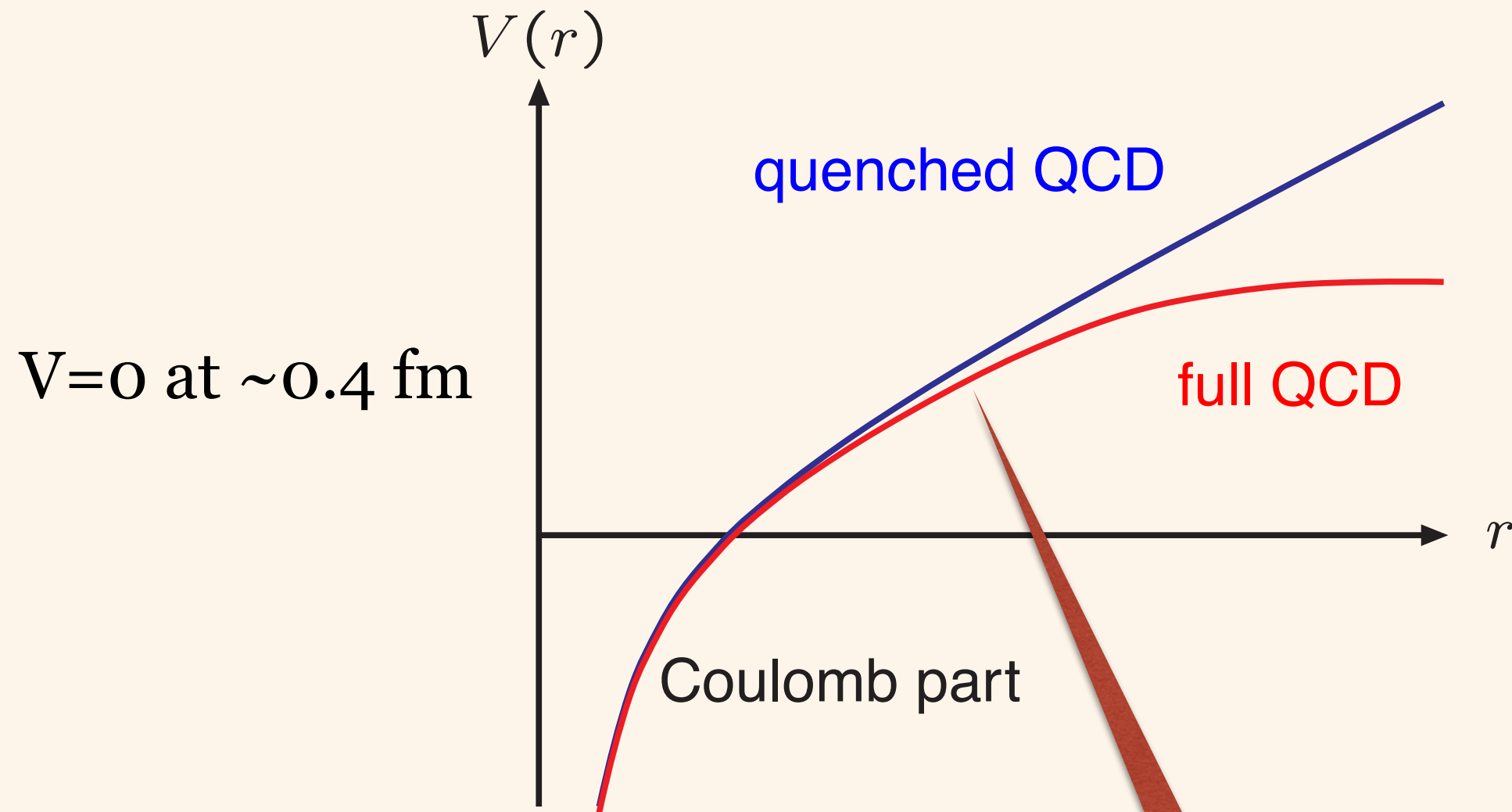
Understand how color interacts within nuclei

- Partonic interactions with medium
  - energy loss in-medium:  $\hat{e}$
  - transverse momentum broadening:  $\hat{q}$

*Method: struck quark from DIS probes nuclei of different sizes*



# Connection to Confinement



Dynamical enforcement of  
confinement begins here

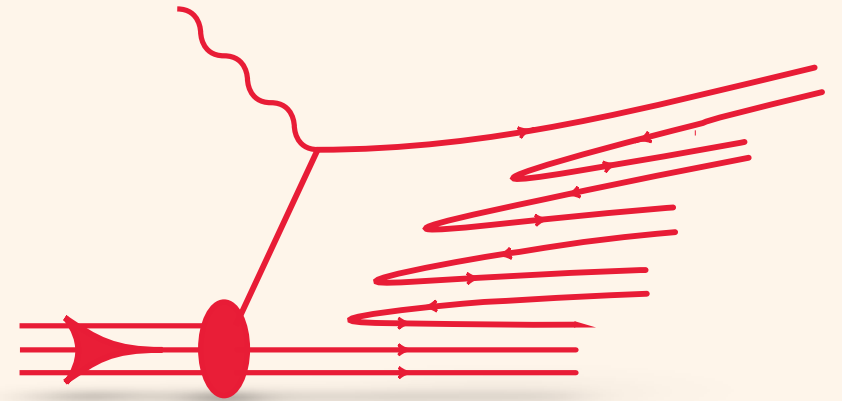
Beyond  $\sim 1$  fm the potential is irrelevant but confinement is still enforced



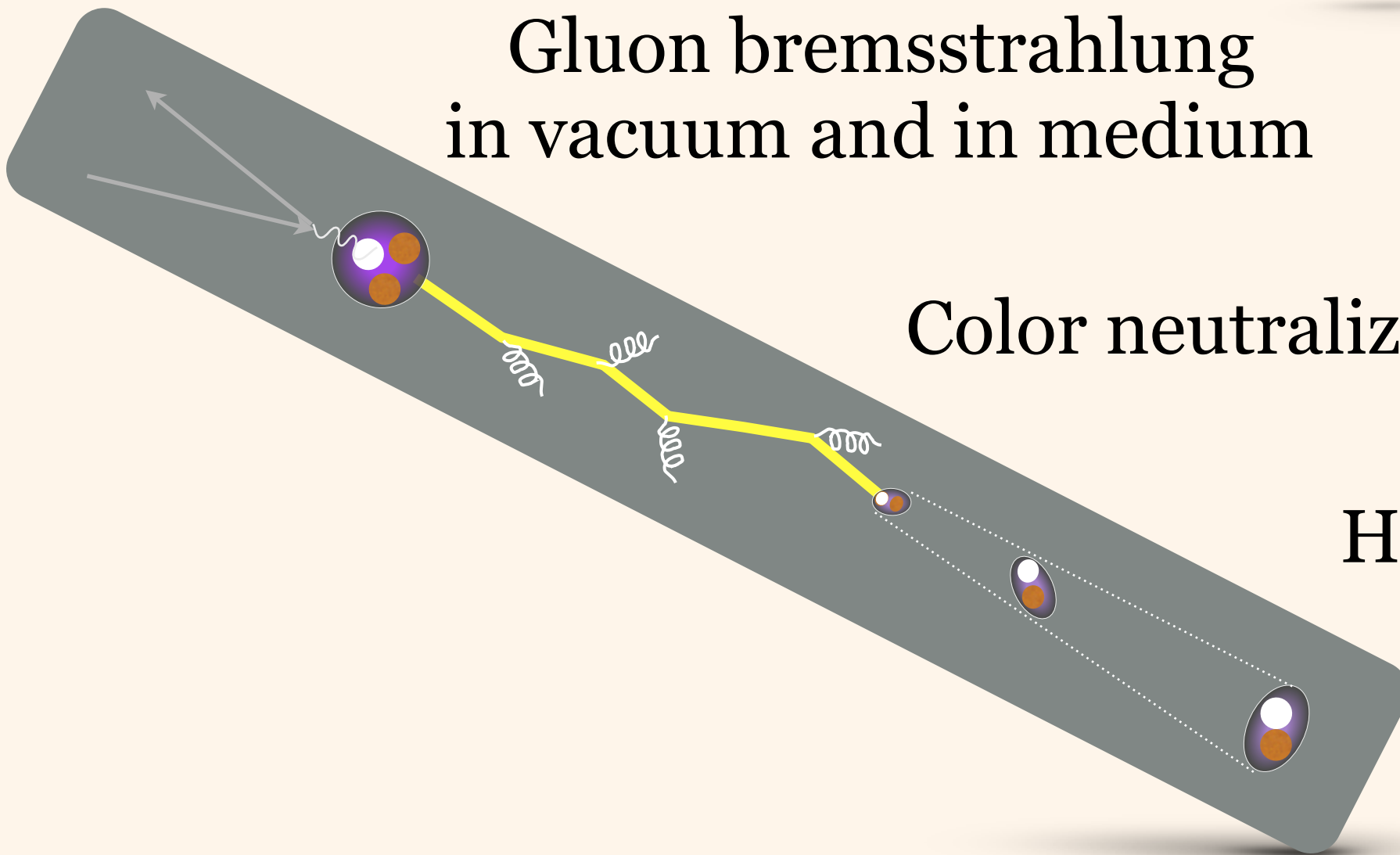
# FUNDAMENTAL QCD PROCESSES

(DIS, pQCD picture)

Partonic elastic scattering  
in medium



Gluon bremsstrahlung  
in vacuum and in medium

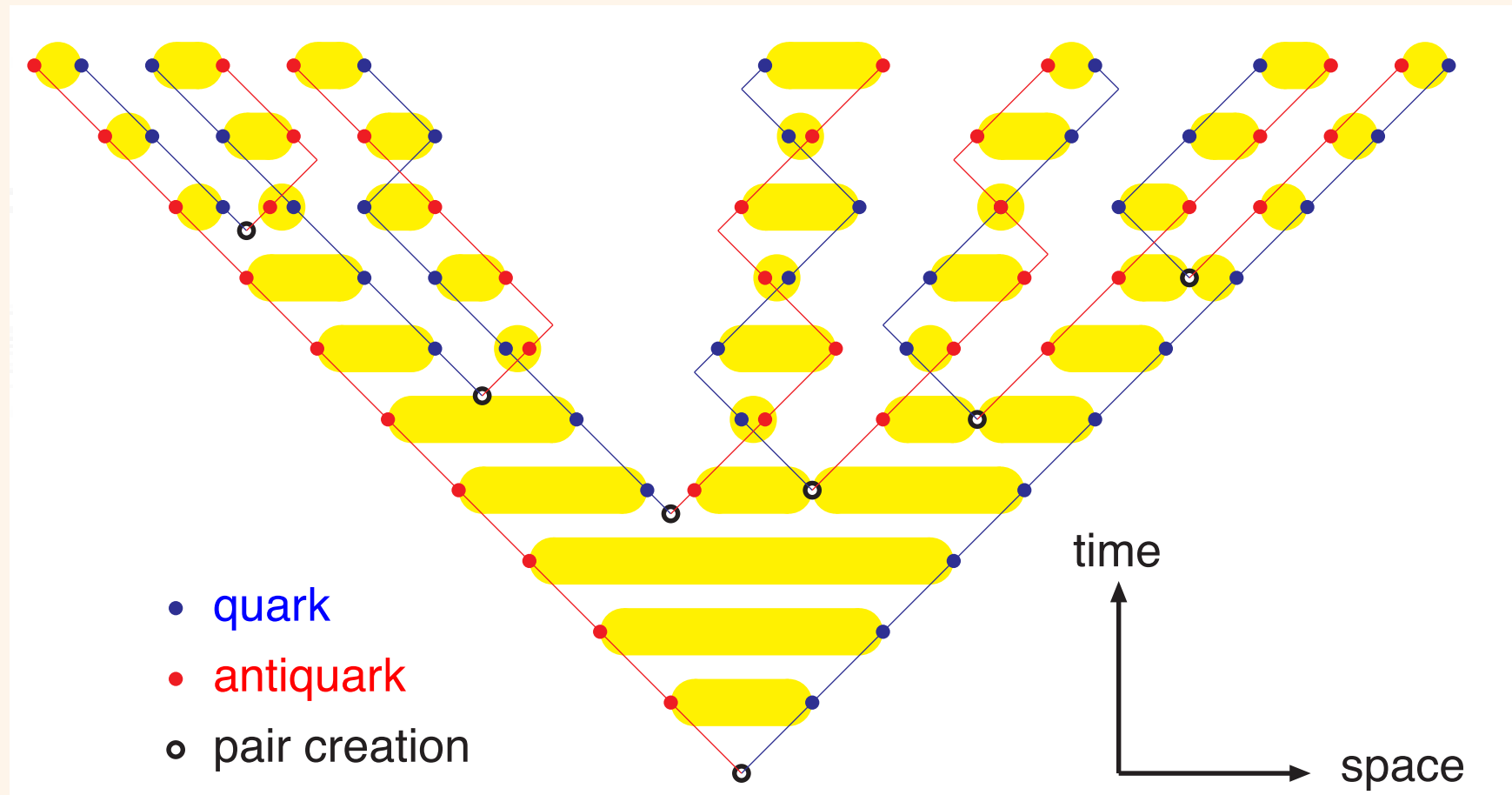


Color neutralization

Hadron formation

We implant this process in large nuclei and compare to deuterium

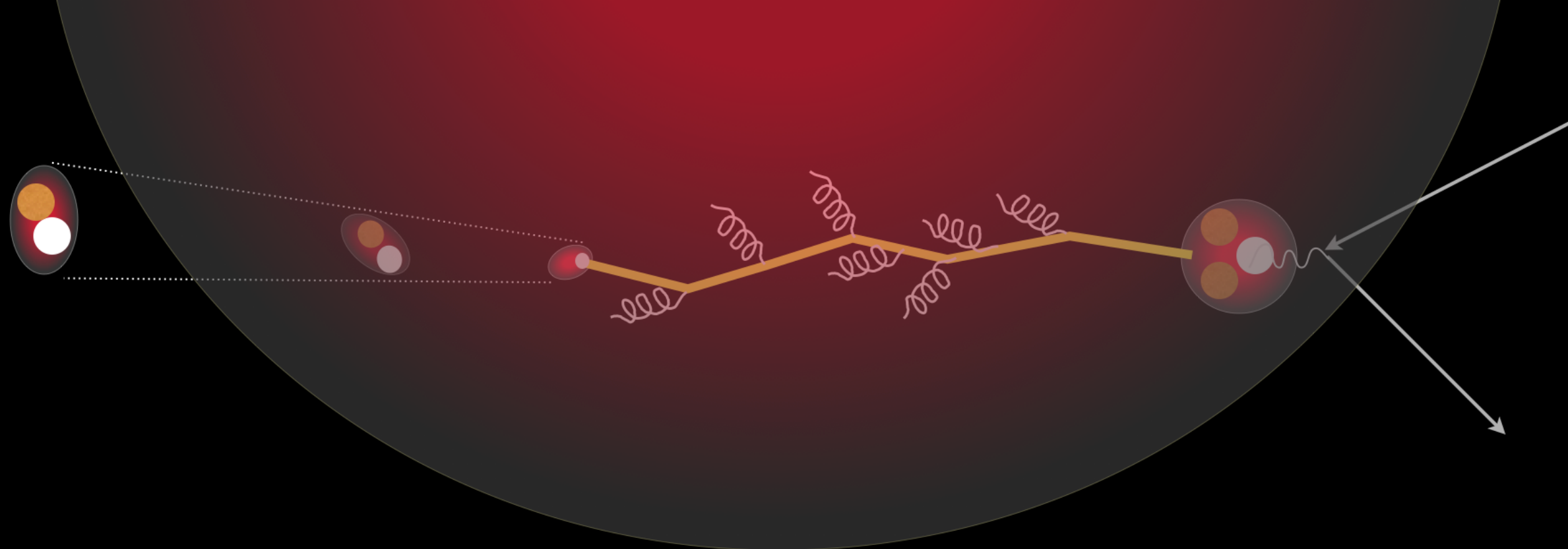
# Lund String Model (~1983)



Remarkably successful model, foundational tool in HEP

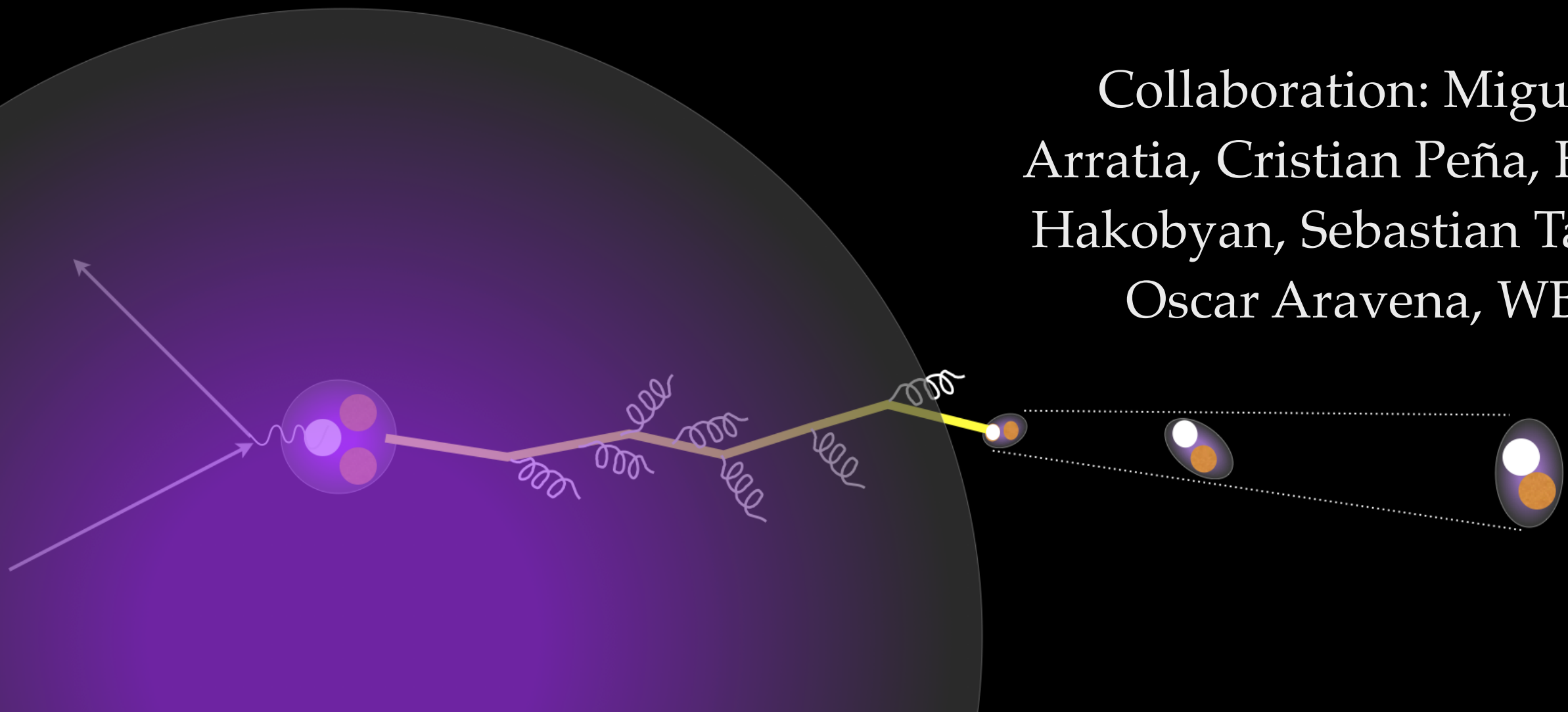
- Alternative physical picture to pQCD: emission of many gluons in vacuum, string as an average; quantitative
- Successful, but few connections to fundamental QCD
- We can *compare* some of our results to the Lund String Model, and other results to pQCD/Color Dipole Model

*We use none of these in our analysis*



# Direct measurement of quark energy loss

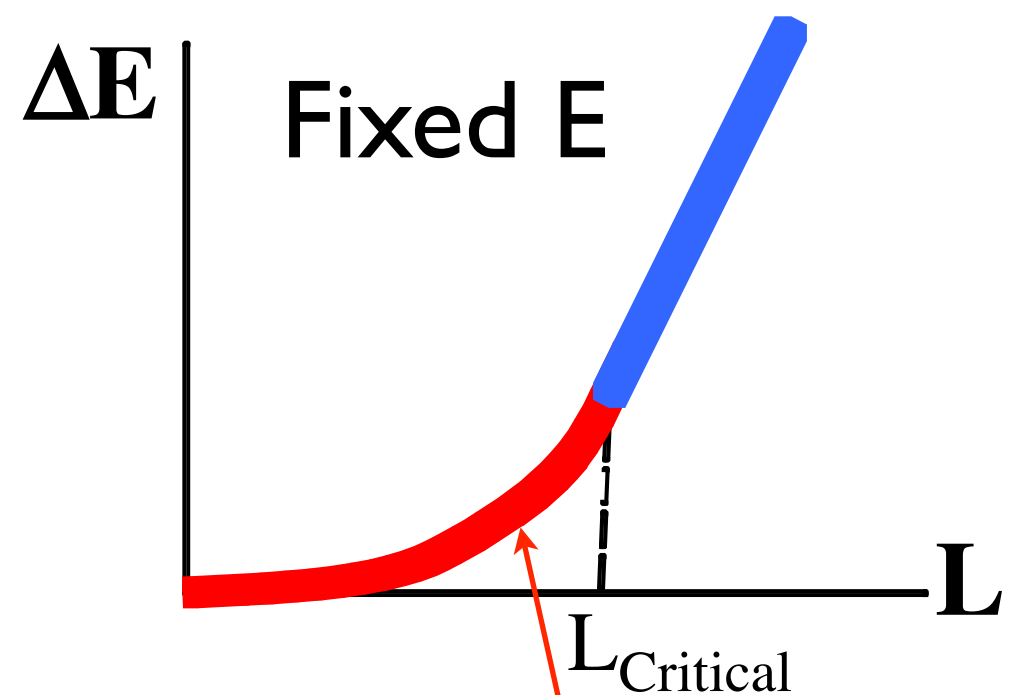
Collaboration: Miguel Arratia, Cristian Peña, Hayk Hakobyan, Sebastian Tapia, Oscar Aravena, WB





$$L < L_{\text{Critical}} \quad -\frac{dE}{dx} \propto L \hat{q}$$

$$L > L_{\text{Critical}} \quad -\frac{dE}{dx} \propto \sqrt{E \hat{q}}$$

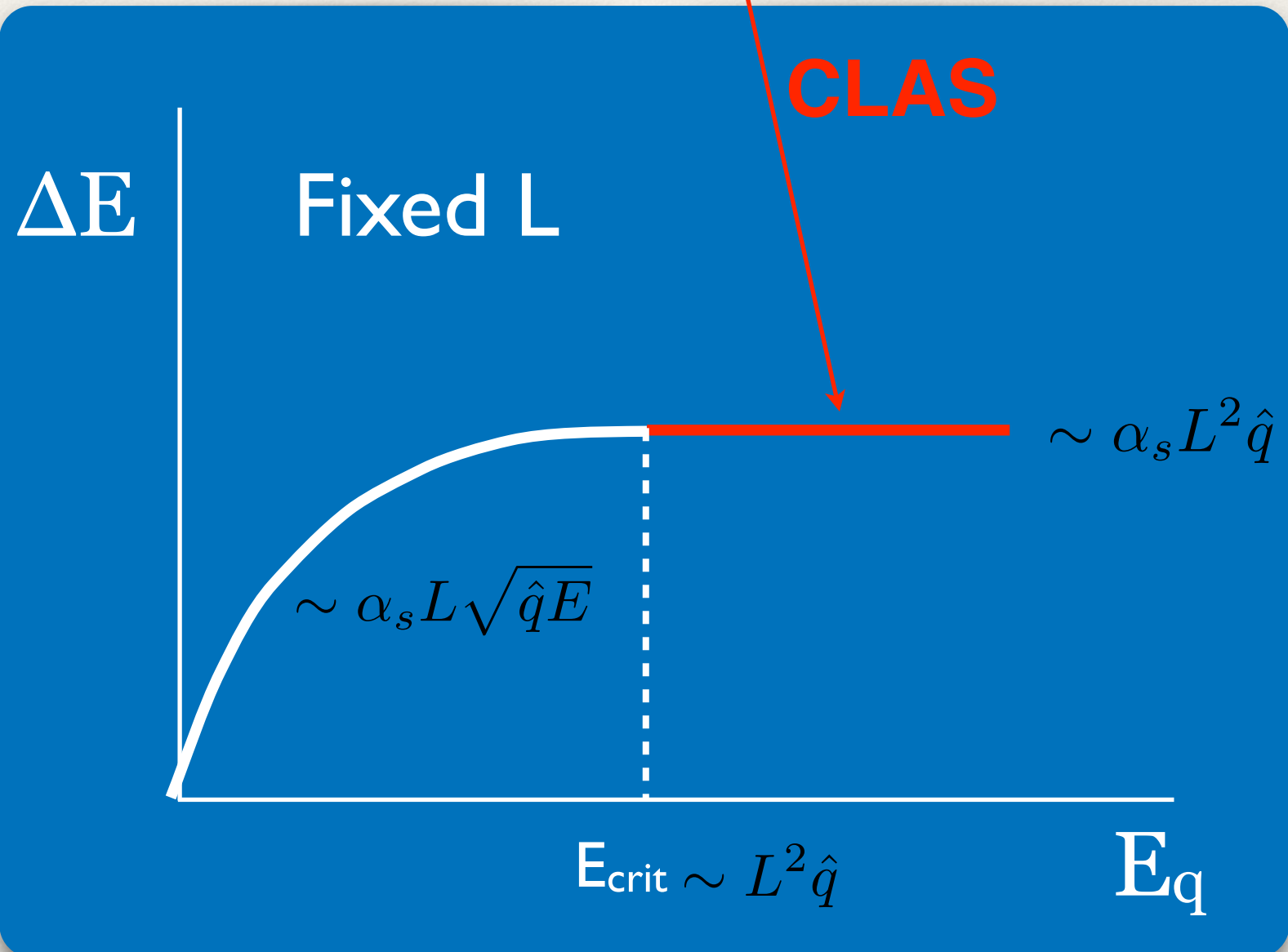


**Partonic energy loss** in pQCD (BDMPS-Z) exhibits a critical system length  $L_c$  and a critical energy  $E_c$

$$L_c \propto \sqrt{\frac{E_q}{\hat{q}}}$$

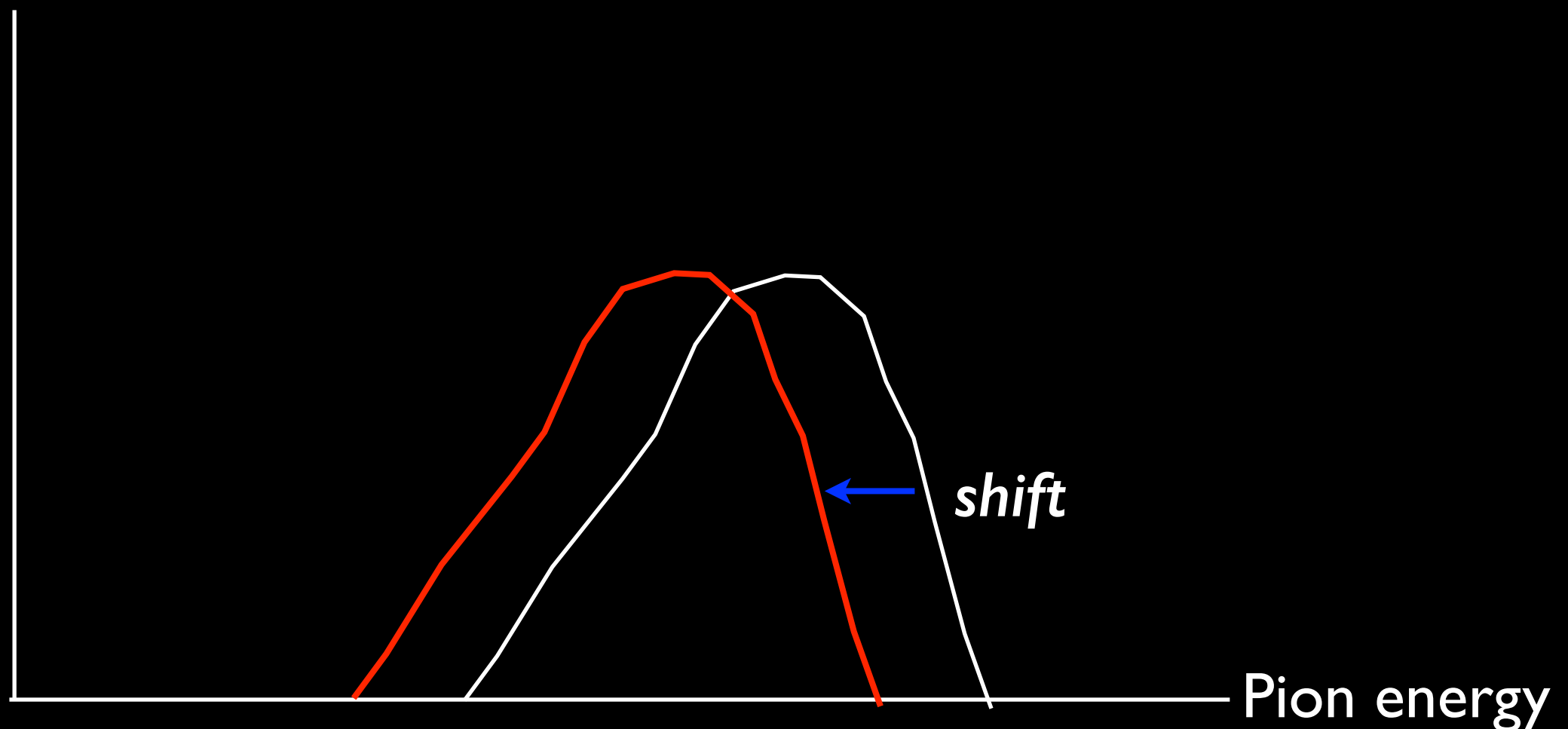
$$E_c \approx 0.4 \cdot \left(\frac{L}{1 \text{ fm}}\right)^2 \text{ GeV}$$

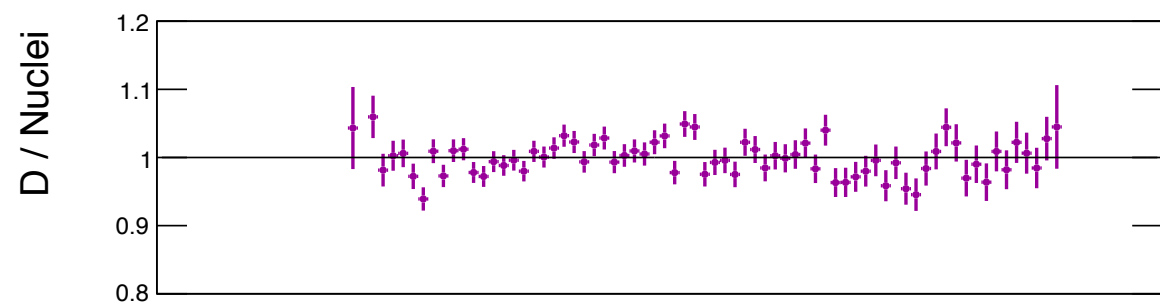
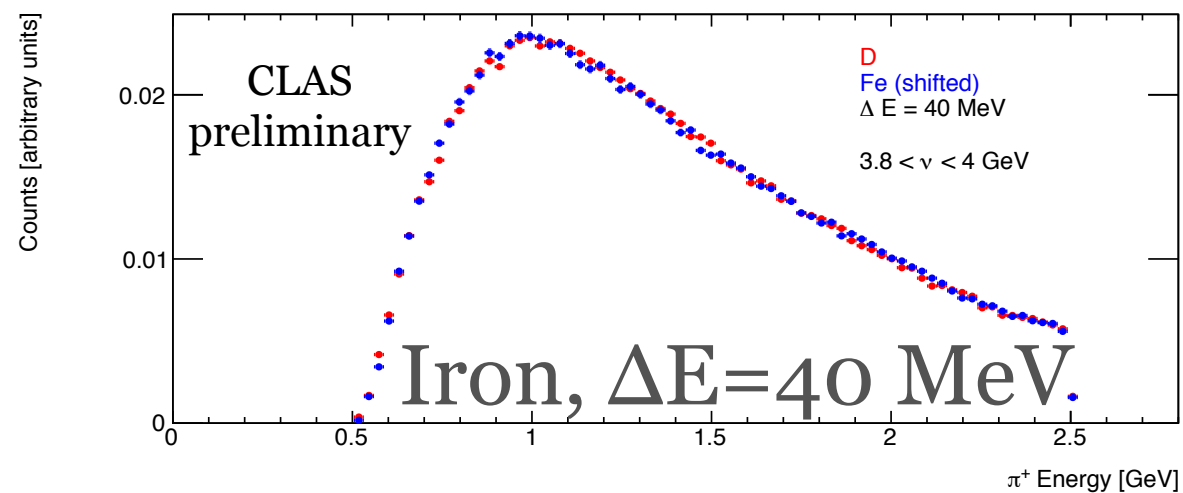
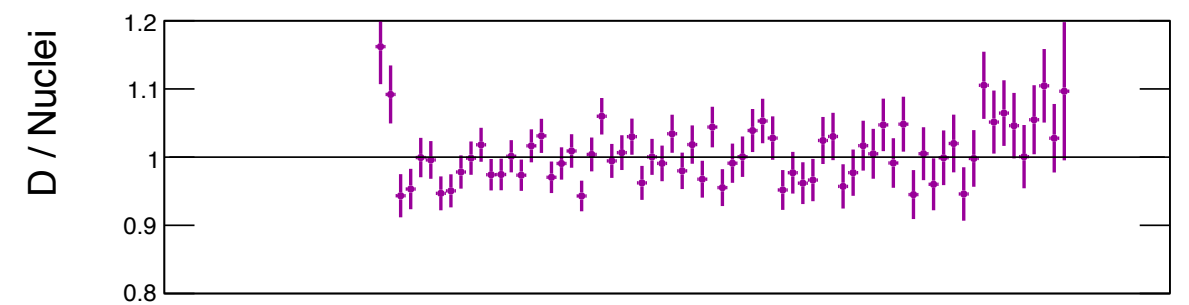
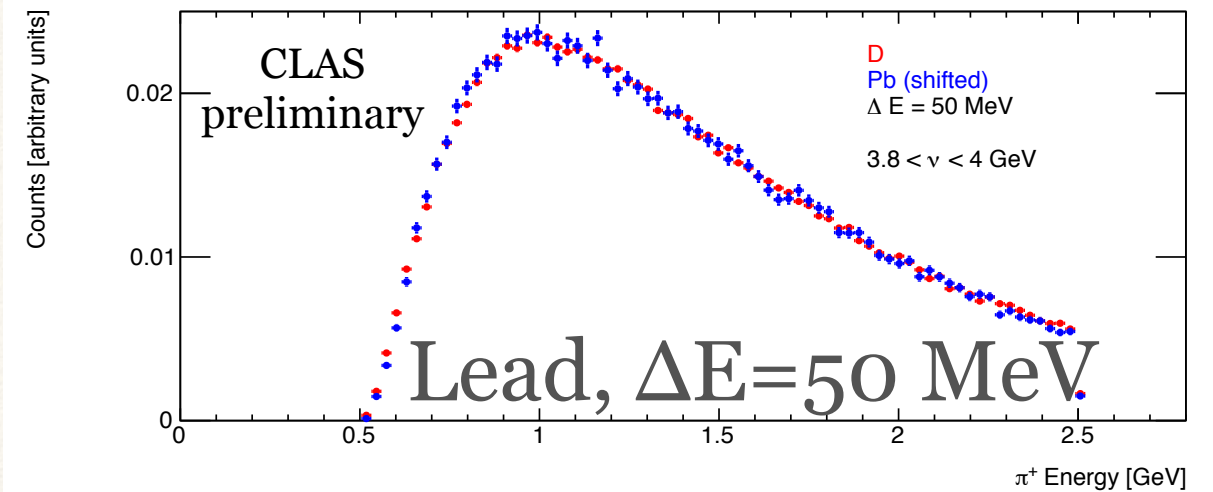
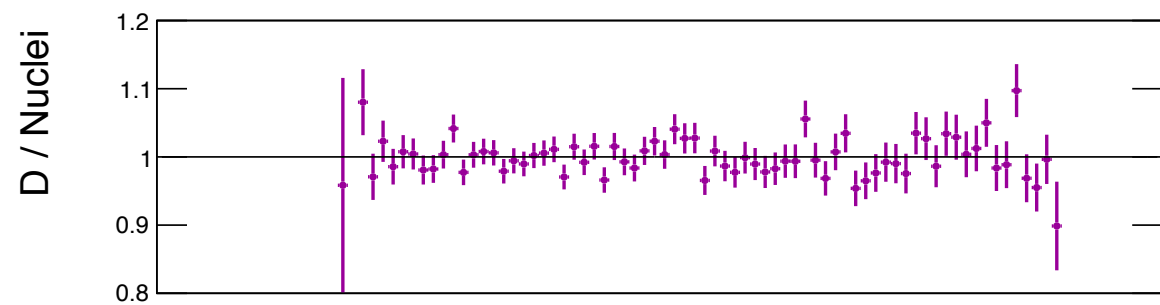
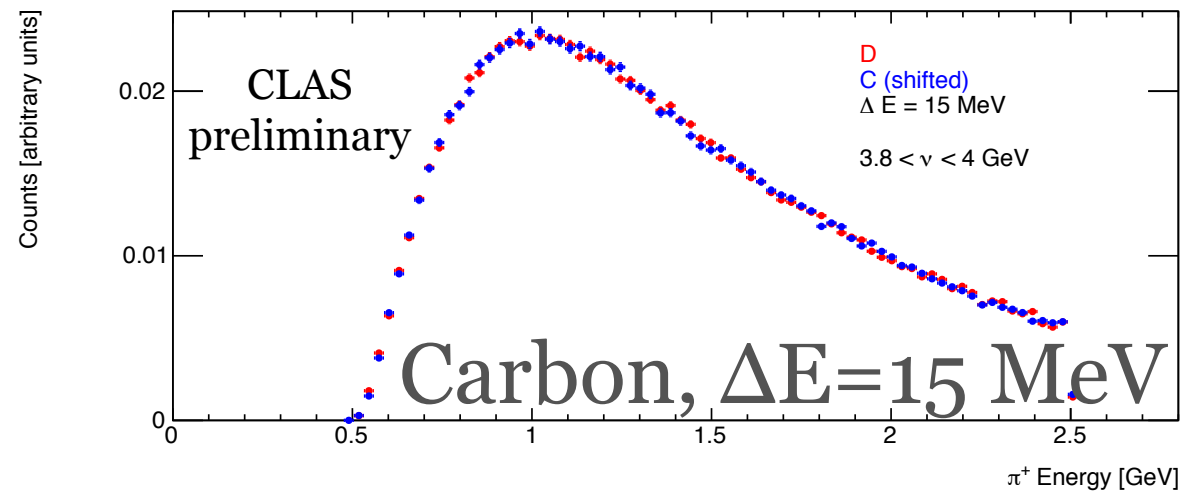
$$-\Delta E_q = \frac{\alpha_s}{4} \Delta k_T^2 \cdot L = \frac{\alpha_s}{4} \hat{q} \cdot L^2$$



# How to *directly* measure quark energy loss?

- Energy loss: *independent of energy* for thin medium
- “Thin enough” = quark energy should be greater than  $E_{\text{crit}}$
- If energy loss is independent of energy, it will produce a **shift** of the energy spectrum, *without distortion of shape*.
- We can look for a **shift** of the Pb energy spectrum compared to that of the deuterium energy spectrum





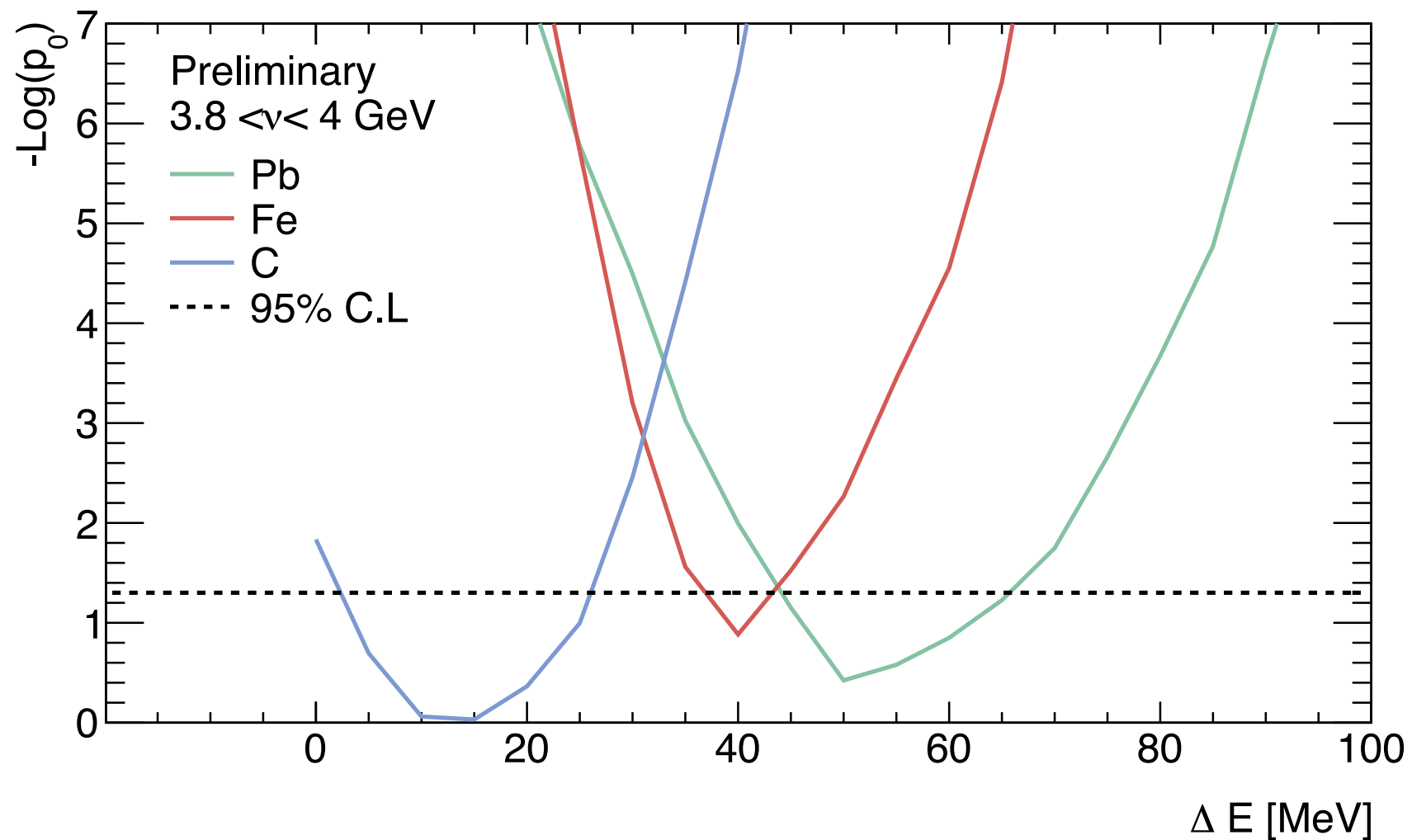
Energy spectrum of  $\pi^+$  produced in C, Fe, Pb compared to that of deuterium, normalized to unity, with energy shifted by  $\Delta E$ .

Acceptance corrected

Cut on  $X_F > 0.1$  is applied

Consistent with simple energy shift + unchanged fragmentation





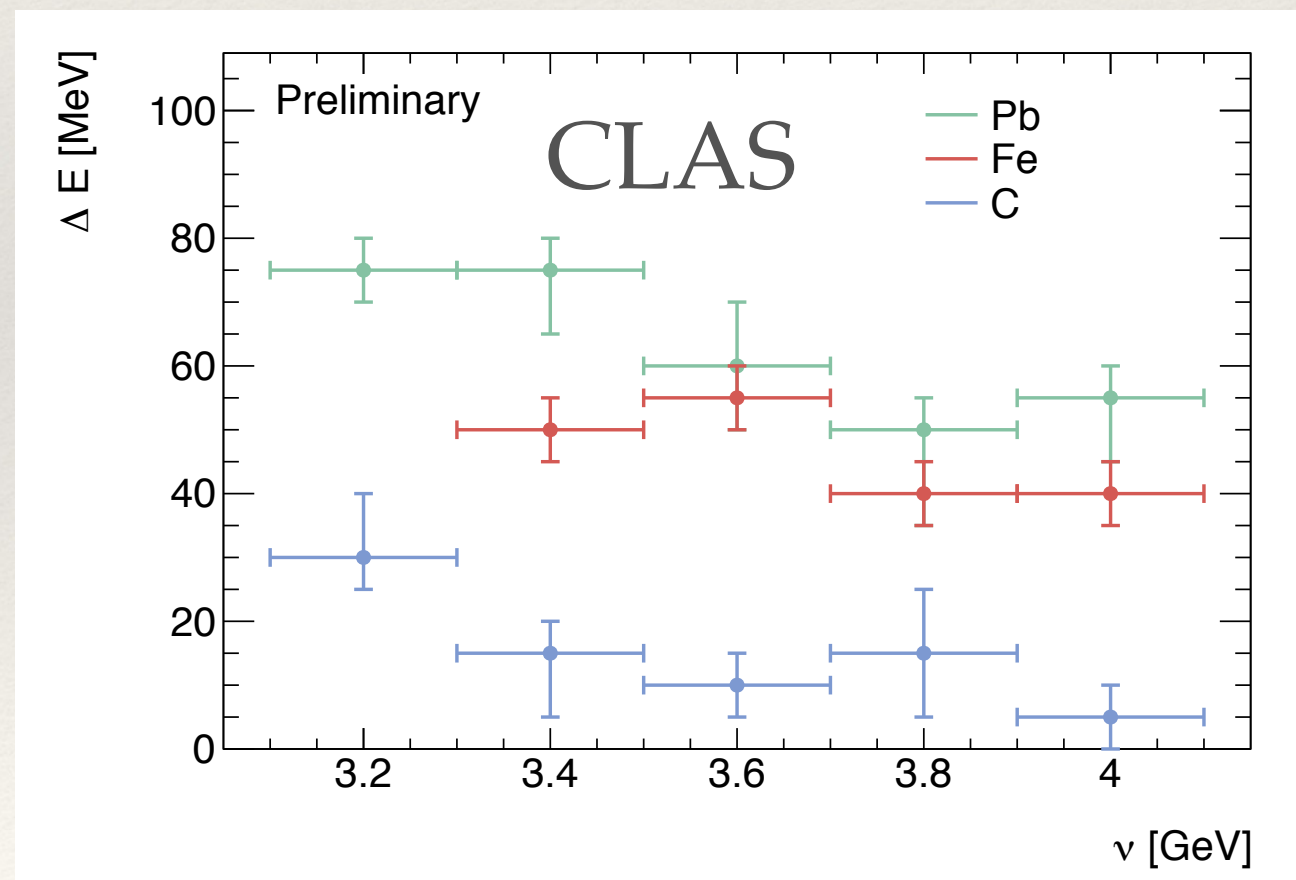
Log of p-values of Kolmogorov-Smirnov test as a function of energy shift  $\Delta E$ : carbon, iron, lead.

Dashed line corresponds to 95% confidence level



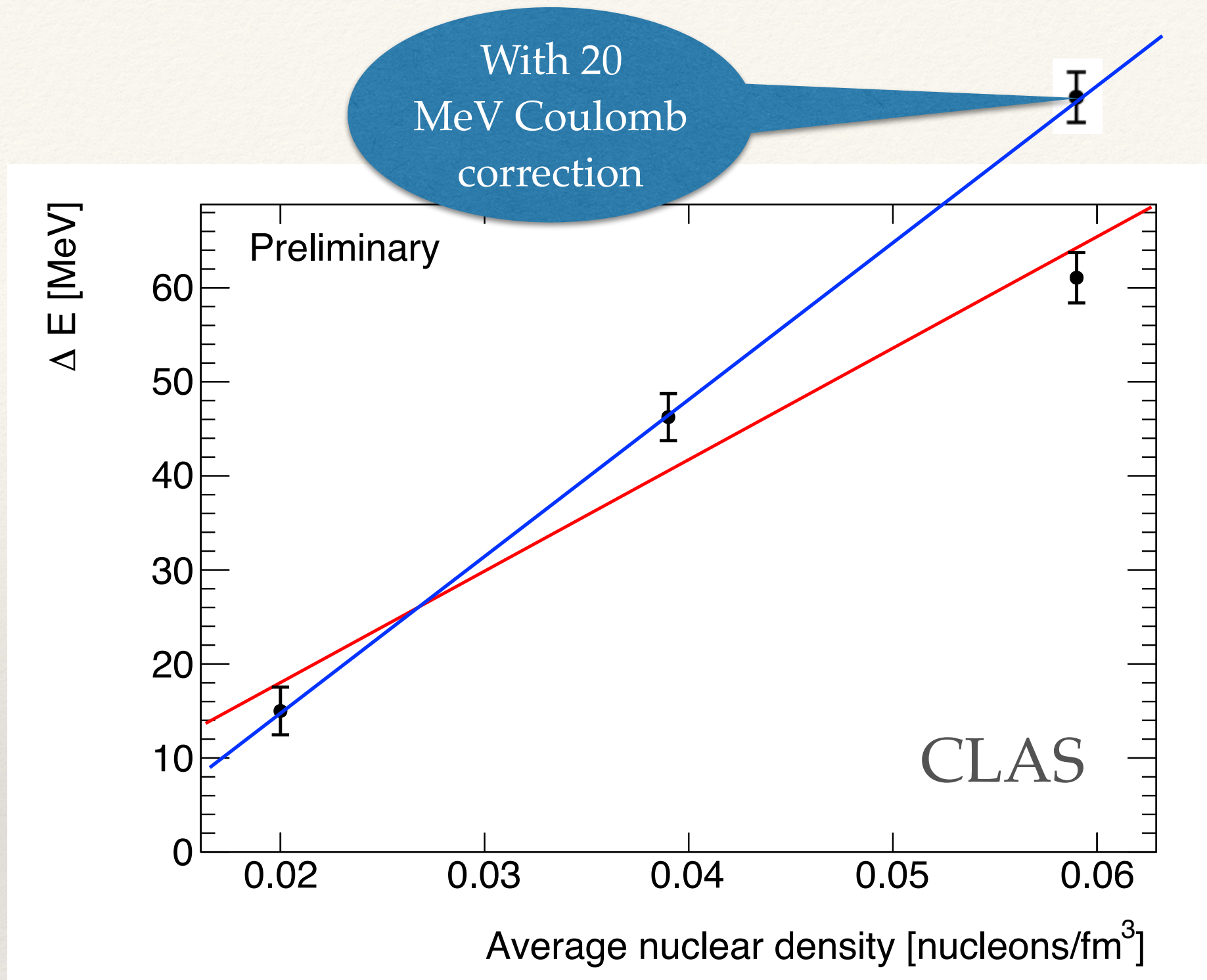
$\nu/\text{GeV}$	Carbon	Iron	Lead
2.4–2.6	—	—	—
2.6–2.8	—	—	—
2.8–3.0	—	—	—
3.0–3.2	—	—	—
3.2–3.4	20–35	—	75
3.4–3.6	10–25	50	70–85
3.6–3.8	10–25	55	50–70
3.8–4.0	5–25	40	45–65
4.0–4.2	5–10	35–40	50–65

Range of possible energy shift in MeV obtained by Kolmogorov-Smirnov test in  $\nu$  intervals



Approximately independent of quark energy, as expected from pQCD for  $\nu > E_{\text{crit}}$





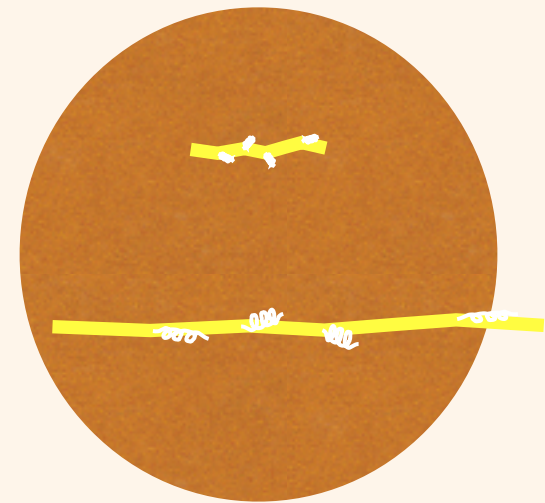
Approximately proportional to density, as expected.  
(fixed pathlength)

Supports the premise that what we measure is ~energy loss!

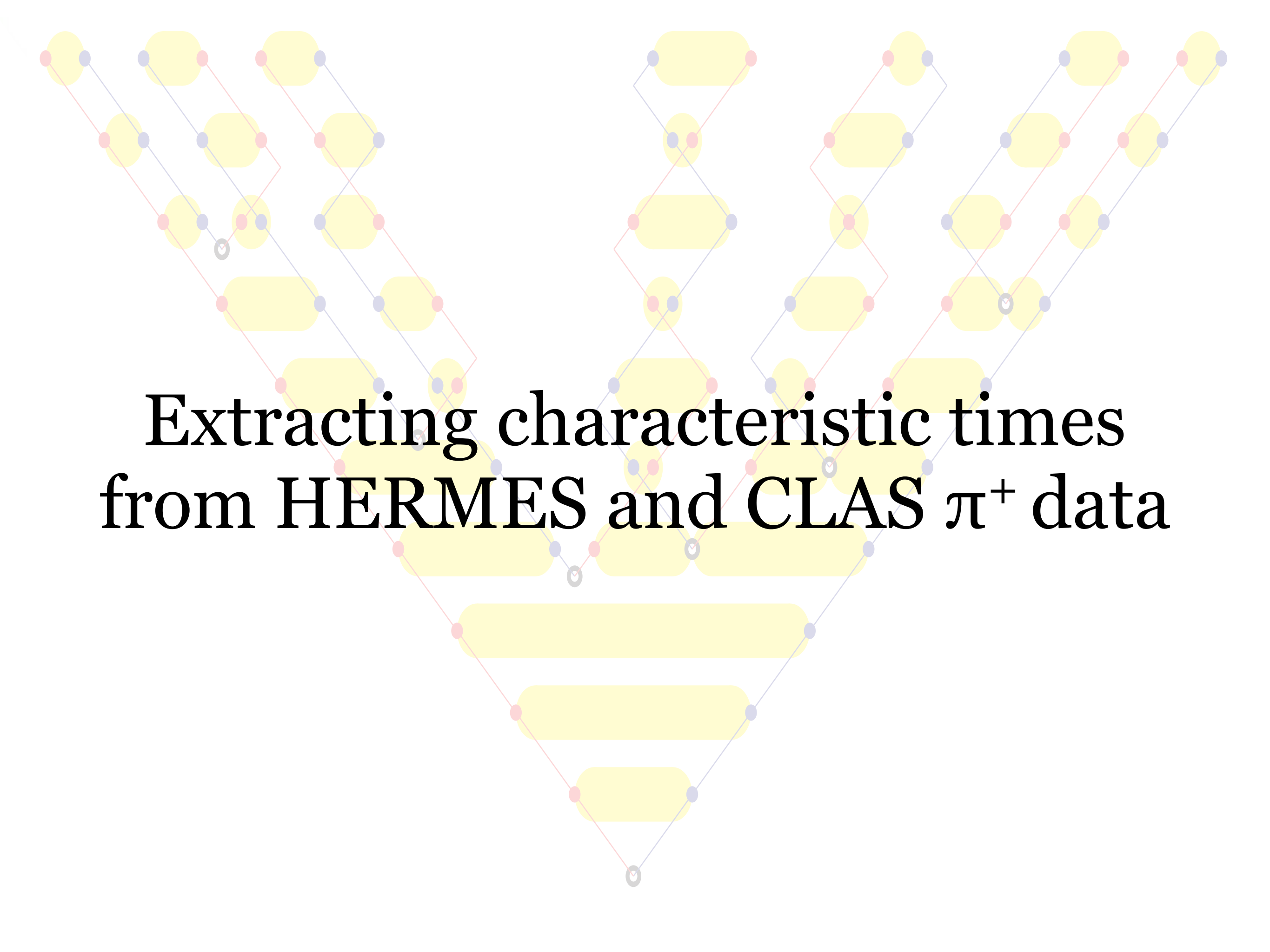


# Direct Measurement of Quark Energy Loss in CLAS: Conclusions

- It is small in magnitude. Why?
  - Best explanation: *short production time*
  - $>500$  MeV vs. 50 MeV in Pb
- It increases with nuclear size. Why?
  - Best explanation: *average nuclear density increases.*
  - Rate of change of virtuality nearly the same in all nuclei, therefore:
    - Path length is short,  $\sim$ independent of nuclear size
    - Nuclear medium has little effect - simple to extrapolate to the vacuum case



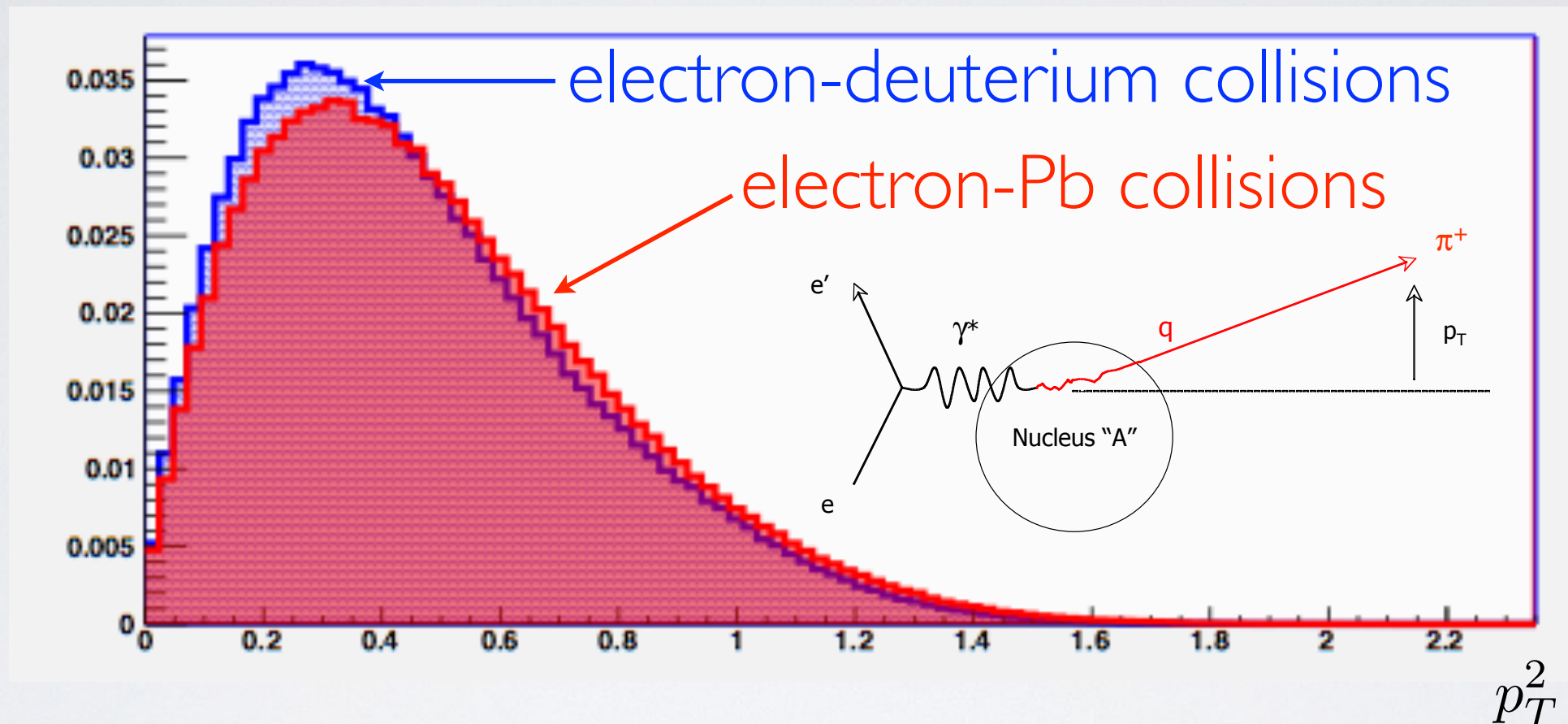
More insights on these ideas in the next section!

The background of the slide features a complex network diagram. It consists of numerous yellow nodes, some of which are circular and others are elongated rectangular shapes. These nodes are interconnected by a web of thin lines, colored in red and blue. The overall structure of the network is somewhat symmetrical and branching, resembling a tree or a complex web. In the center of the image, there is a large text overlay in a black serif font. The text reads: "Extracting characteristic times from HERMES and CLAS  $\pi^+$  data".

Extracting characteristic times  
from HERMES and CLAS  $\pi^+$  data

Observable:  $p_T$  broadening

$$\Delta p_T^2 \equiv \langle p_T^2 \rangle_A - \langle p_T^2 \rangle_D$$



$p_T$  broadening is a tool: sample the gluon field using a colored probe:

$$\Delta p_T^2 \propto G(x, Q^2) \rho L$$

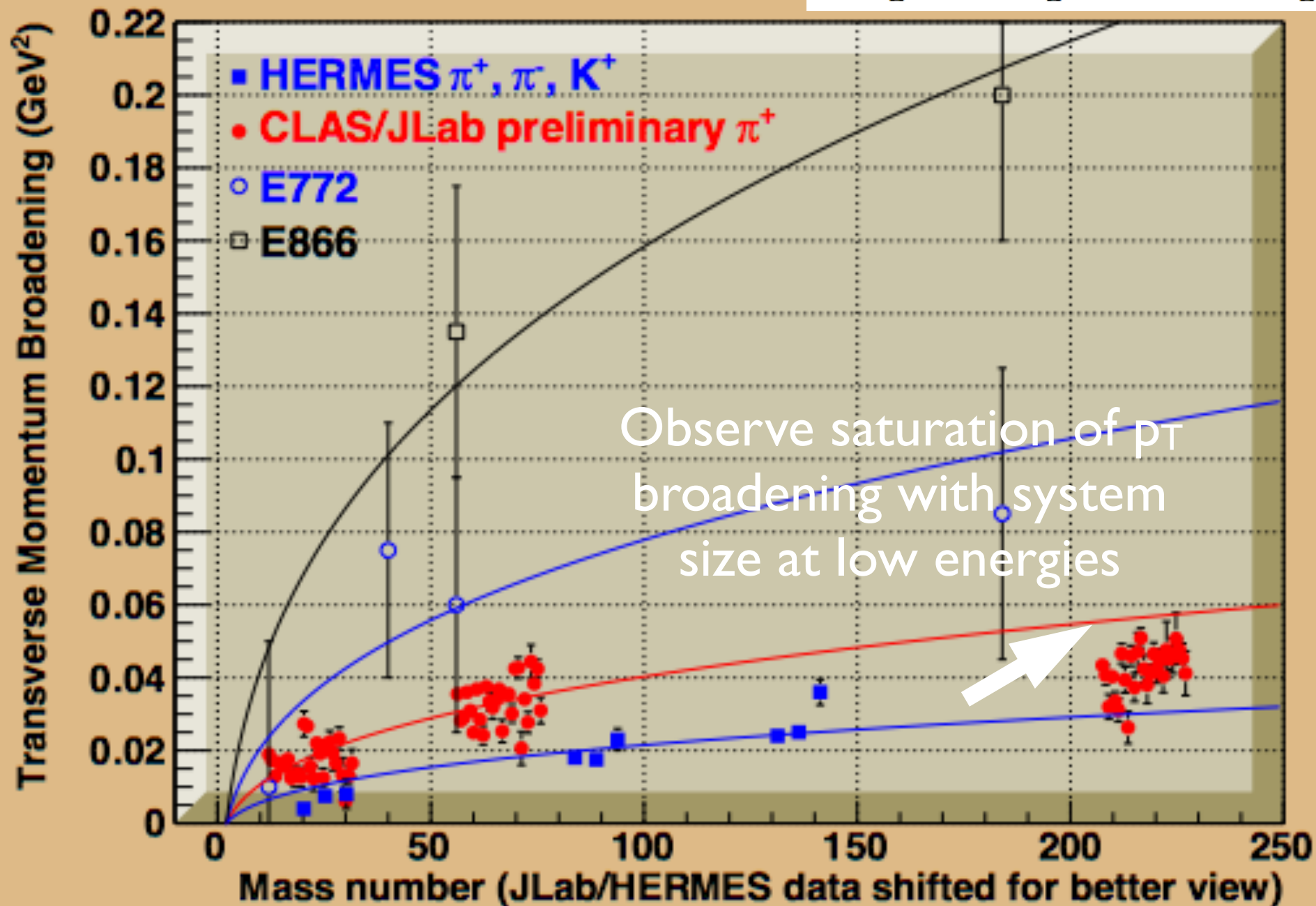
and radiative energy loss:

$$-\frac{dE}{dx} = \frac{\alpha_s N_c}{4} \Delta p_T^2$$



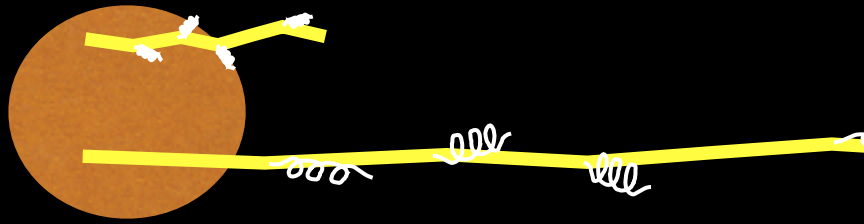
# $p_T$ broadening data - Drell-Yan and DIS

$$\Delta p_T^2 = \langle p_T^2 \rangle_A - \langle p_T^2 \rangle_D$$

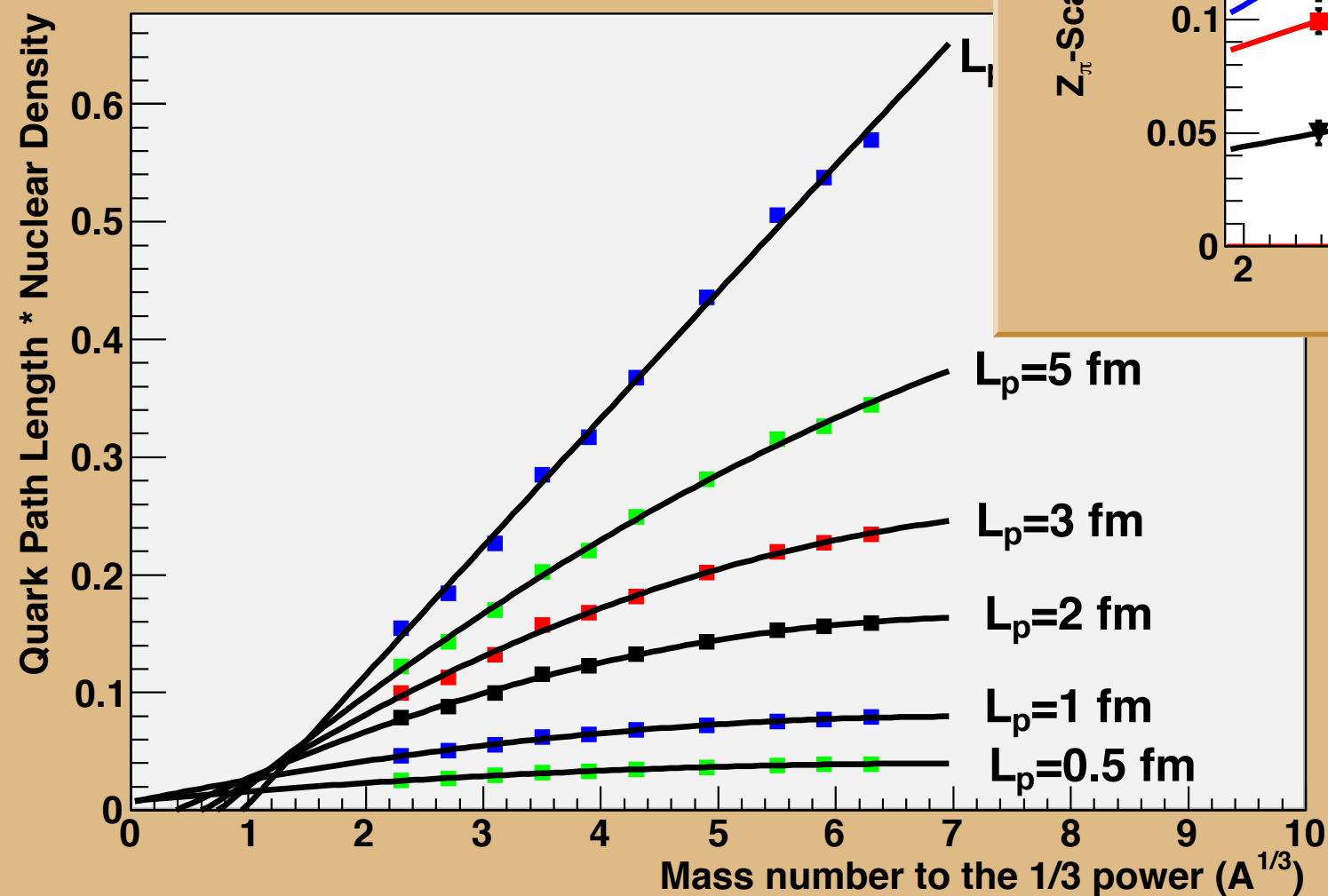


- New, precision data with identified hadrons!
- CLAS  $\pi^+$ : 81 four-dimensional bins in  $Q^2$ ,  $\nu$ ,  $z_h$ , and  $A$
- Intriguing *saturation*: production length or something else?

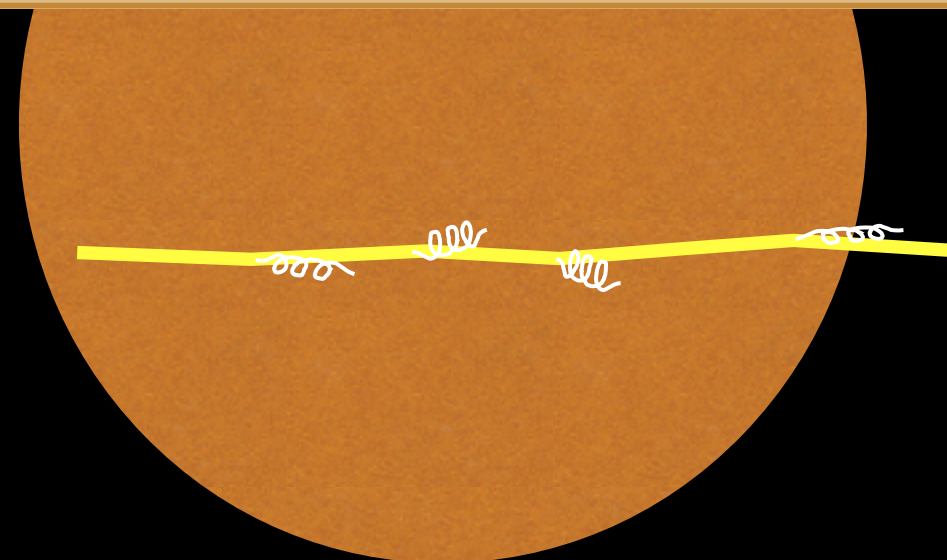
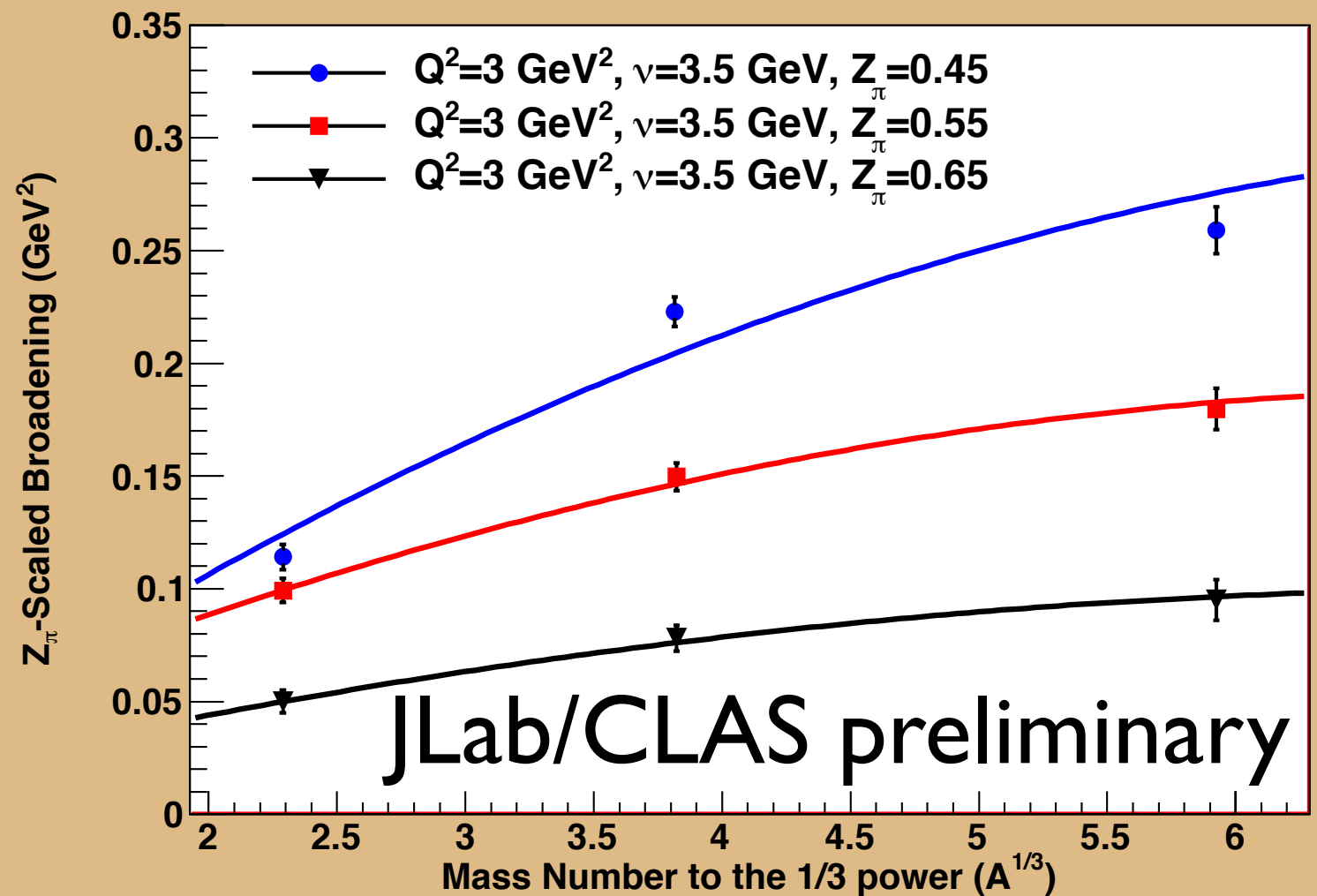
# Production Time Extraction - Geometrical Effects



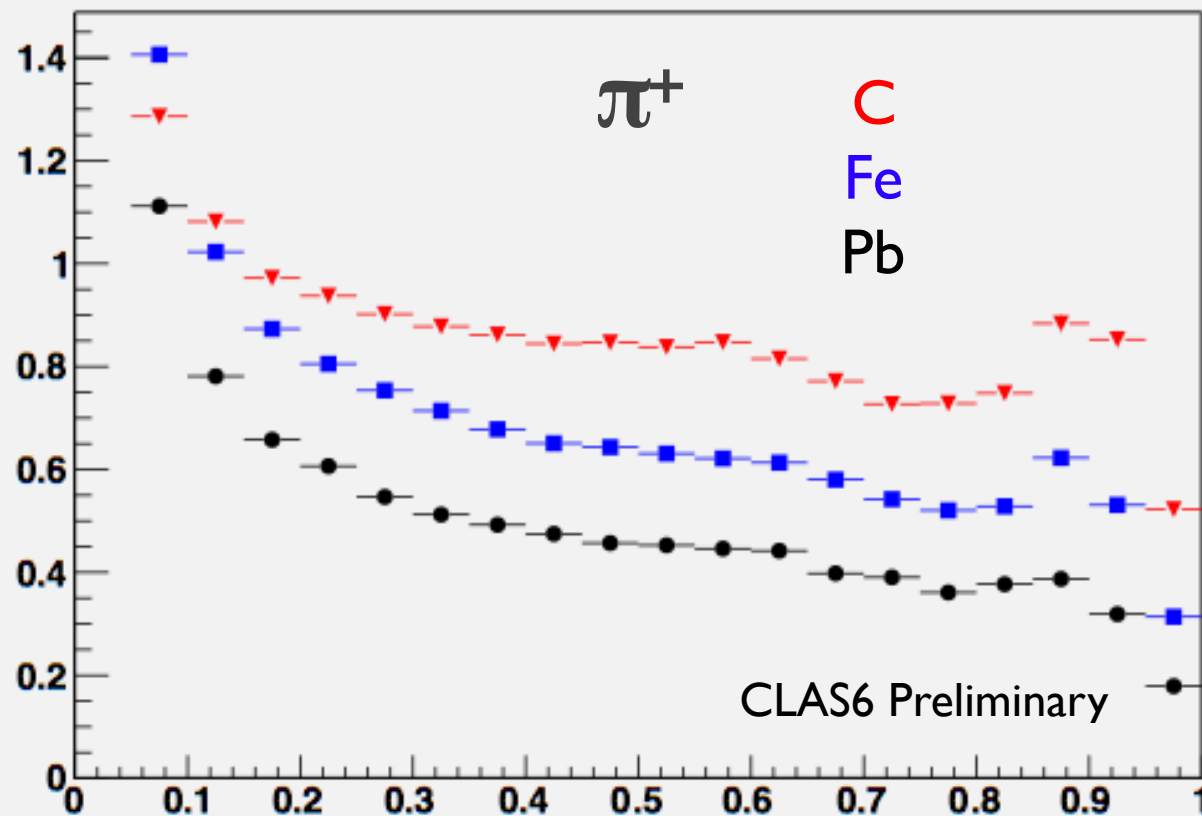
Quark Path Length \* Nuclear Density vs.  $A^{1/3}$



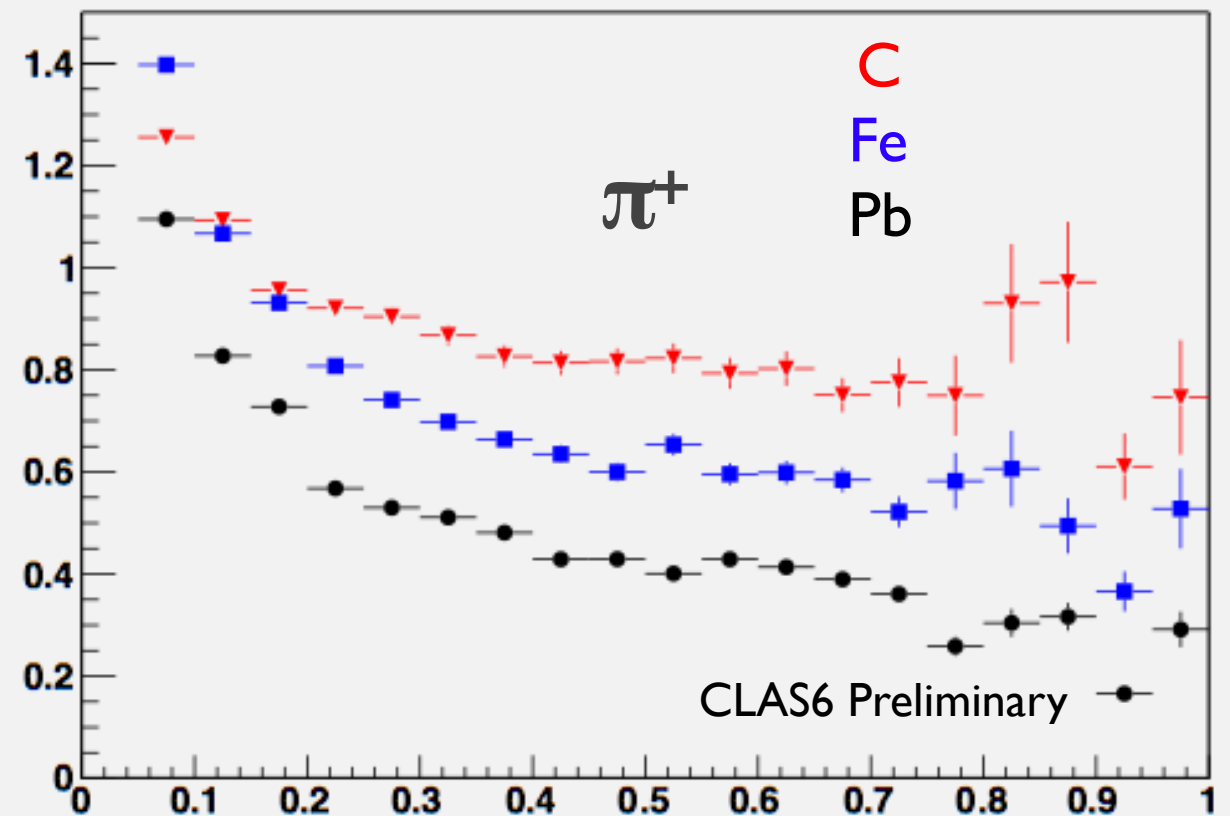
Fits of Z-Scaled Broadening vs.  $A^{1/3}$



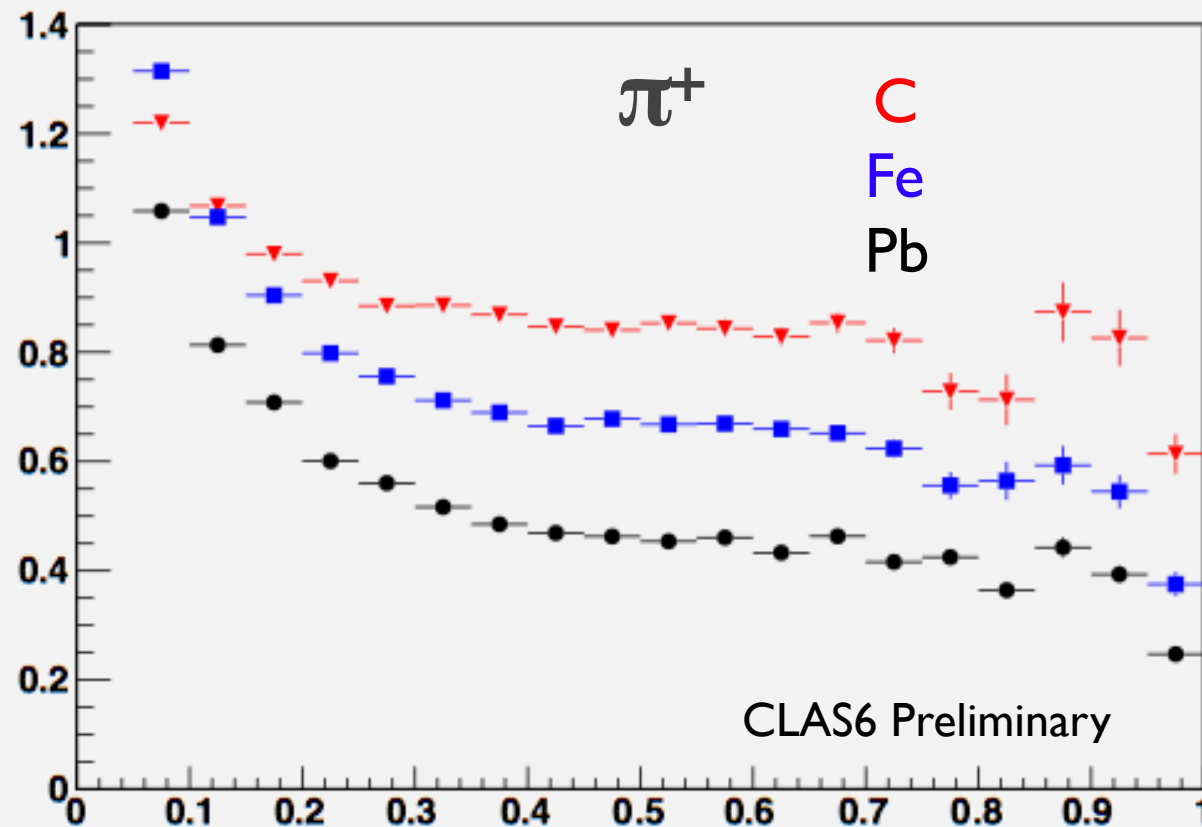
$1.0 < Q^2 < 2.0$   $2.2 < \nu < 2.8$



$3.0 < Q^2 < 4.0$   $3.4 < \nu < 4.0$



$2.0 < Q^2 < 3.0$   $3.4 < \nu < 4.0$



Data from CLAS6 and CLAS12 will provide the ultimate low- $\nu$  studies in up to 4-fold differential multiplicity ratios. EIC will have overlap and will provide the crucial high- $\nu$  studies.

CLAS6:  $\pi^+$  ( $K^0$ ,  $\pi^0$ ,  $\pi^-$ )



# Geometrical Model approach

- Struck quark absorbs energy and momentum of  $\gamma^*$
- Pathlength in medium: *either* parton *or* prehadron
- $p_T$  broadening of final hadron: *only* from partonic multiple scattering,  $\propto L_p$
- Inelastic prehadron interactions ‘attenuate’ produced hadrons,  $\propto L_h$
- Partonic energy loss can also ‘attenuate’ produced hadrons (but we find this is a small effect)
- We extract the dynamical behavior, *separating geometry from dynamics*



# Coupled fit of HERMES data

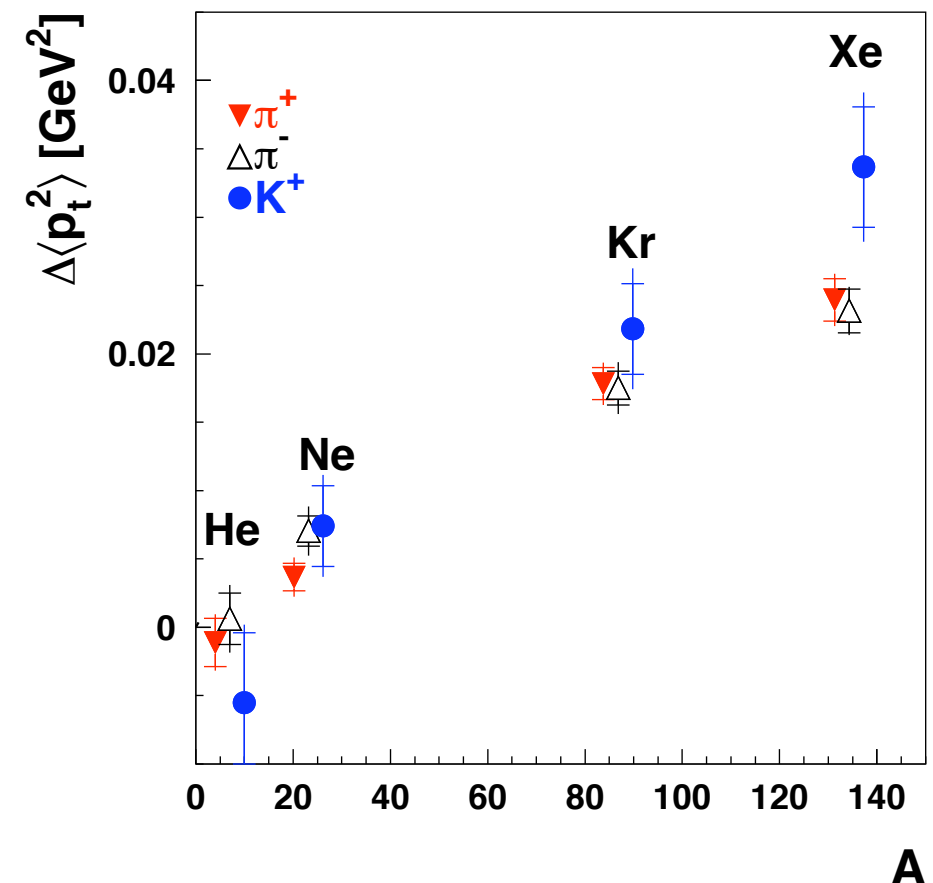
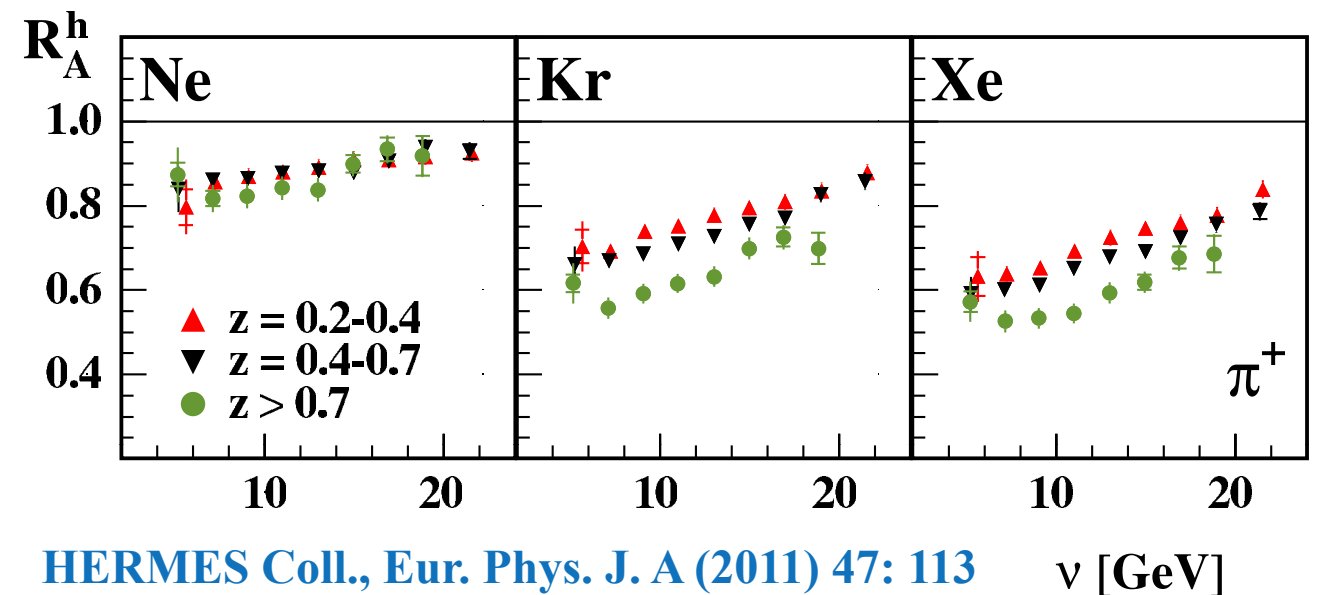
Multiplicity ratio

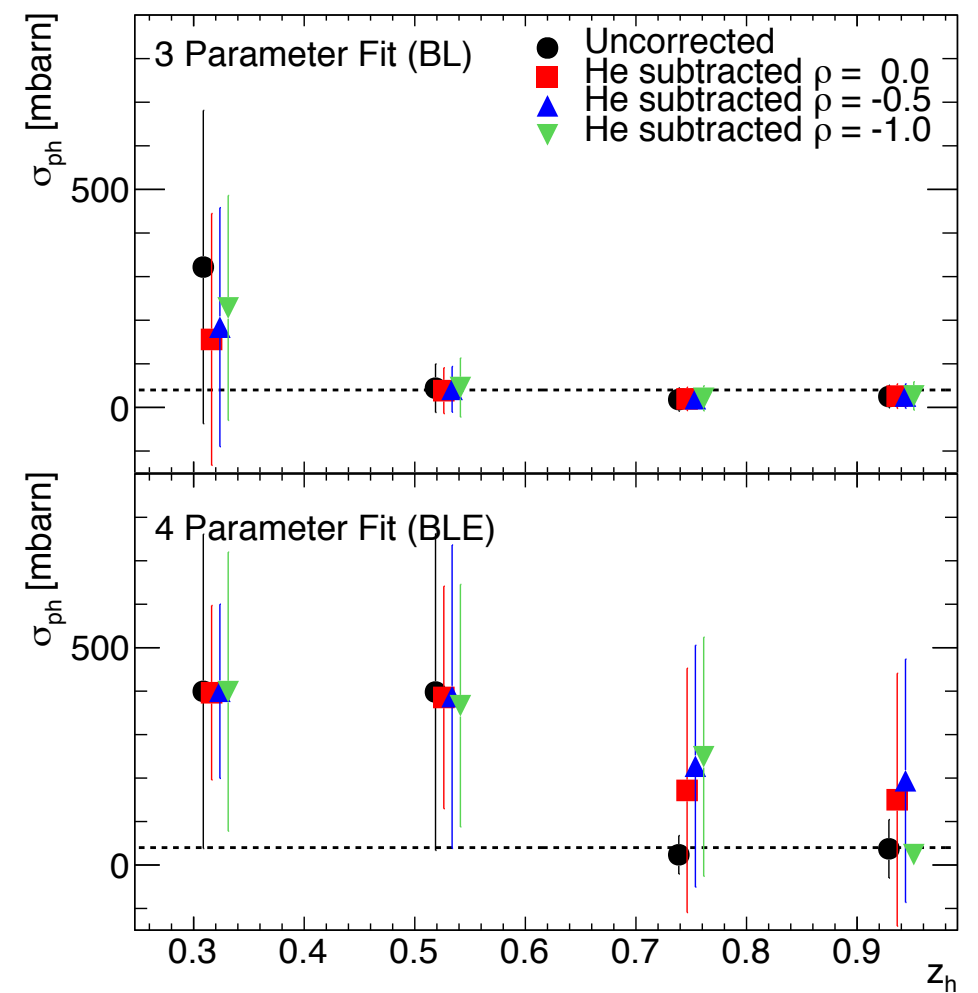
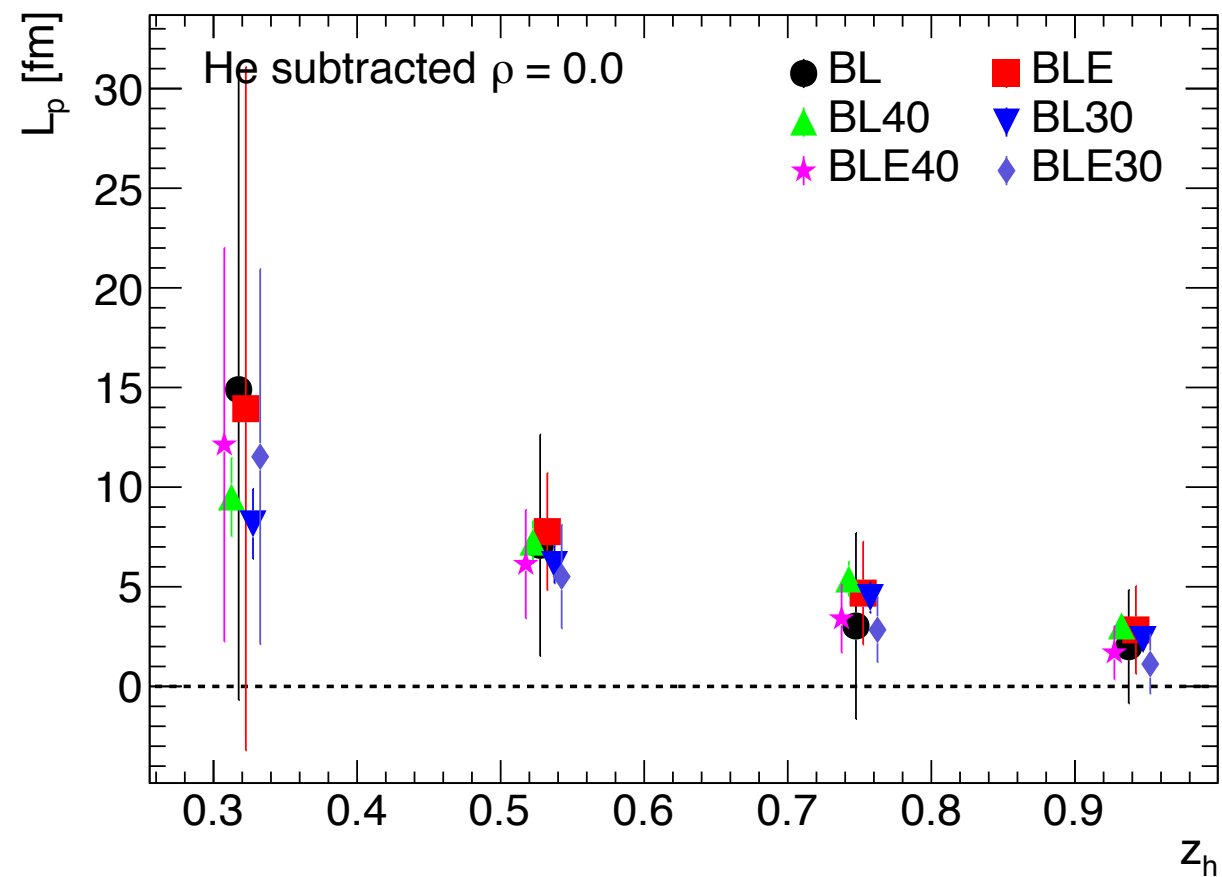
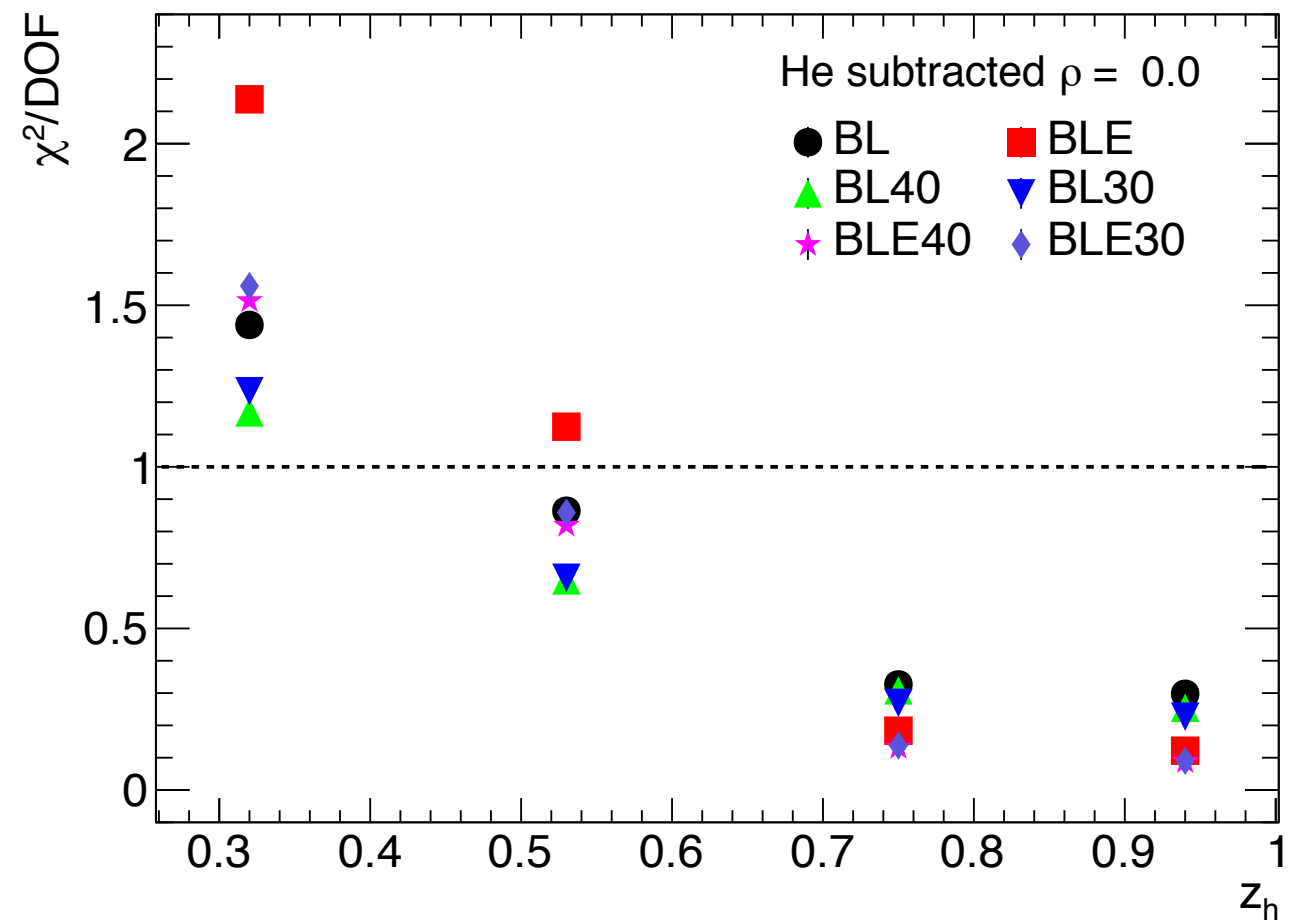
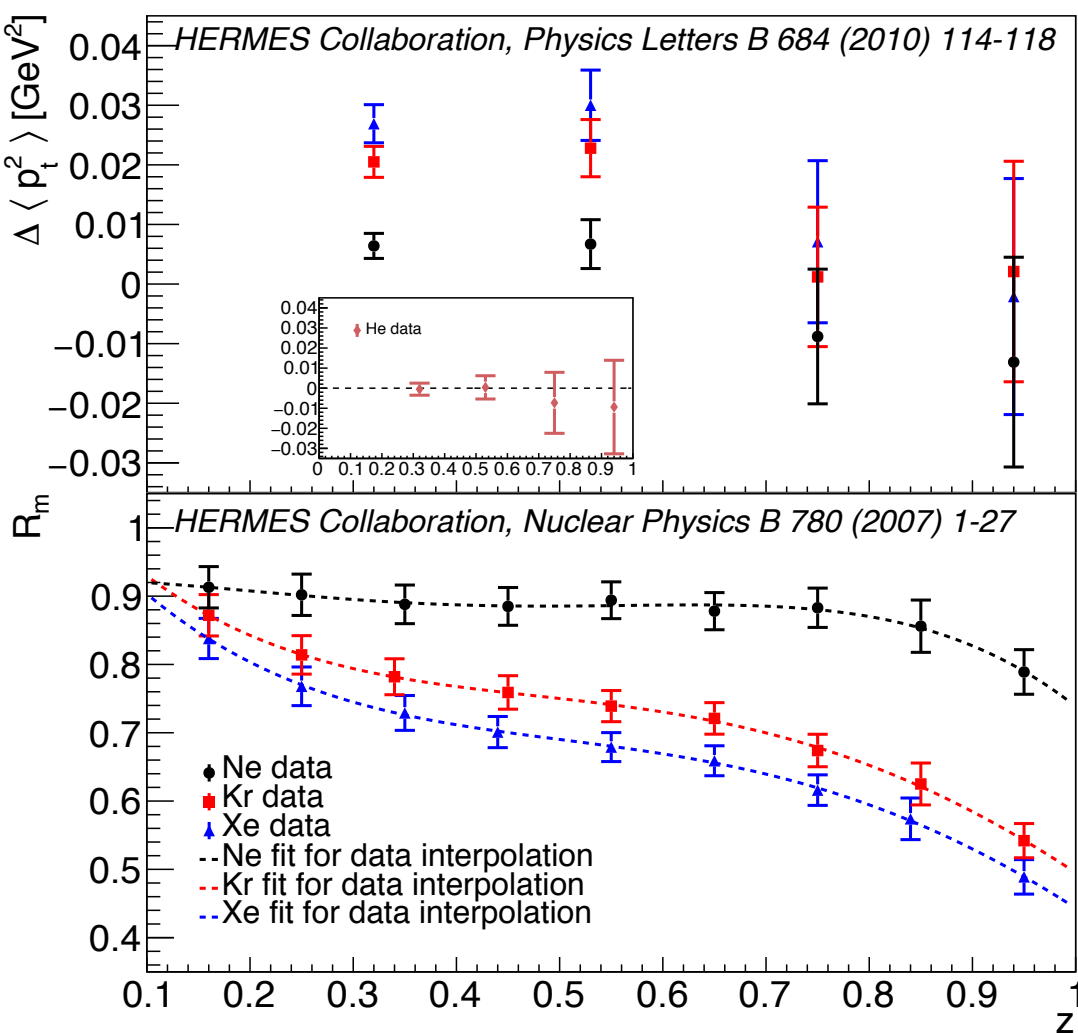
$$R_M^h(Q^2, \nu, z, p_T) \equiv \frac{\frac{1}{N_e(Q^2, \nu)} \cdot N_h(Q^2, \nu, z, p_T)|_A}{\frac{1}{N_e(Q^2, \nu)} \cdot N_h(Q^2, \nu, z, p_T)|_p}$$

$p_T$  broadening

$$\Delta p_T^2(Q^2, \nu, z) \equiv \langle p_T^2(Q^2, \nu, z) \rangle |_A - \langle p_T^2(Q^2, \nu, z) \rangle |_p$$

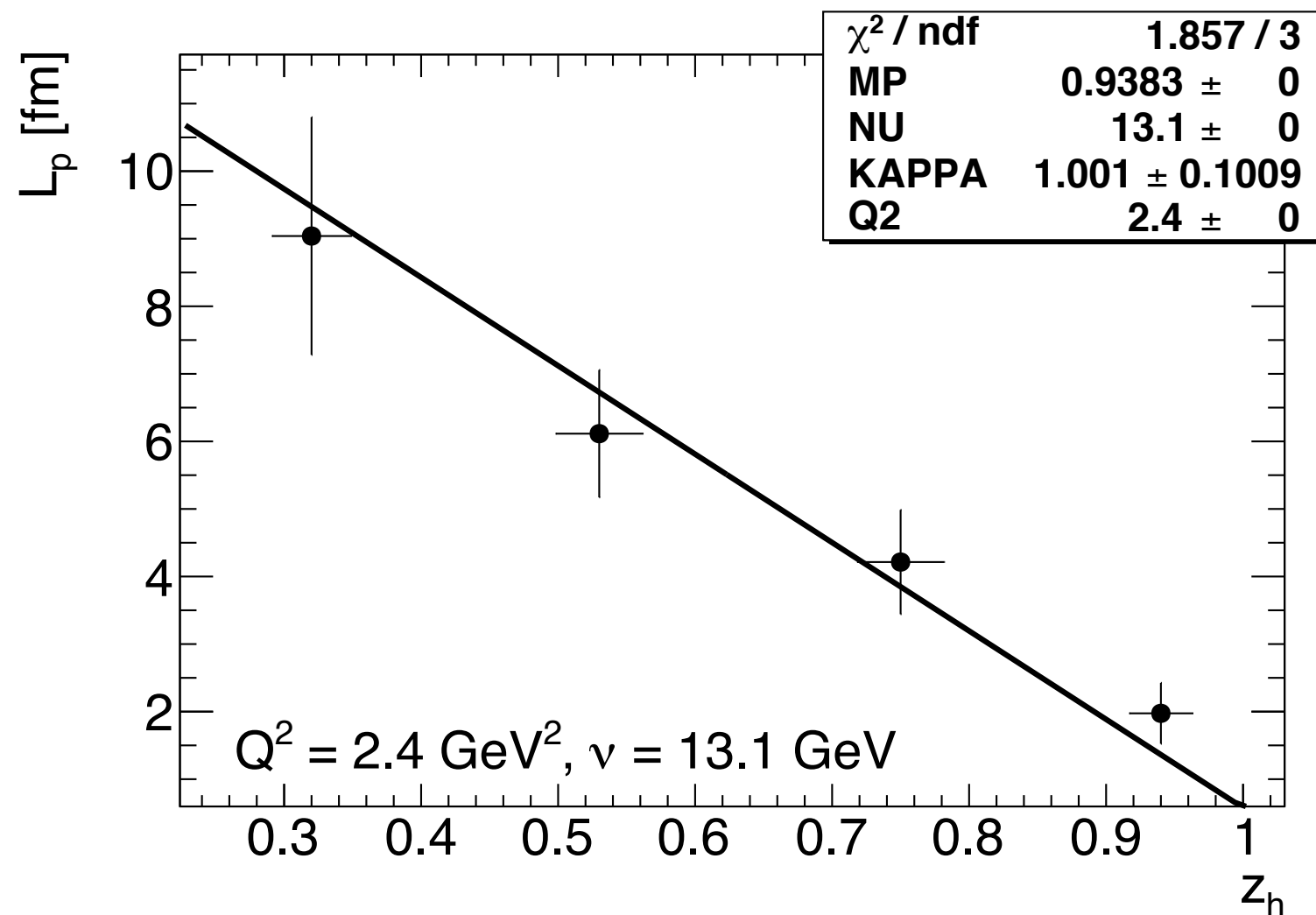
*We fit both observables simultaneously*





# Fit of HERMES $L_p$ results to Lund Model form

A fit of our  
HERMES  
results  
to the  
Lund  
model  
form



This is a  
strong  
validation  
of our  
model

We recover the known value of the string constant  
completely independently!

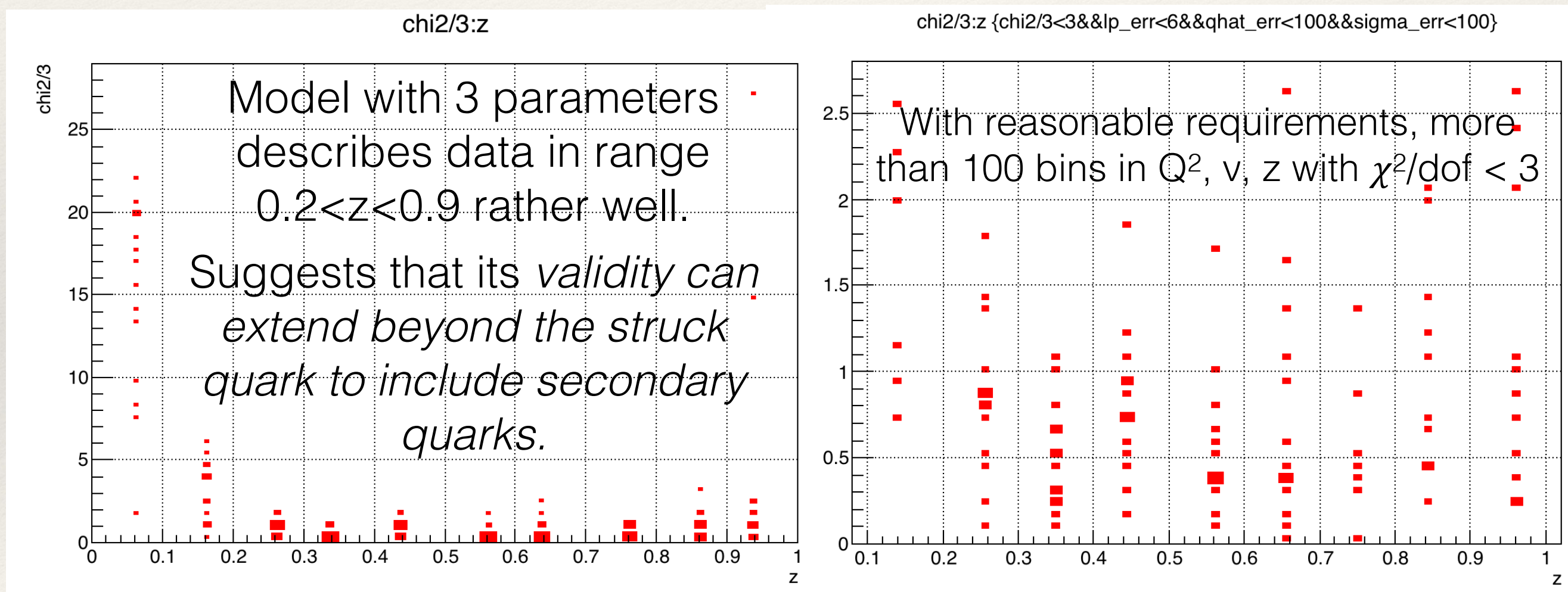
Light cone Lund String Model form for lab frame:

$$l_p = \frac{1}{2\mathcal{K}} \cdot \left( M_p + \nu + \sqrt{\nu^2 + Q^2} - 2 \cdot \nu \cdot z' \right)$$



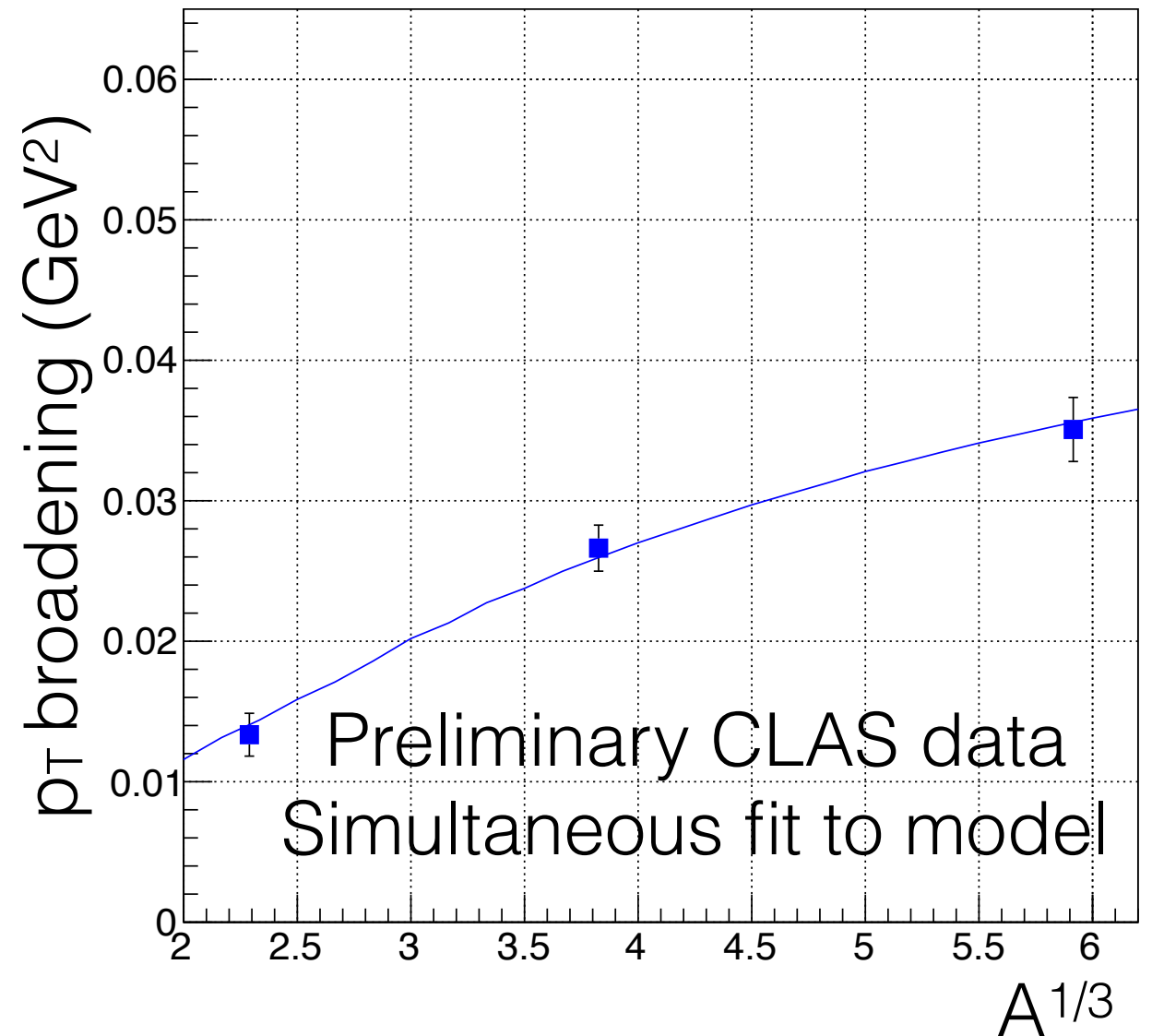
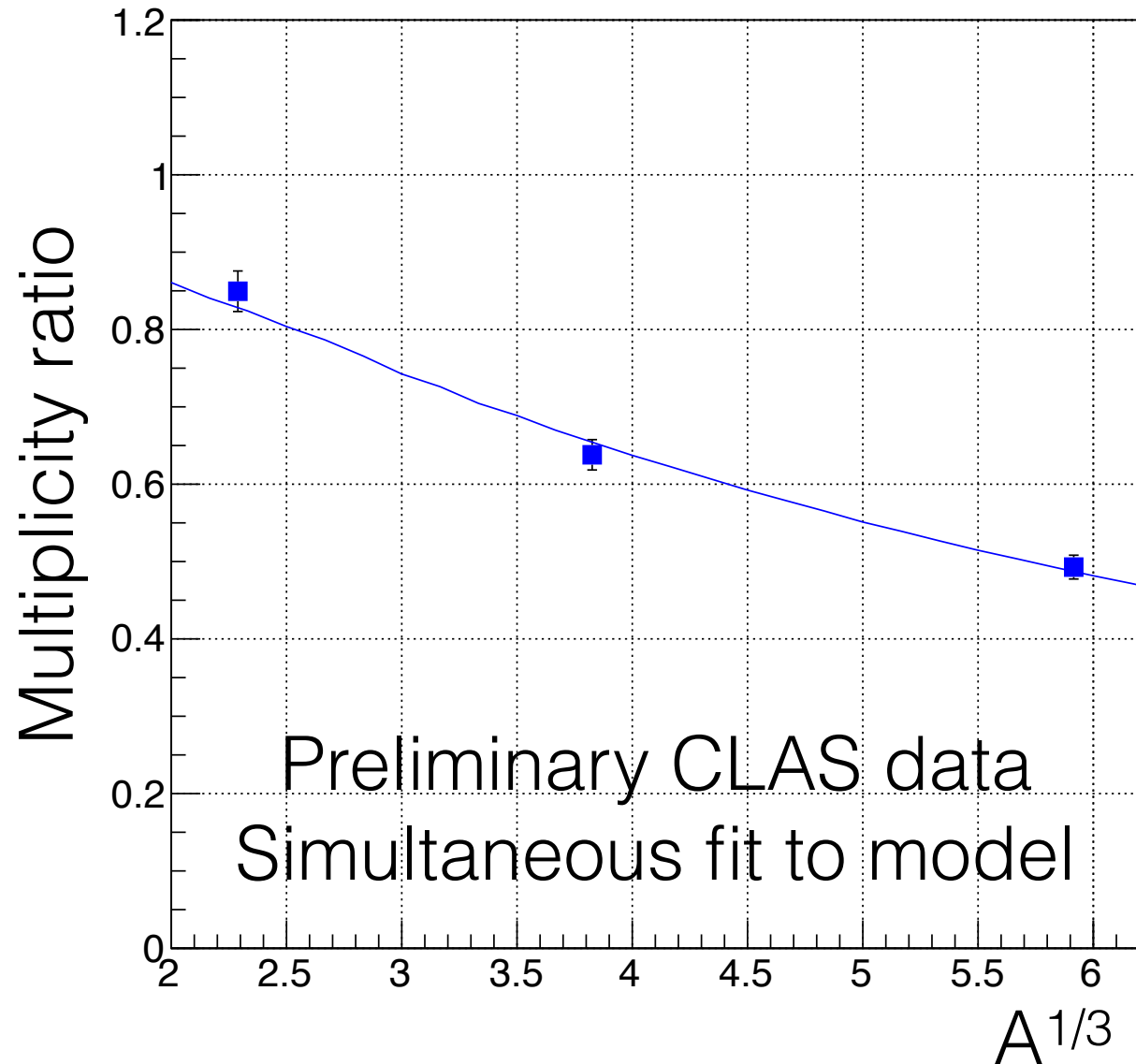
# Virtual quark lifetime extraction: 3-parameter geometric model applied to CLAS 5 GeV data

$\chi^2/\text{dof}$  vs.  $z$



Good values of chi-squared/degrees of freedom for > 100 bins  
*Exploratory study*

# Example of fit (one of 150 bins in $x$ , $Q^2$ , and $z$ )



$\langle x \rangle = 0.166$ ,  $\langle Q^2 \rangle = 1.17 \text{ GeV}^2$ , ( $\langle v \rangle = 3.76 \text{ GeV}$ ),  $\langle z \rangle = 0.445$

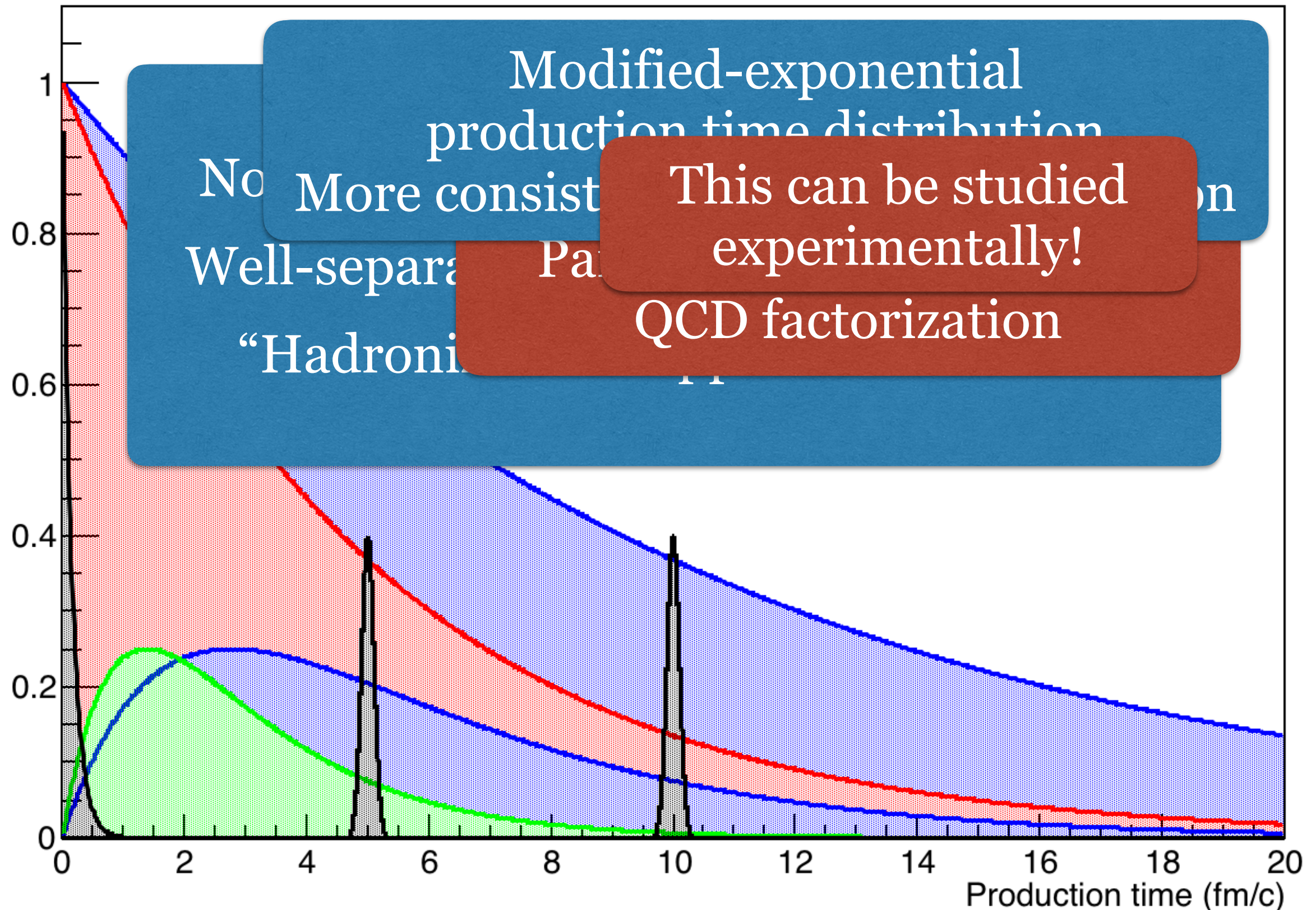
**$L_p = 1.8 \pm 0.4 \text{ fm}$**

$\chi^2/\text{dof} = 0.5$

Simultaneous fit *couples*  $p_T$  broadening to multiplicity ratio

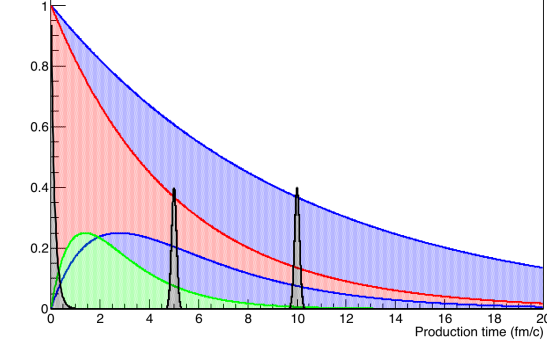


# Three possible distributions of production time

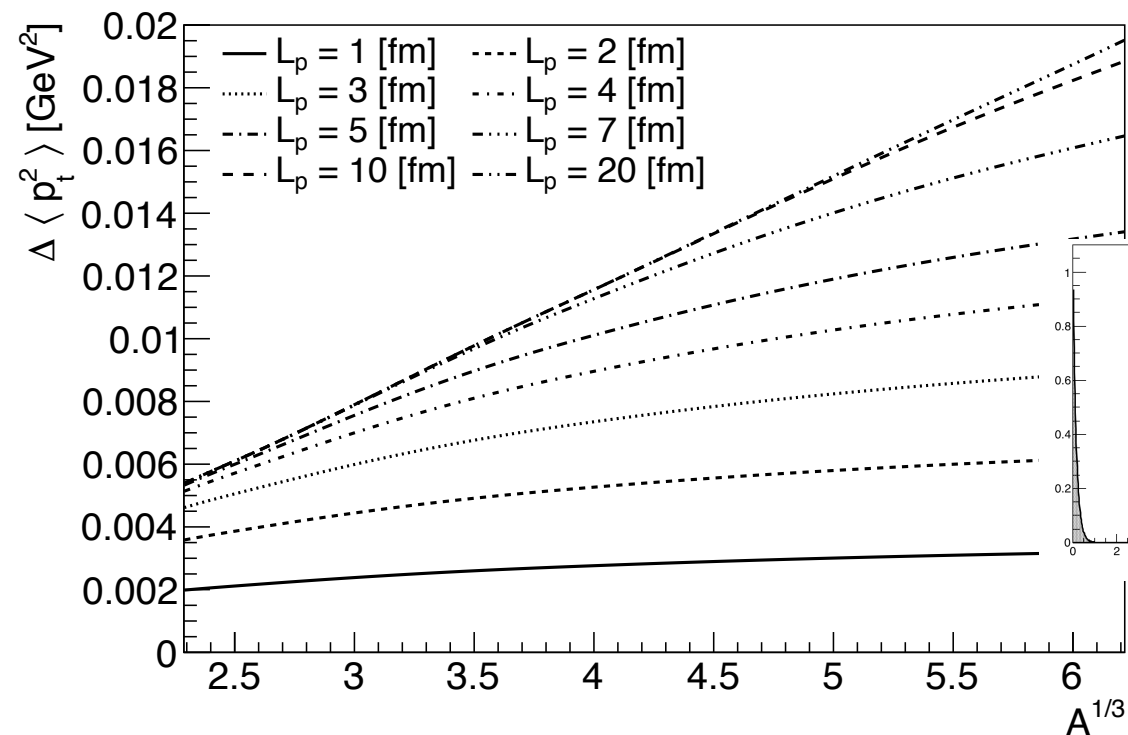




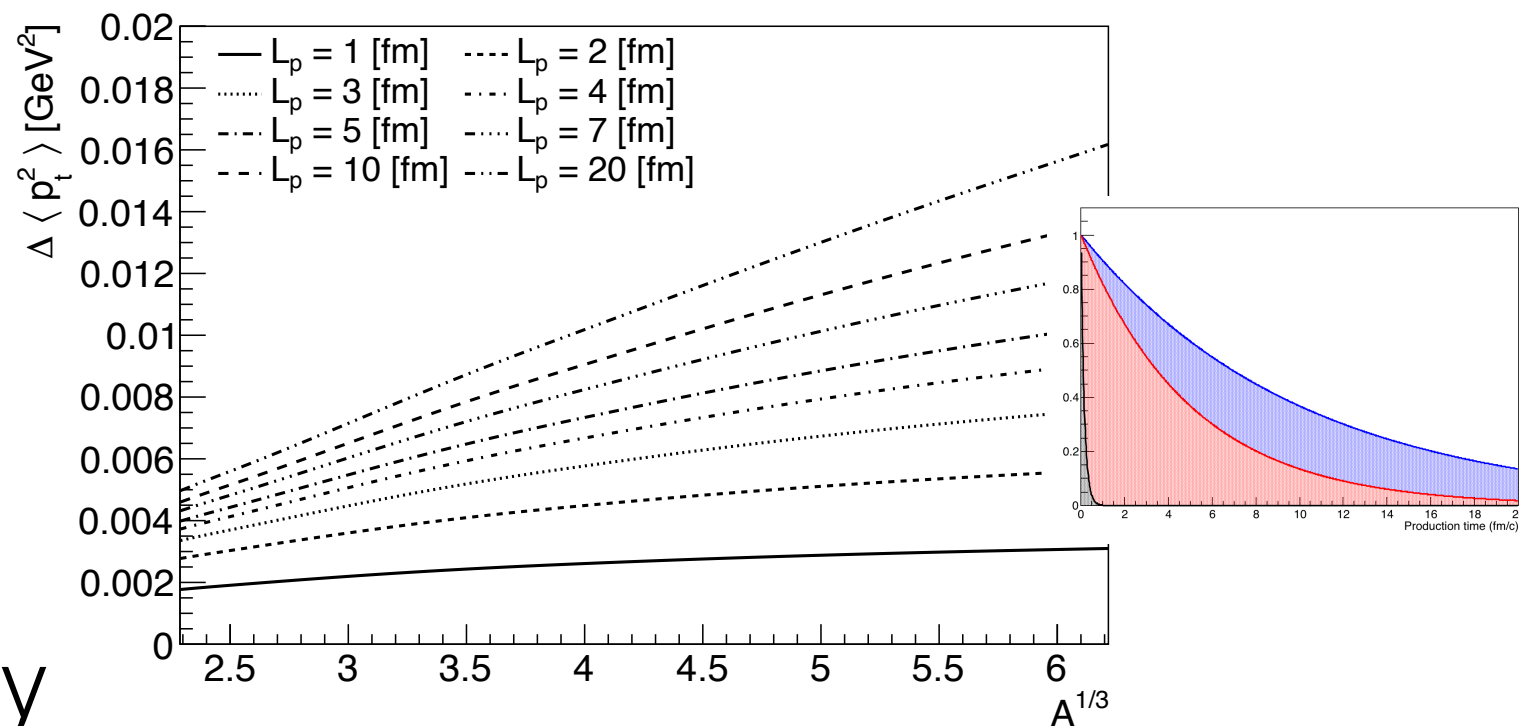
# Effect of production length distribution on $p_T$ broadening



Fixed production time



Exponential production time distribution



QCD factorization  
Relevance at high energy  
Relevance to EIC!



# Tests of exponential distribution hypothesis for quark lifetime

## CLAS Exploratory Study with 5 GeV Data

### Exponential distribution of quark lifetime

103 points, chisquared=69.2, **chisq/dof = 0.685** MEDIUM event selection.  
FCN=69.2253 FROM MINOS STATUS=SUCCESSFUL 10 CALLS 63 TOTAL  
EDM=2.30163e-20 STRATEGY= 1 ERROR MATRIX ACCURATE

EXT NO.	PARAMETER NAME	VALUE	ERROR	STEP SIZE	FIRST DERIVATIVE
1	p0	1.07864e+00	4.83476e-01	-0.00000e+00	6.52690e-07
2	p1	9.33423e-01	2.45714e-01	2.45714e-01	7.34350e-11

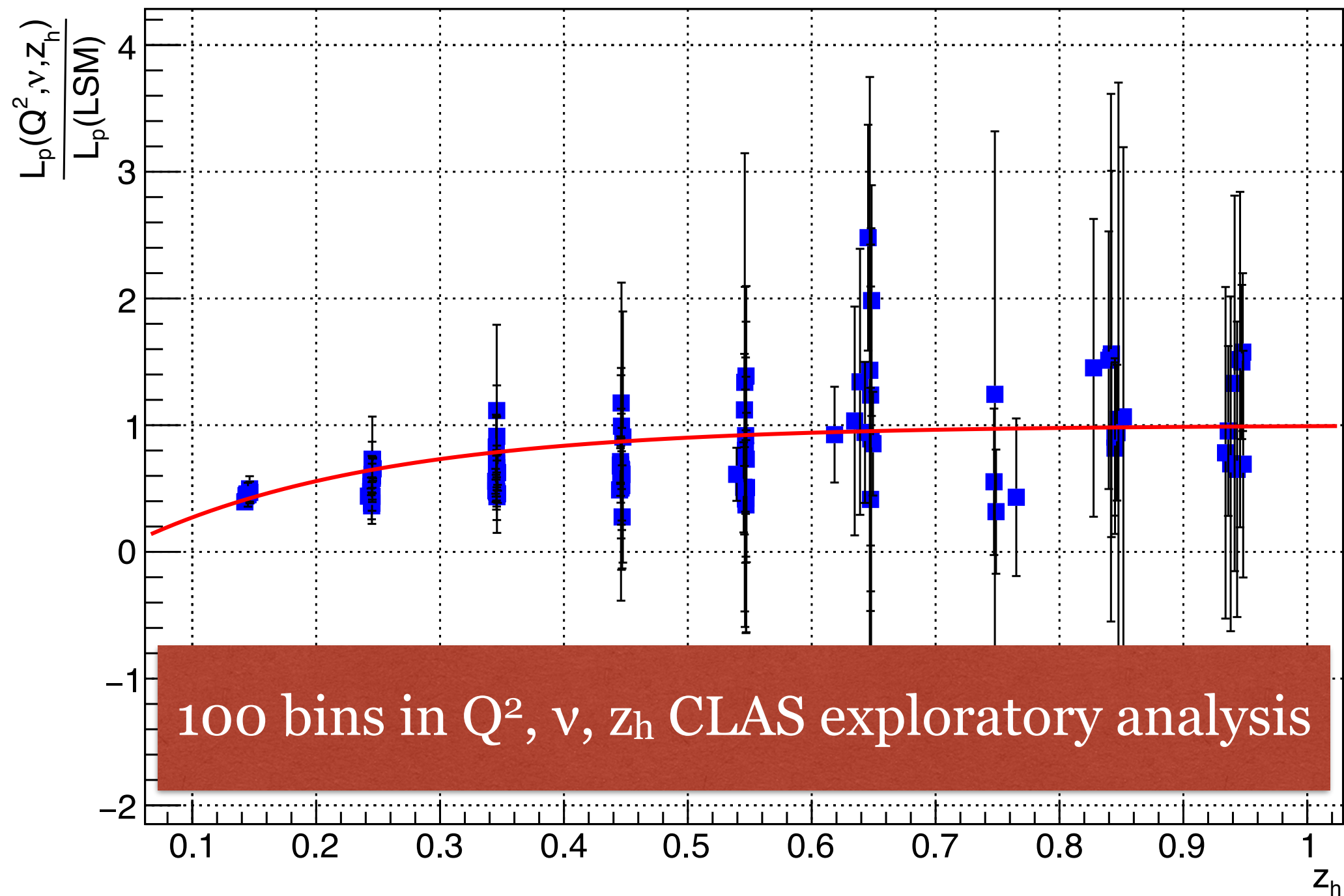
### Single value of quark lifetime

88 points, chisquared=289.5, **chisq/dof = 3.36** MEDIUM event selection.  
FCN=289.533 FROM MINOS STATUS=SUCCESSFUL 8 CALLS 63 TOTAL  
EDM=3.95499e-19 STRATEGY= 1 ERROR MATRIX ACCURATE

EXT NO.	PARAMETER NAME	VALUE	ERROR	STEP SIZE	FIRST DERIVATIVE
1	p0	1.95920e+00	2.75776e-01	-0.00000e+00	8.75252e-07
2	p1	3.95062e-01	1.37012e-01	1.37012e-01	-3.09899e-10

The data clearly prefer an exponential distribution

# CLAS Exploratory Analysis $\approx$ Lund String Model



$L_p(Q^2, \nu, z_h)$  from CLAS analysis similar to values from the Lund String Model for  $z_h > 0.4$



# Space-time characteristics of the struck quark

Assume: Single-photon exchange, no quark-pair production

“JLab” example:  $Q^2 = 3 \text{ GeV}^2$ ,  $\nu = 3 \text{ GeV}$ . ( $x_{\text{Bj}} \sim 0.5$ )

Struck quark absorbs virtual photon energy  $\nu$  and momentum

$$p_{\gamma^*} = |\vec{p}_{\gamma^*}| = \sqrt{(\nu^2 - Q^2)}.$$

- Neglect any initial momentum/mass of quark
- Immediately after the interaction, quark mass  $m_q = Q = \sqrt{Q^2}$ .
- Gamma factor is therefore  $\gamma = \nu/Q$ ,  $\beta = p_{\gamma^*}/\nu$ .

JLab example:  $\gamma = 1.73$ ,  $\beta = 0.82$

Rigorous?  $\gamma$ ,  $\beta$  allow:

1. extrapolations to EIC kinematics,
2. test of time dilation in CLAS fits, and
3. direct comparison between JLab and HERMES fits

# Time dilation test of the CLAS exploratory fit

If what we are measuring is the lifetime of the virtual struck quark as we assert, it must be subject to time dilation

We can test this by plotting **production time vs. relativistic  $\gamma$**

However, there are also dynamical processes affecting production time too (“predicted by Niedermayer in 1980’s”)

Try: fit  $L_p/\gamma$  in 100 bins to determine other dependencies.

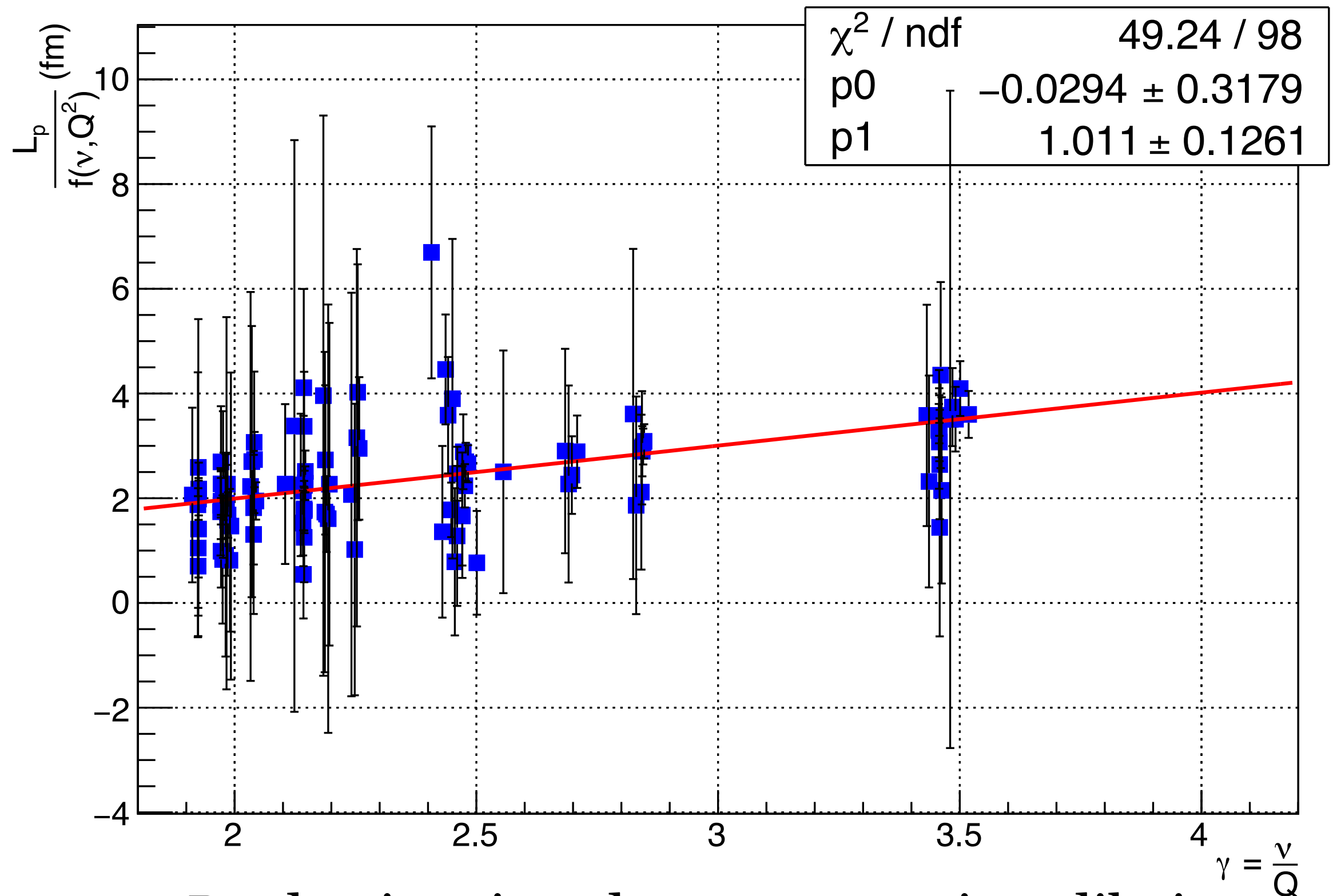
Find:

$$L_p/\gamma = 0.88 + 0.18*Q^2 - 0.16*\nu \equiv f(\nu, Q^2)$$

(Parameter uncertainties 10-20%, chisquared/dof  $\sim 0.5$ )

Next, fit  $L_p/f(\nu, Q^2)$  to see if dependence on  $\gamma$  is consistent with time dilation. Results:

# Time dilation test of the CLAS exploratory fit



Production time demonstrates time dilation

Average slope of  $L_p$  vs  $\gamma$  is  $1 \pm 0.1$ !



# Extrapolation from HERMES to EIC and CLAS

Using the prescription  $\gamma=v/Q$  and  $\beta = p_{\gamma^*}/v$ , we can extrapolate:

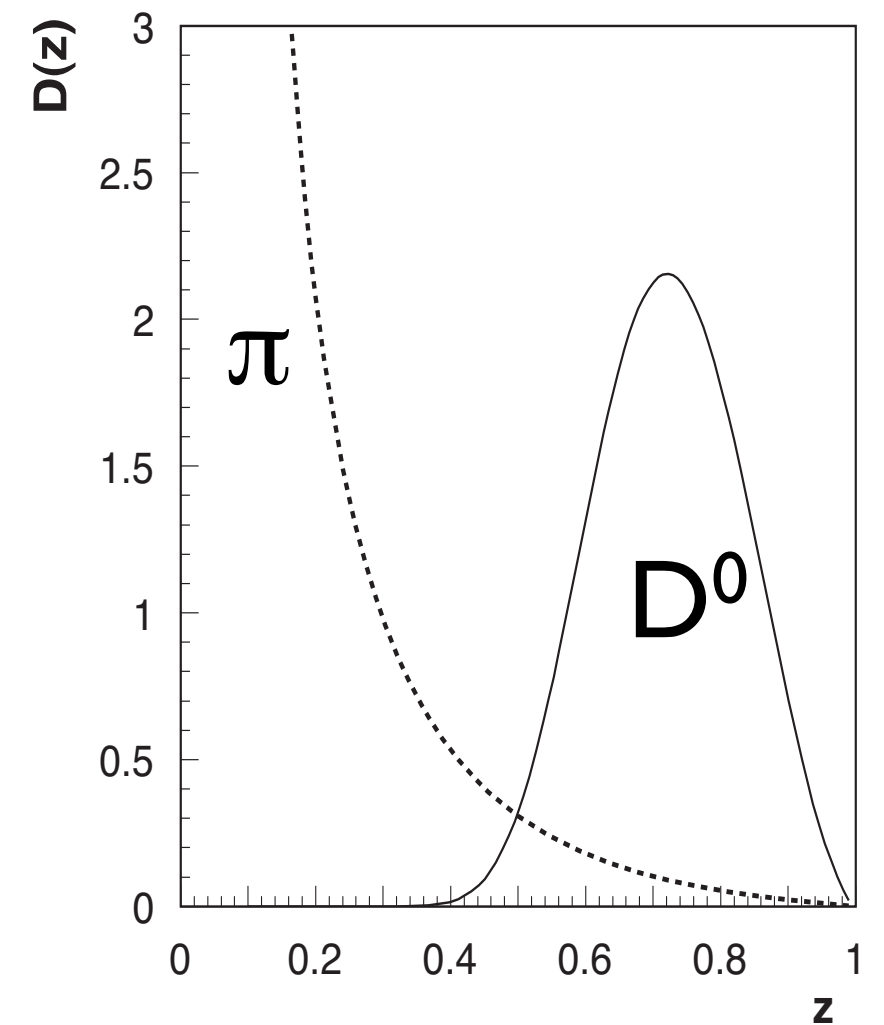
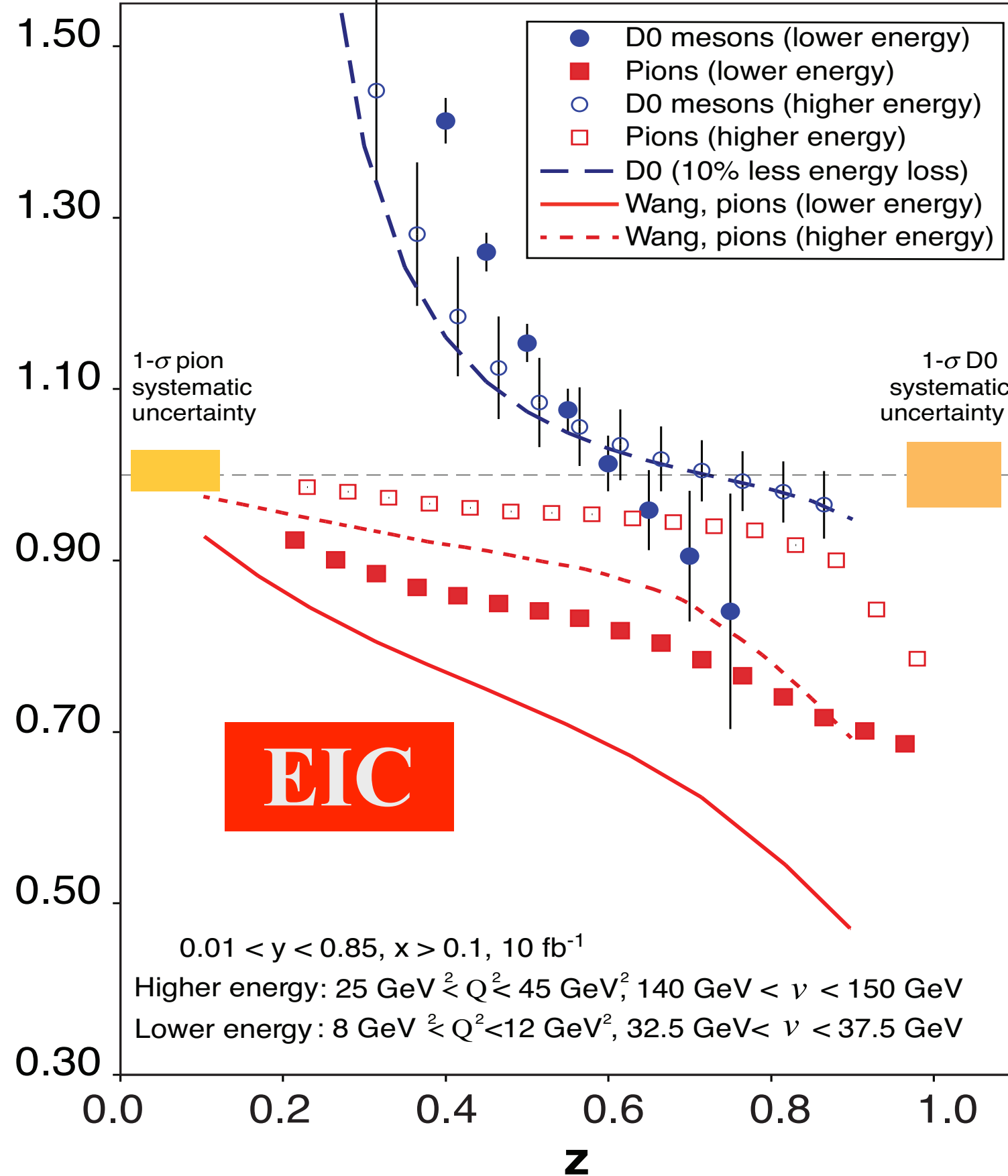
Q <sup>2</sup>	nu	beta*gamma	lp, z=0.32	lp, z=0.53	lp, z=0.75	lp, z=0.94	Experiment	x
2.40	14.50	9.31	8.57				HERMES	0.09
2.40	13.10	8.40		6.39			HERMES	0.10
2.40	12.40	7.94			4.63		HERMES	0.10
2.30	10.80	7.05				2.40	HERMES	0.11
3.00	4.00	2.08	1.92	1.58	1.21	0.71	CLAS	0.40
7.00	7.00	2.45	2.26	1.86	1.43	0.83	CLAS12	0.53
1.00	4.00	3.87	3.57	2.95	2.26	1.32	CLAS	0.13
2.00	9.00	6.28	5.79	4.78	3.66	2.14	CLAS12	0.12
12.00	32.50	9.33	8.59	7.10	5.44	3.18	EIC	0.20
8.00	37.50	13.22	12.17	10.06	7.71	4.50	EIC	0.11
45.00	140.00	20.85	19.20	15.86	12.15	7.10	EIC	0.17
27.00	150.00	28.85	26.57	21.96	16.82	9.82	EIC	0.10

At EIC we can study a wide range of production lengths!

# The Breakthrough Potential of EIC

- Solving the heavy quark puzzle via heavy meson production (see following slides)
- Precision time dilation tests over a wide range in  $\gamma$
- pQCD enhanced non-linear broadening (see following)
- Flavor dependencies of formed hadrons
- $L_p$  distribution determination

Multiplicity Ratio



**EIC**  
Year 1

Definitive comparisons of  
light quark and heavy quark  
energy loss

Access to very strong, unique light quark energy loss signature via  $D^0$  heavy meson. Compare to s and c quark energy loss in  $D_s^+$



# NEW THEORY DEVELOPMENT

- T. Liou, A.H. Mueller, B. Wu: Nuclear Physics A 916 (2013) 102–125, arXiv:1304.7677
- Old: multiple scattering  $\rightarrow$  gluon emission, = energy loss

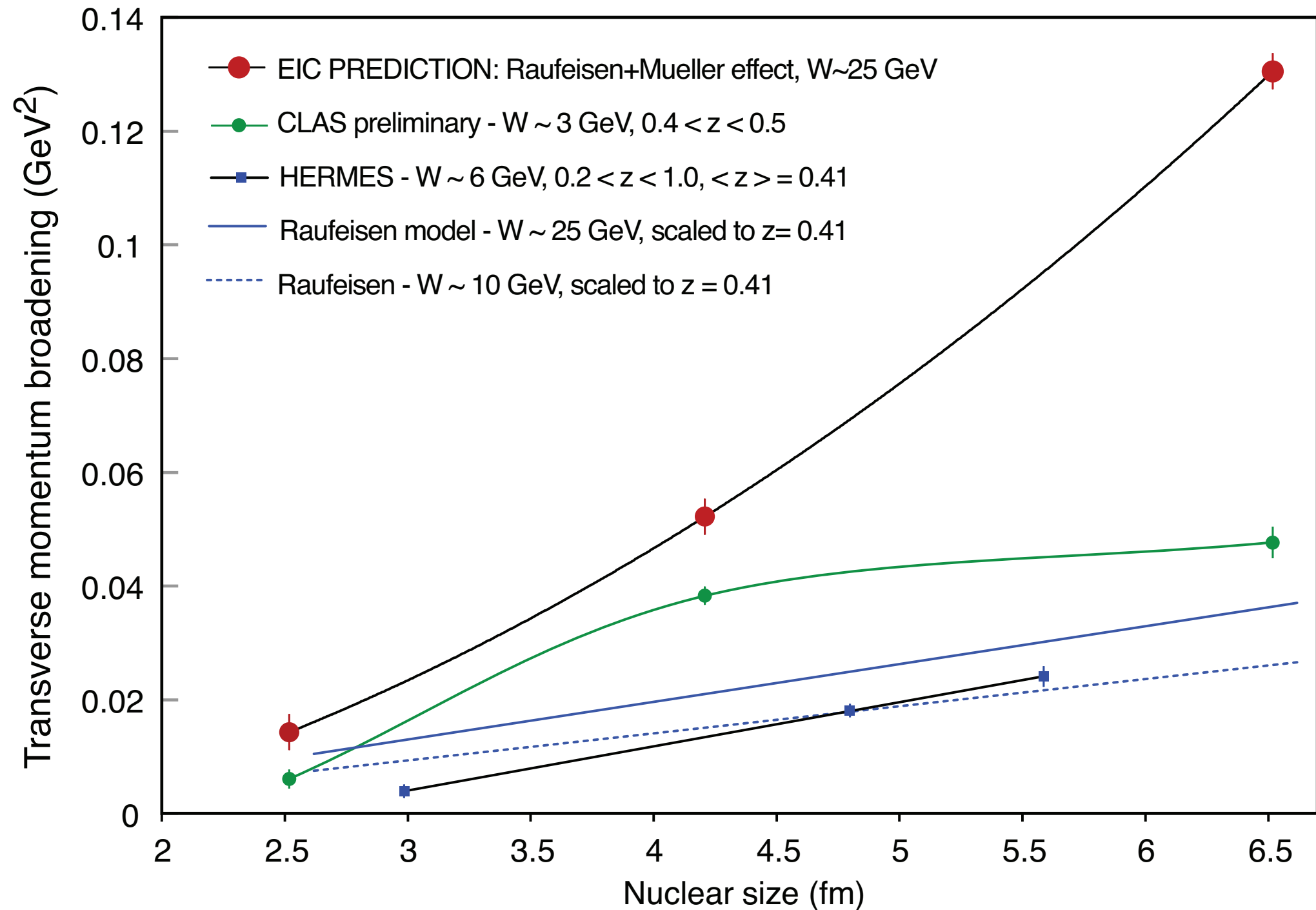
$$-\frac{dE}{dx} = \frac{\alpha_s N_c}{4} \Delta p_T^2 \propto \hat{q} L$$

- New: this energy loss creates *more*  $p_T$  broadening

$$\Delta p_T^2 = \frac{\alpha_s N_c}{8\pi} \hat{q} L \left[ \ln^2 \frac{L^2}{l_0^2} + \dots \right]$$

$\rightarrow$  predicts a non-linear relationship between  $p_T$  broadening and  $L$ .  
we can look for this at EIC!

# pQCD description of quark energy loss on $p_T$ broadening



DIS channels: *stable* hadrons, accessible with 11 GeV

JLab future experiment PR12-06-117

Actively underway with existing 5 GeV data

HERMES

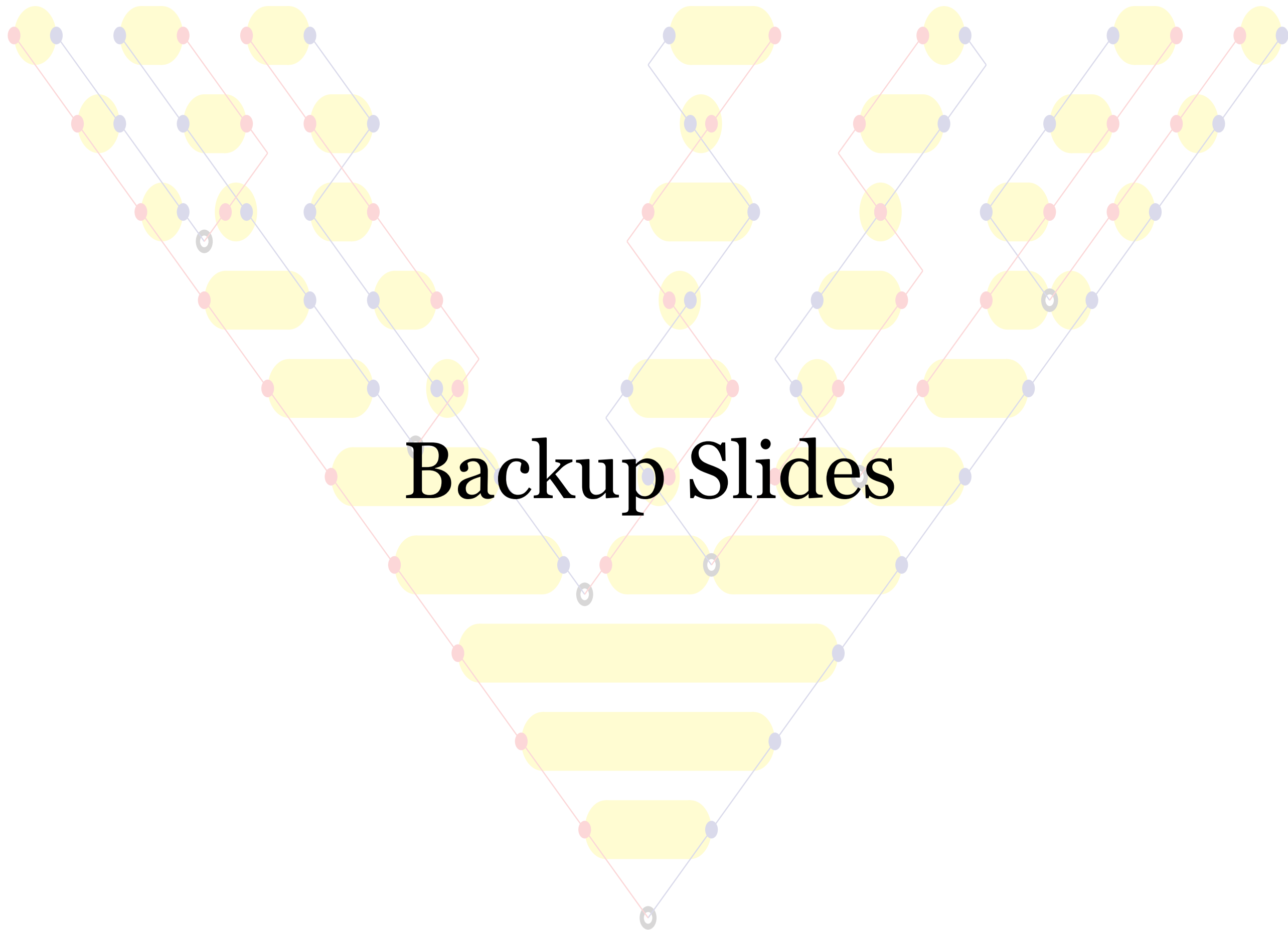
EIC: heavy  
mesons and  
baryons; wide  
kinematic  
range!

meson	c $\tau$	mass	flavor content	baryon	c $\tau$	mass	flavor content
$\pi^0$	25 nm	0.13	ud	$p$	stable	0.94	ud
$\pi^+, \pi^-$	7.8 m	0.14	ud	$\bar{p}$	stable	0.94	ud
$\eta$	170 pm	0.55	uds	$\Lambda$	79 mm	1.1	uds
$\omega$	23 fm	0.78	uds	$\Lambda(1520)$	13 fm	1.5	uds
$\eta'$	0.98 pm	0.96	uds	$\Sigma^+$	24 mm	1.2	us
$\phi$	44 fm	1	uds	$\Sigma^-$	44 mm	1.2	ds
f1	8 fm	1.3	uds	$\Sigma^0$	22 pm	1.2	uds
$K^0$	27 mm	0.5	ds	$\Xi^0$	87 mm	1.3	us
$K^+, K^-$	3.7 m	0.49	us	$\Xi^-$	49 mm	1.3	ds

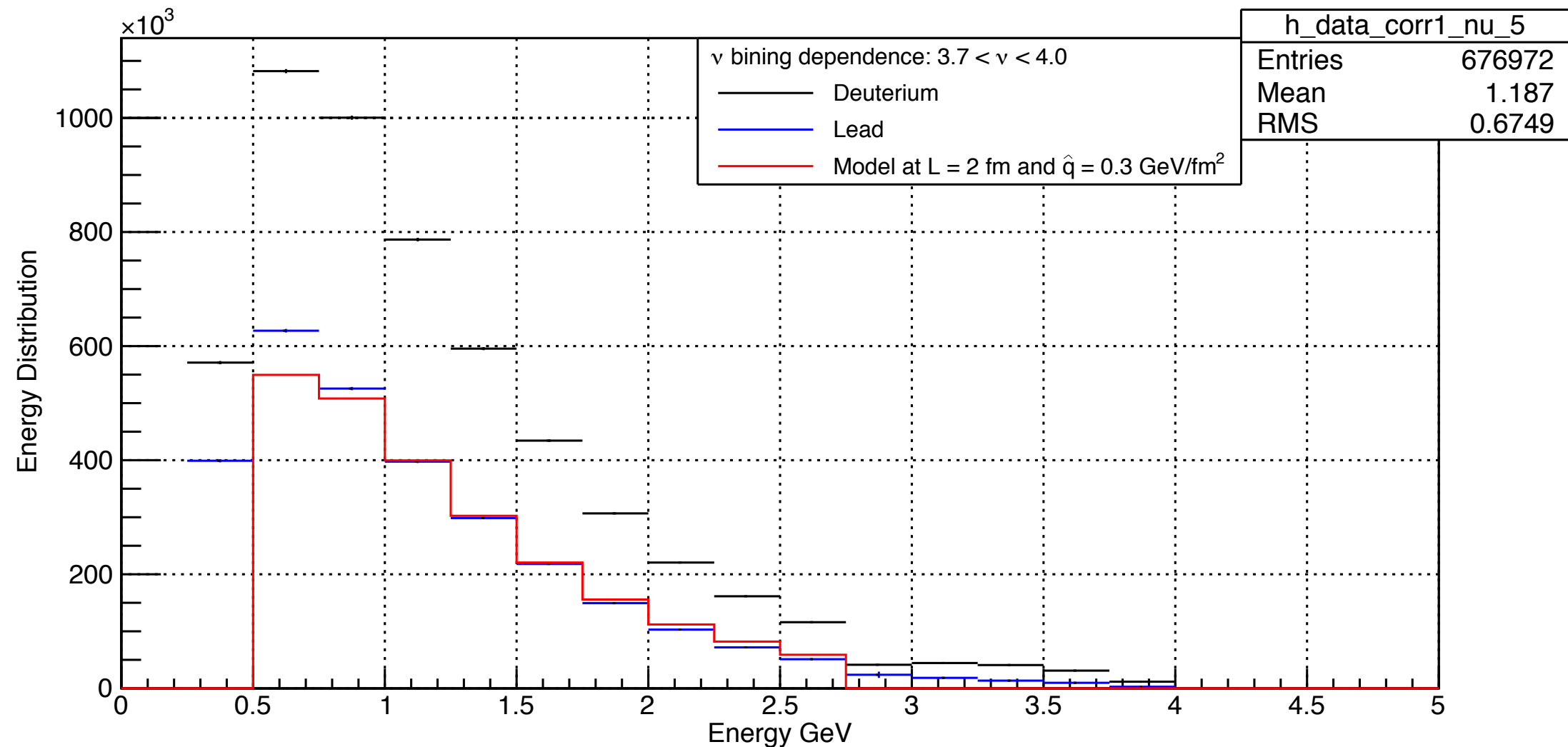


# Conclusions

- Completely new categories of physics analyses. A broad and deep program of studies for the future!
- First direct measurement of quark energy loss
- Extracting characteristic times: semi-inclusive DIS
  - ➔ HERMES data - we measure the production time, and independently obtain the Lund string constant of  $1 \text{ GeV}/\text{fm}$
  - ➔ CLAS - (exploratory) observation of time dilation, sensitivity to production length distribution form, comparison to HERMES results through Lorentz boost
  - ➔ Clear connections to confinement, QCD factorization, Electron Ion Collider, higher energies
- Much much more in future: **12 GeV** and **EIC**:
  - ➔ Heavy quark puzzle; time dilation; pQCD enhanced broadening; flavor dependences;  $L_p$  distribution studies



# Comparison to Arleo model



Black: Deuterium, Blue: Lead data  
Red: Energy loss model evolution of deuterium to lead,  
production length 2 fm



# Geometric model description I

- Propagating quark causes  $p_T$  broadening of hadron
- Propagating (pre-)hadron “disappears” when it undergoes an inelastic interaction with cross section  $\sigma$
- Implemented as a Monte Carlo calculation in  $x, y, z, L_p$
- Simultaneous fit of  $p_T$  broadening and multiplicity ratio
- Realistic nuclear density, integrated along path w/ GSL

Path of quark is  
divided into  
“partonic phase” and  
“hadronic phase”





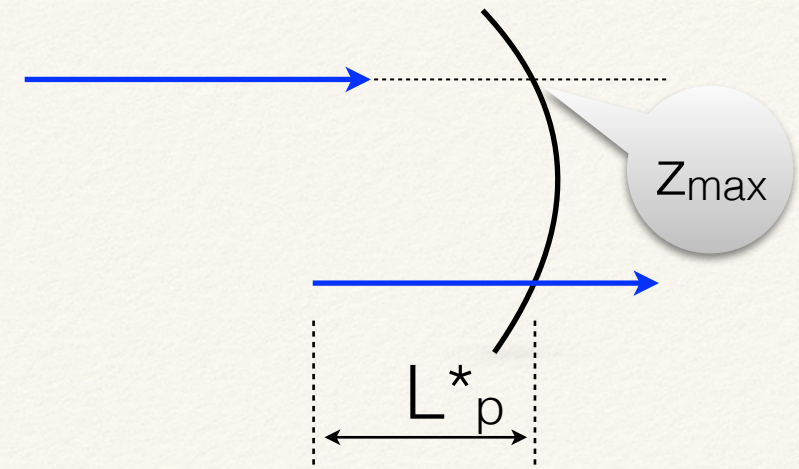
# Geometric model description II

Model implemented with 3, 4 or 5 **parameters**:

1. **q-hat** parameter (transport coefficient) that sets the scale of  $p_T$  broadening
2. Production length  **$L_p$** : distance over which  $p_T$  broadening and energy loss occur. Assumed exponential form.
3. **Cross section** for prehadron to interact with nucleus.
4. **Shift in z** caused by quark energy loss in medium
5. Average **distance between scatterings** or “mean free path”  $l_0$   
(alternative form of  $p_T$  broadening, proportional to  $L_p \cdot \log^2(L_p/l_0)$ )



# Geometric model description III



$$\langle \Delta p_T^2 \rangle = \langle \hat{q}_0 \int_{z=z_0}^{z=z_0+L_p^*} \rho(x_0, y_0, z) dz \rangle_{x_0, y_0, z_0, L_p}$$

$L_p$  is distributed as exponential

$x_0, y_0, z_0$  thrown uniformly in sphere, weighted by  $\rho(x, y, z)$

$L_p^* = L_p$  except where truncated by integration sphere

$$\langle R_M \rangle = \langle \exp(-\sigma \int_{z=z_0+L_p}^{z=z_{max}} \rho(x, y, z) dx dy dz) \rangle_{x_0, y_0, z_0, L_p}$$

The above are computed sequentially (same  $x_0, y_0, z_0, L_p$ )

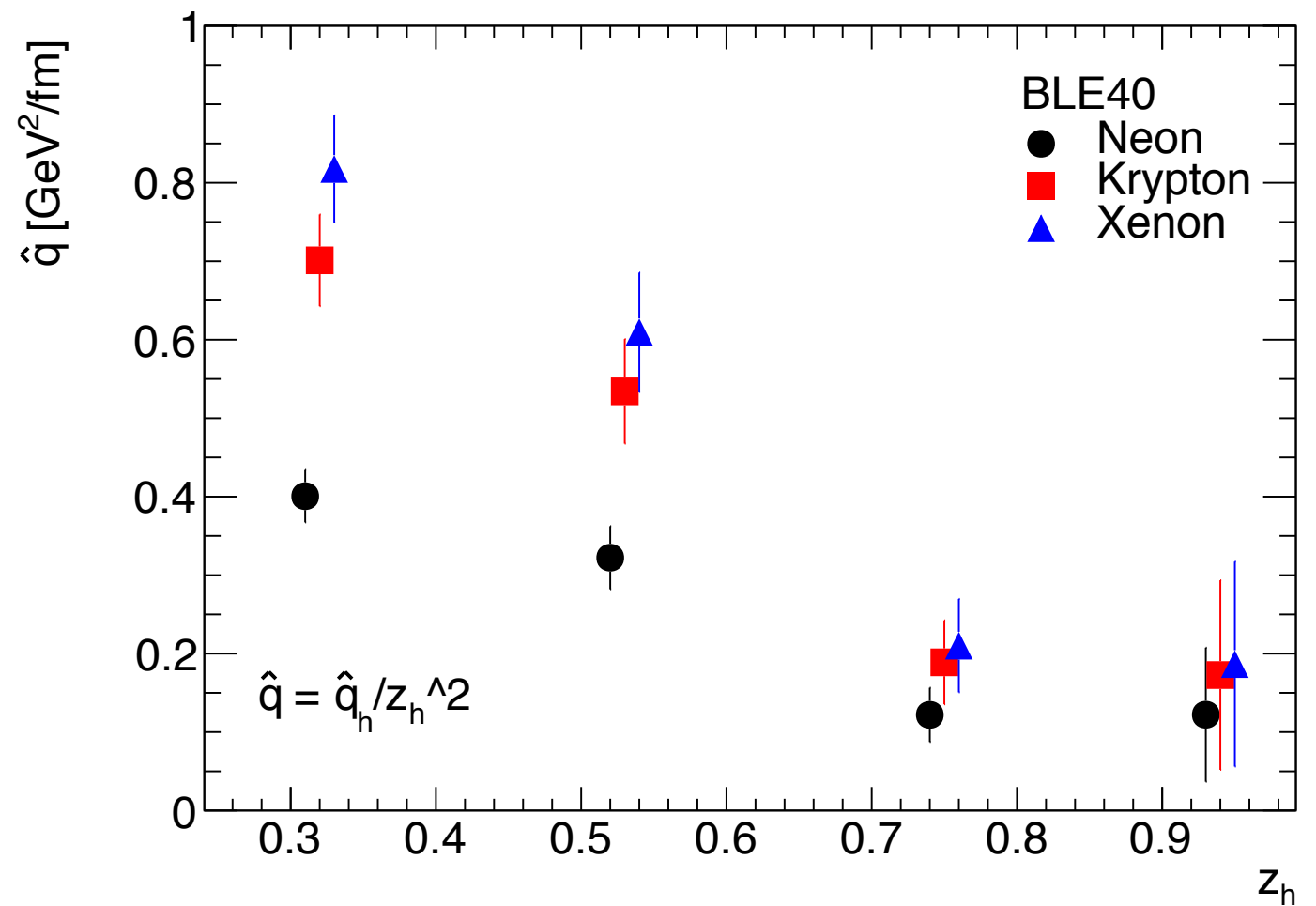
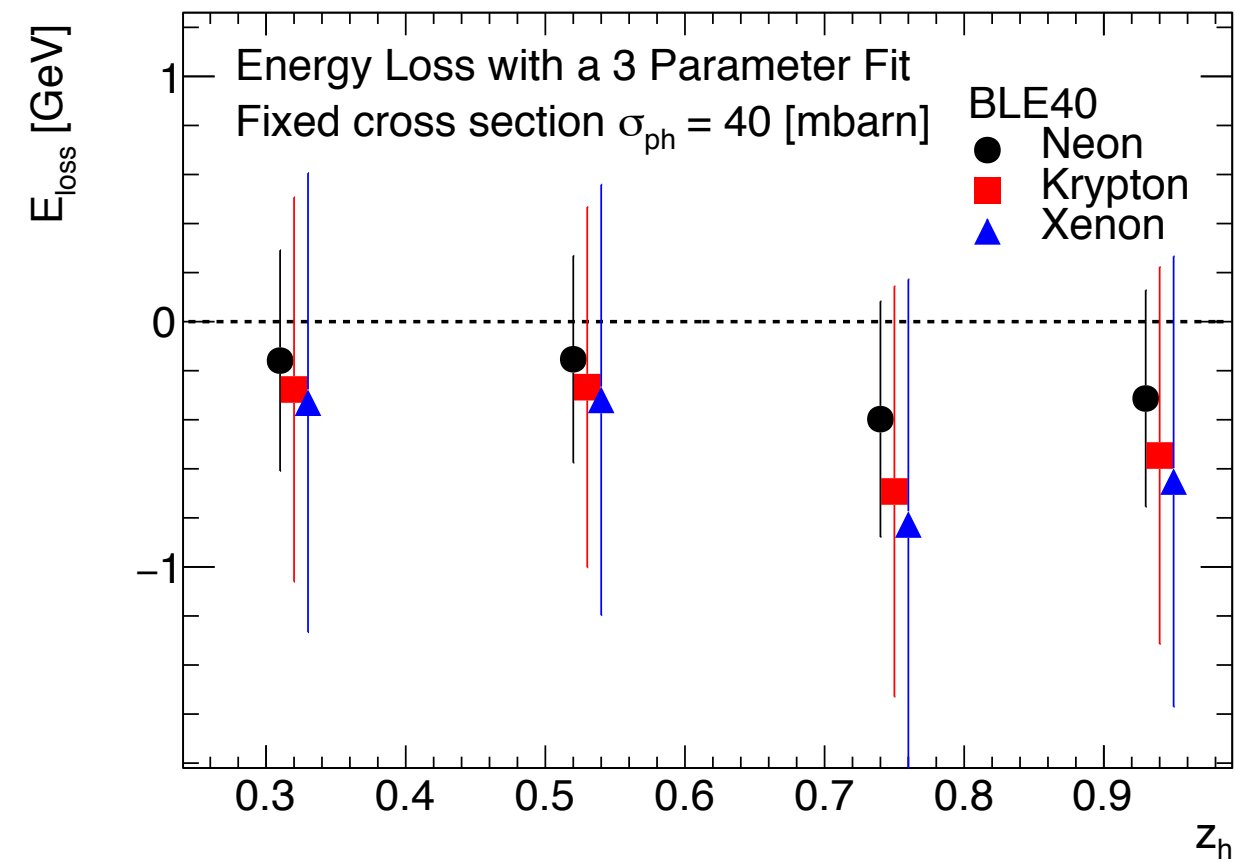
Data in  $(x, Q^2, z)$  bin: fitted to model, 3 parameters:  $\hat{q}_0, \langle L_p \rangle, \sigma$

No dynamical information is assumed; it emerges from fit

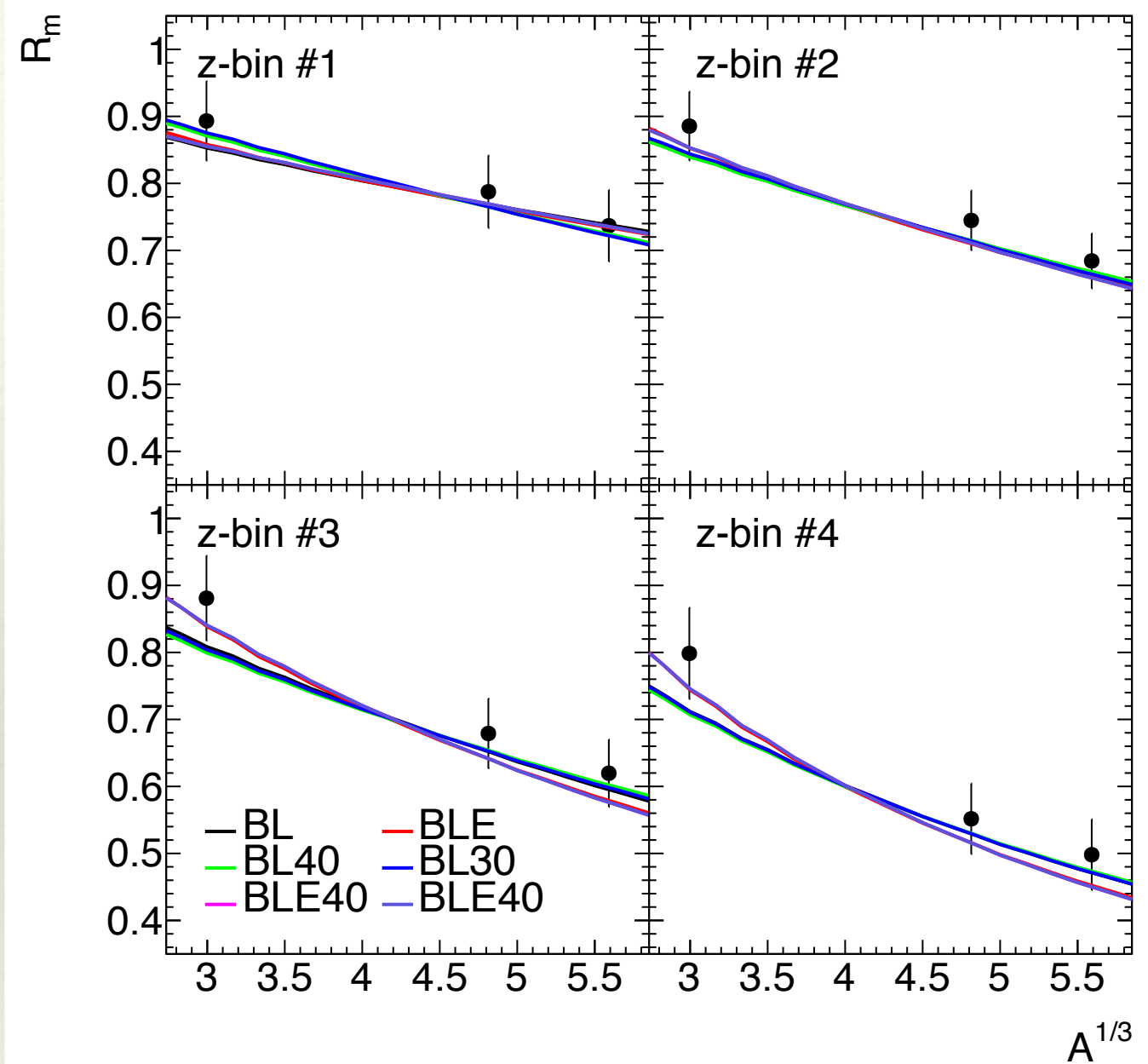
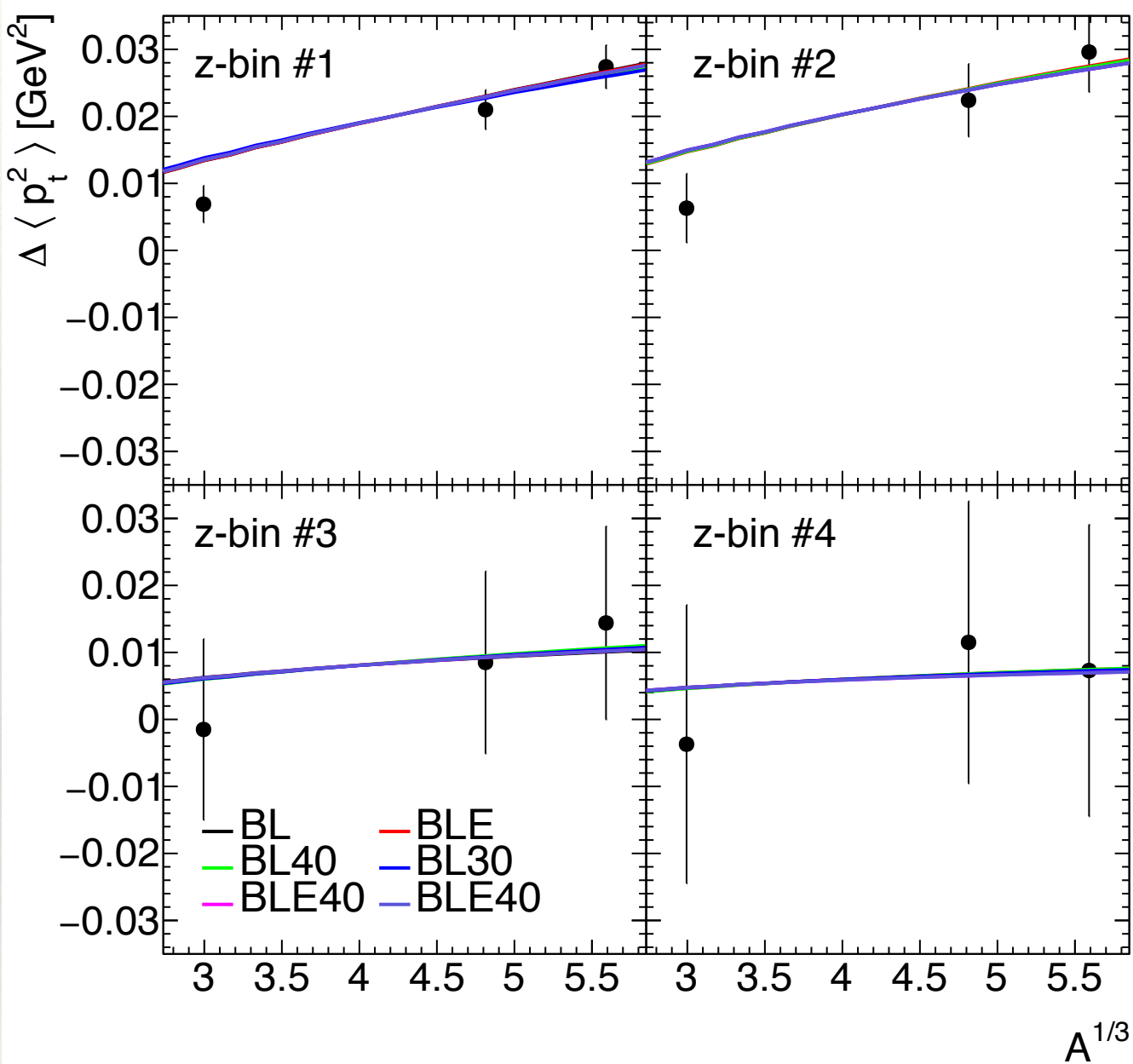
Systematic errors: 3% for multiplicity ratio, 4% for  $p_T$  broadening



# Energy loss and $\hat{q}$ from our HERMES analysis



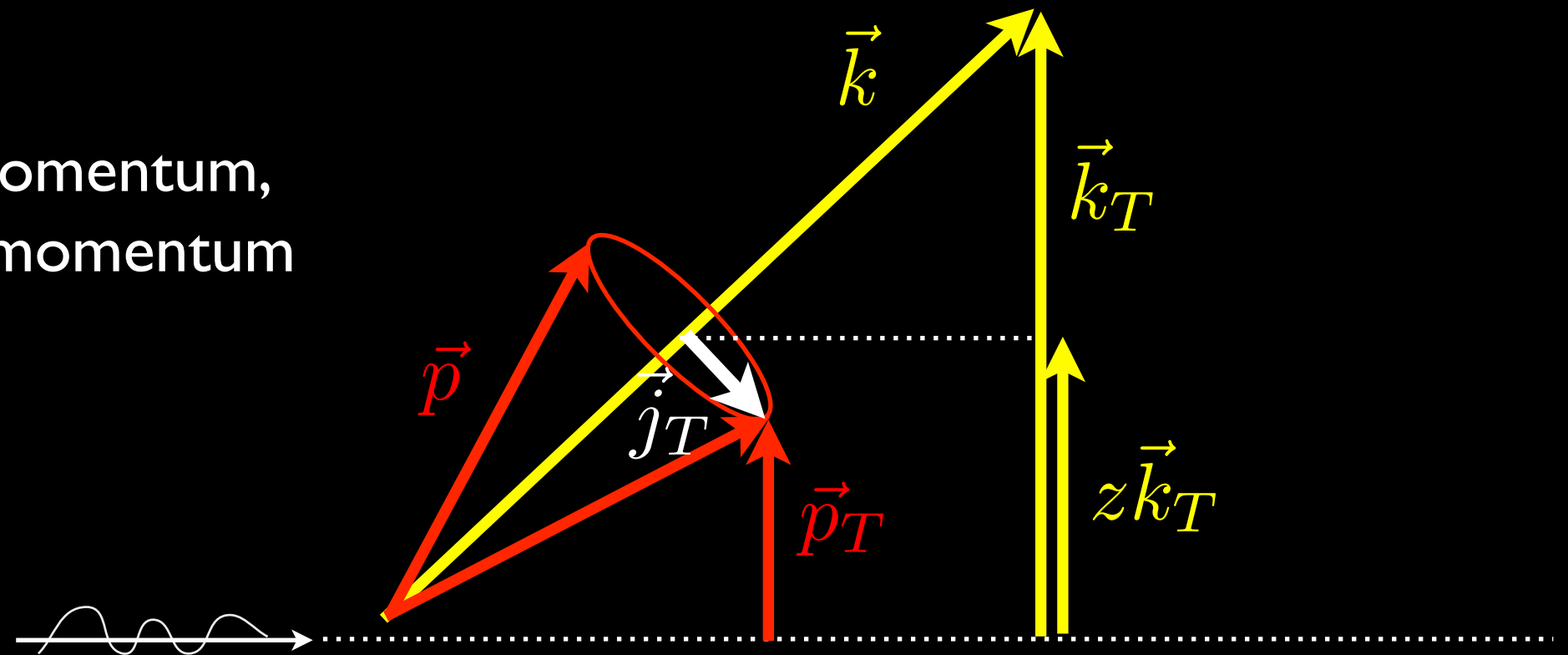
# Fits of data from our HERMES analysis



# Quark $k_T$ broadening vs. *hadron* $p_T$ broadening

The  $k_T$  broadening experienced by a quark is “*diluted*” in the fragmentation process

$k$  is the *quark* momentum,  
 $p$  is the *hadron* momentum



$$\vec{p}_T = z\vec{k}_T + \vec{j}_T$$

$$\langle p_T^2 \rangle = \langle z^2 k_T^2 \rangle + \langle j_T^2 \rangle$$

$$\Delta \langle p_T^2 \rangle = \Delta \langle z^2 k_T^2 \rangle + \Delta \langle j_T^2 \rangle \sim 0$$

$$\Delta \langle p_T^2 \rangle \approx z^2 \Delta \langle k_T^2 \rangle$$

Verified for pions to 5-10% accuracy for vacuum case,  $z=0.4-0.7$ , by Monte Carlo studies



Basic questions at low energies:

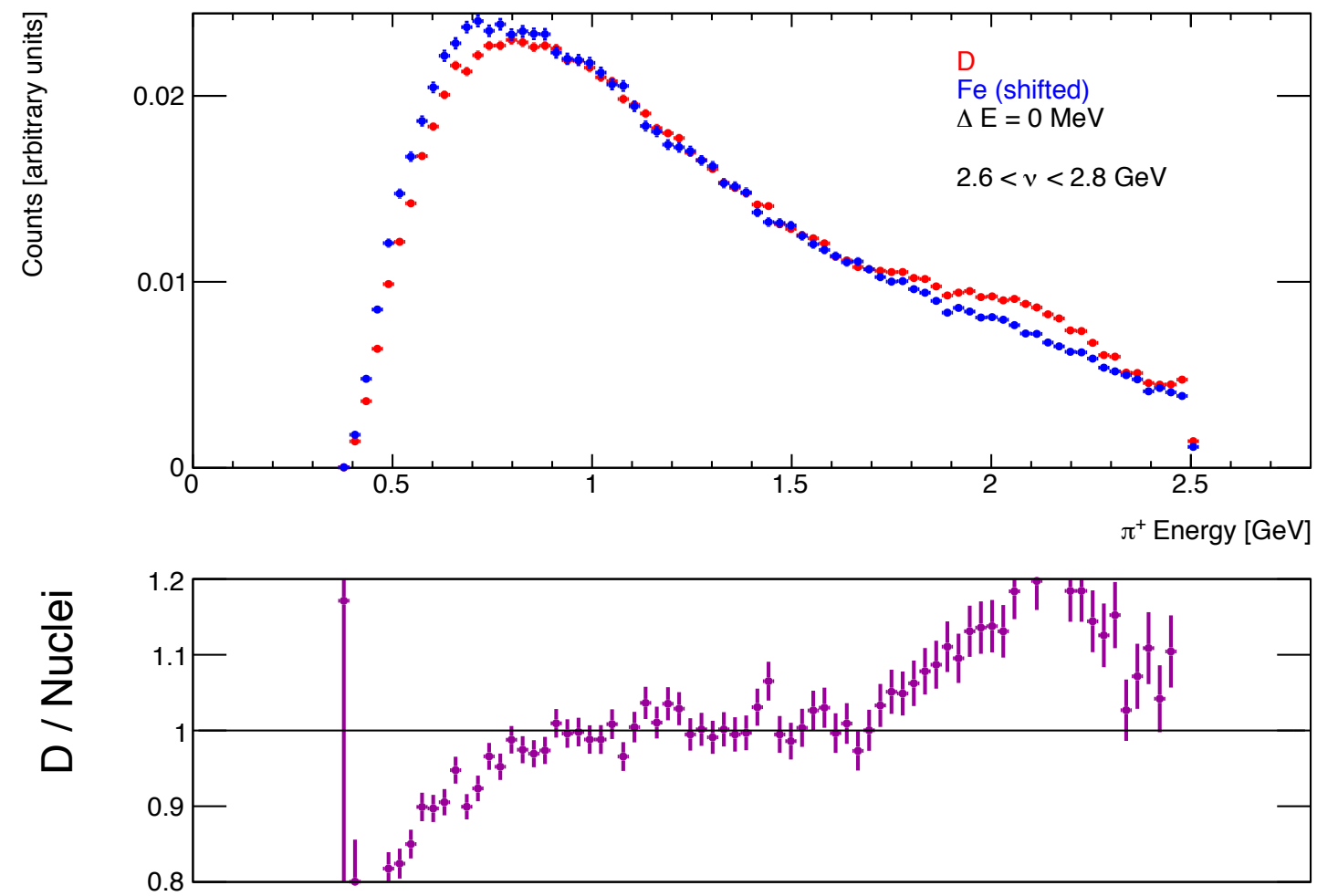
Partonic processes dominate, or hadronic? in which kinematic regime? classical or quantum?

Can identify dominant hadronization mechanisms, uniquely? what are the roles of flavor and mass?

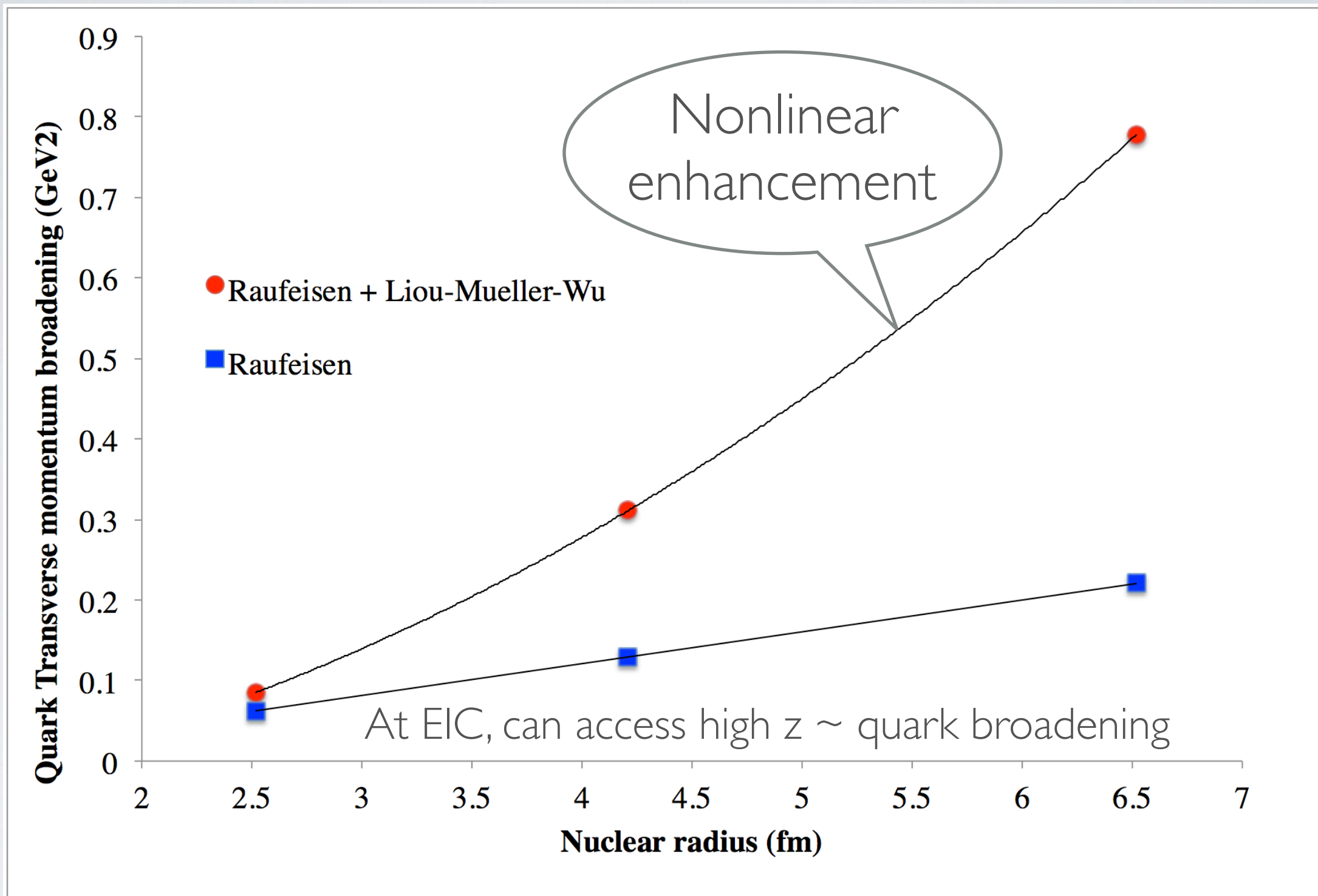
What can we infer about fundamental QCD processes by observing the interaction with the nucleus?

*If  $p_T$  broadening uniquely signals the partonic stage, can use this as one tool to answer these questions*

# What happens if $v$ is too low



# QUARK $K_T$ BROADENING

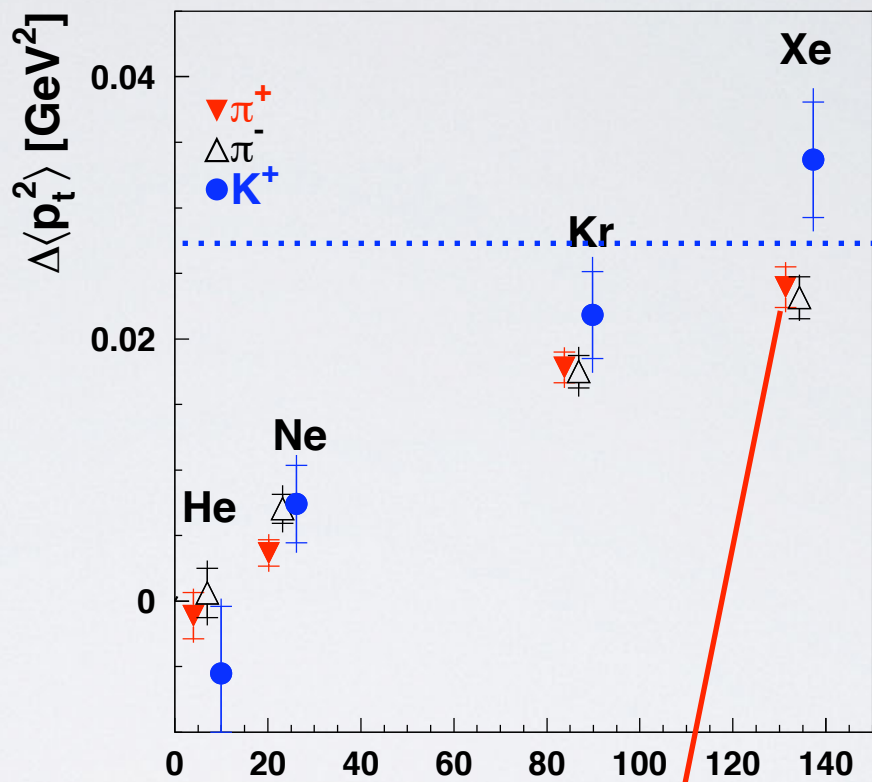
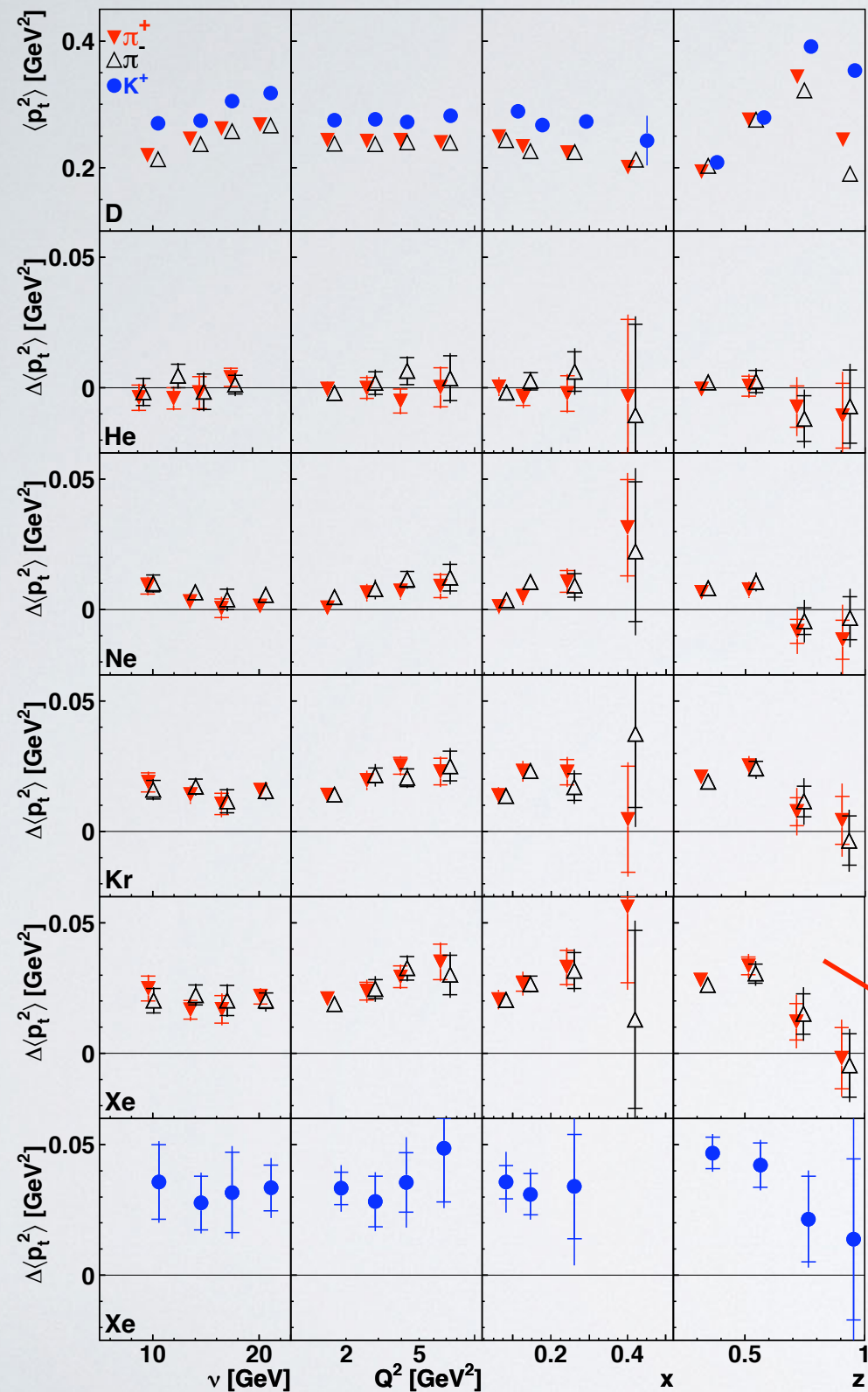


Jörg Raufeisen (Physics Letters B 557 (2003) 184–191) =

Dolejsi, Hüfner, Kopeliovich, Johnson, Tarasov, Baier, Dokshitzer, Mueller, Peigne, Schiff, Zakharov, Guo<sup>2</sup>, Luo, Qiu, Sterman, Majumder, Wang<sup>2</sup>, Zhang, Kang, Zing, Song, Gao, Liang, Bodwin, Brodsky, Lepage, Michael, Wilk...color dipole, BDMPS-Z, higher-twist, etc.

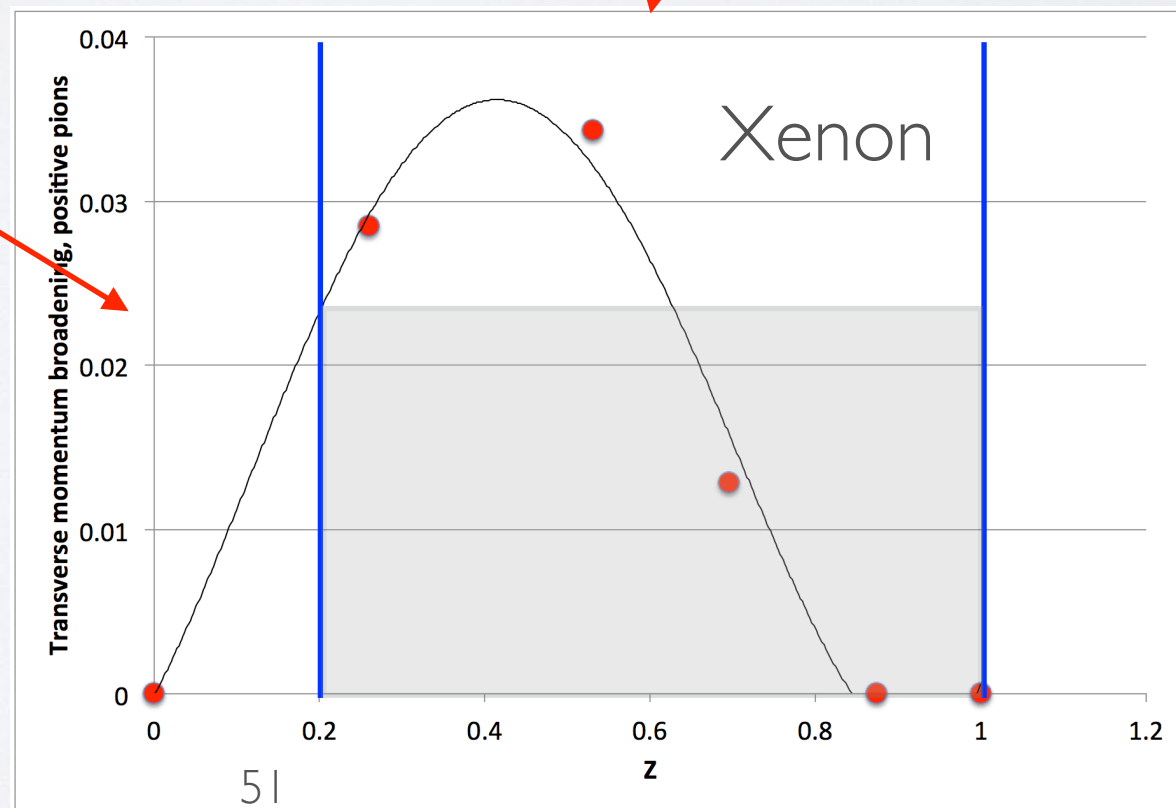


# $P_T$ BROADENING FROM HERMES



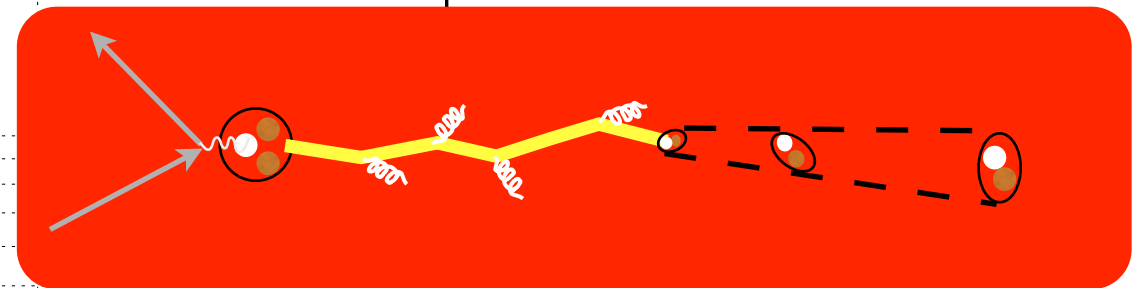
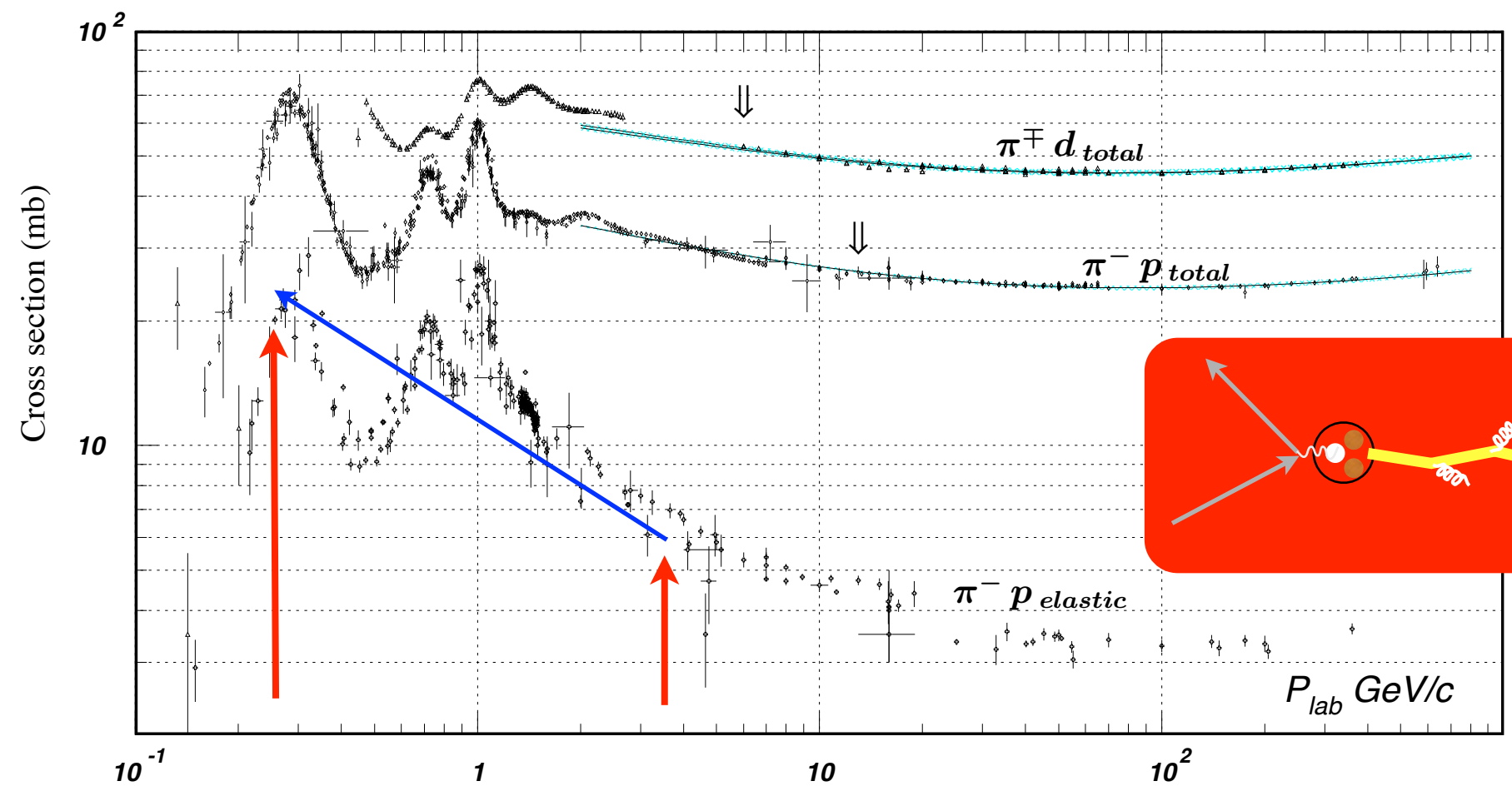
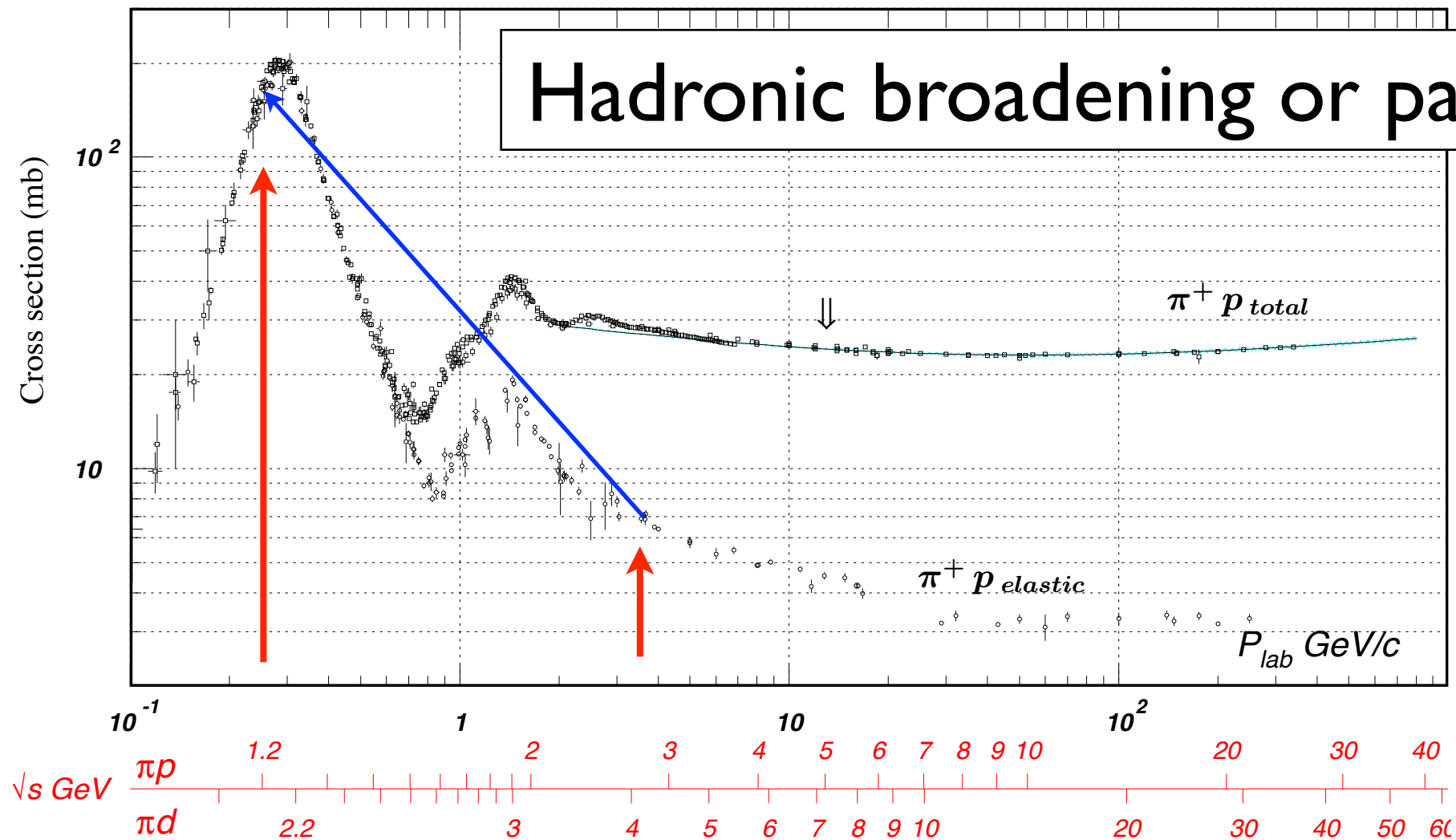
Q: Why is CLAS broadening larger than HERMES broadening?

A: Averaging over  $z$  results in reduced broadening. CLAS results are binned in  $\nu$ ,  $Q^2$ , and  $z$ . (Fermi motion seems to be small, suppressed by a factor of  $\{z \cdot x_{Bj}\}^2$ )



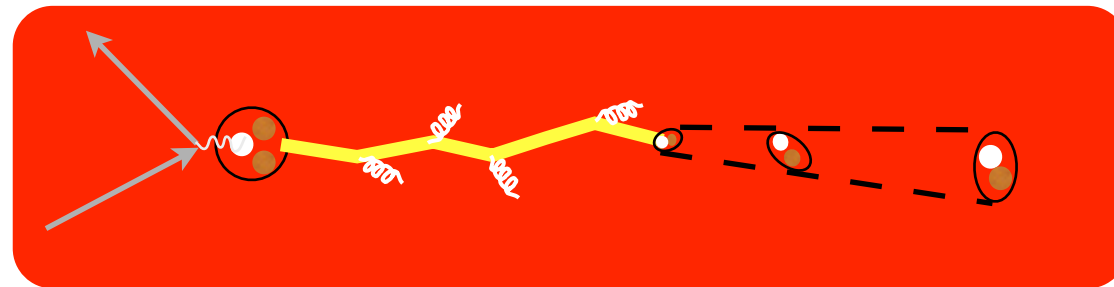
1-D distributions, integrated over all other variables

# Hadronic broadening or partonic broadening?

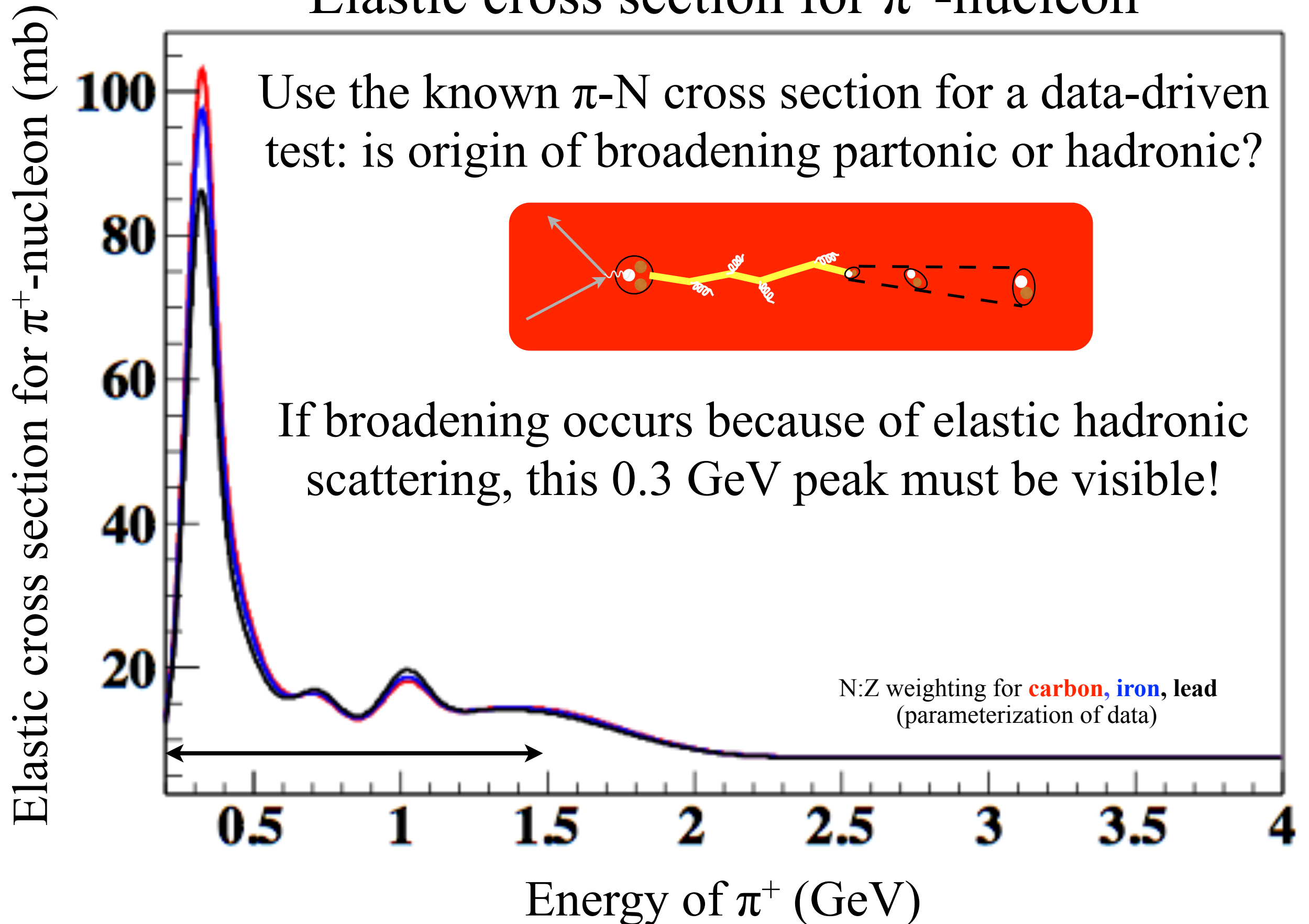


# Elastic cross section for $\pi^+$ -nucleon

Use the known  $\pi$ -N cross section for a data-driven test: is origin of broadening partonic or hadronic?



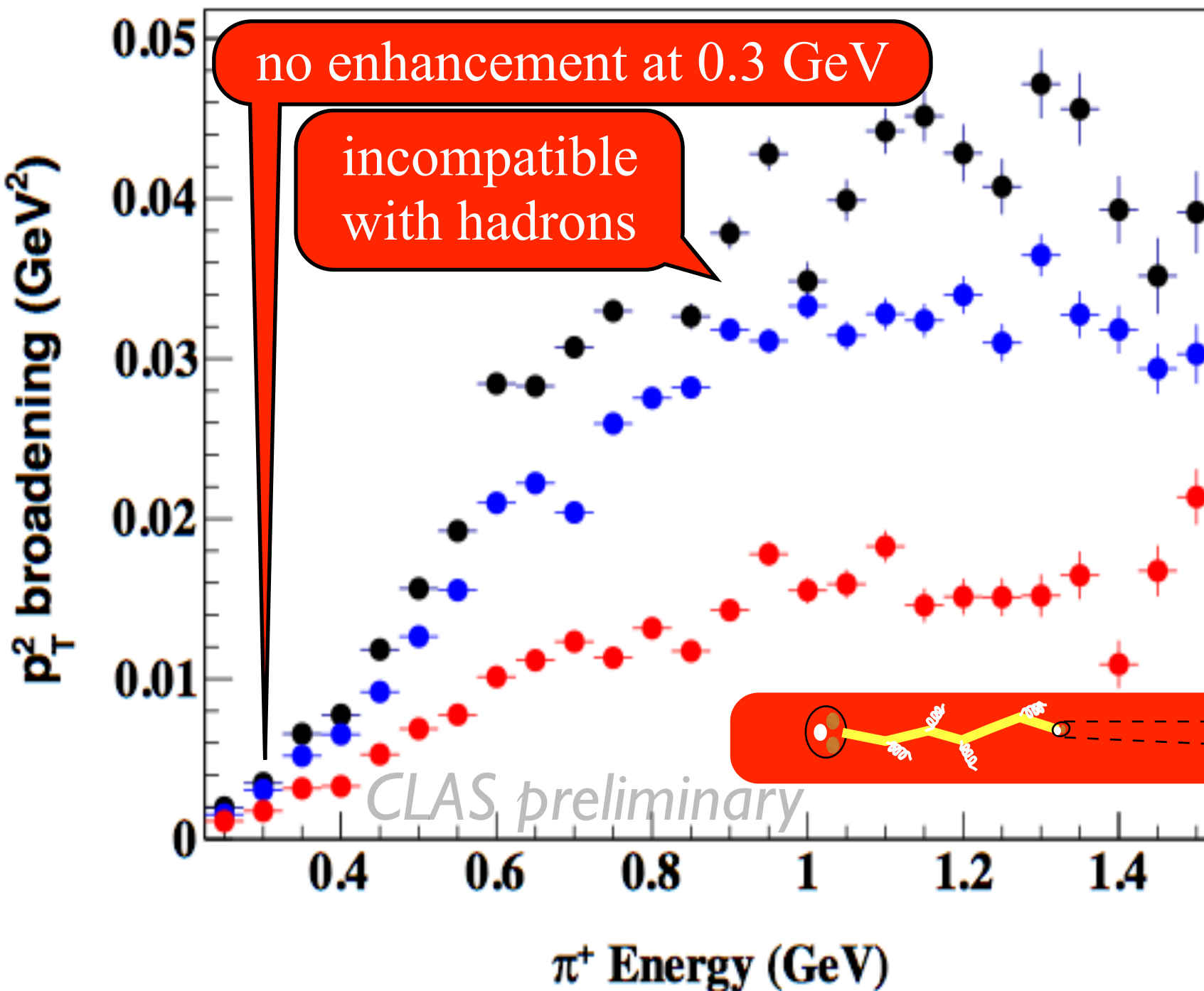
If broadening occurs because of elastic hadronic scattering, this 0.3 GeV peak must be visible!



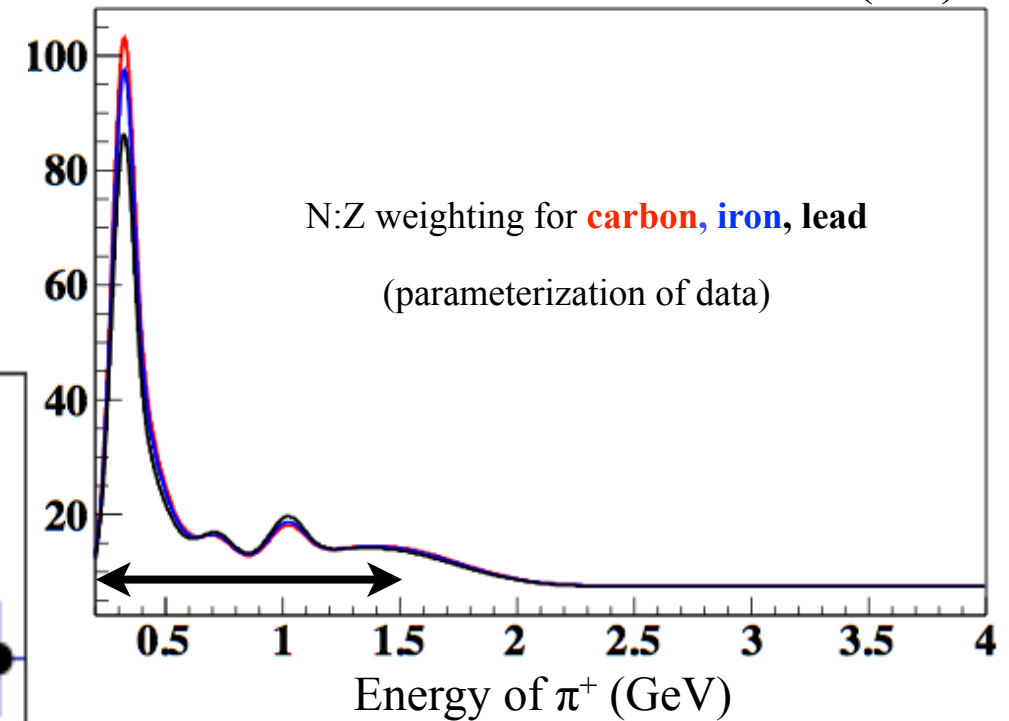


# $p_T^2$ Broadening vs. Hadron Energy

$2.0 < Q^2 < 3.0 \text{ GeV}^2$   $3.4 < \nu < 4.0 \text{ GeV}$



Elastic cross section for  $\pi^+$ -nucleon (mb)



No visible evidence of hadronic elastic scattering?  
Suggests:

1) formation length is very long

2) broadening is purely partonic