

*New description of anharmonic
double- γ excitations in
deformed nuclei using the IBM.*

José Enrique García Ramos^{1,2}, C.E. Alonso², J.M. Arias², and
P. Van Isacker³

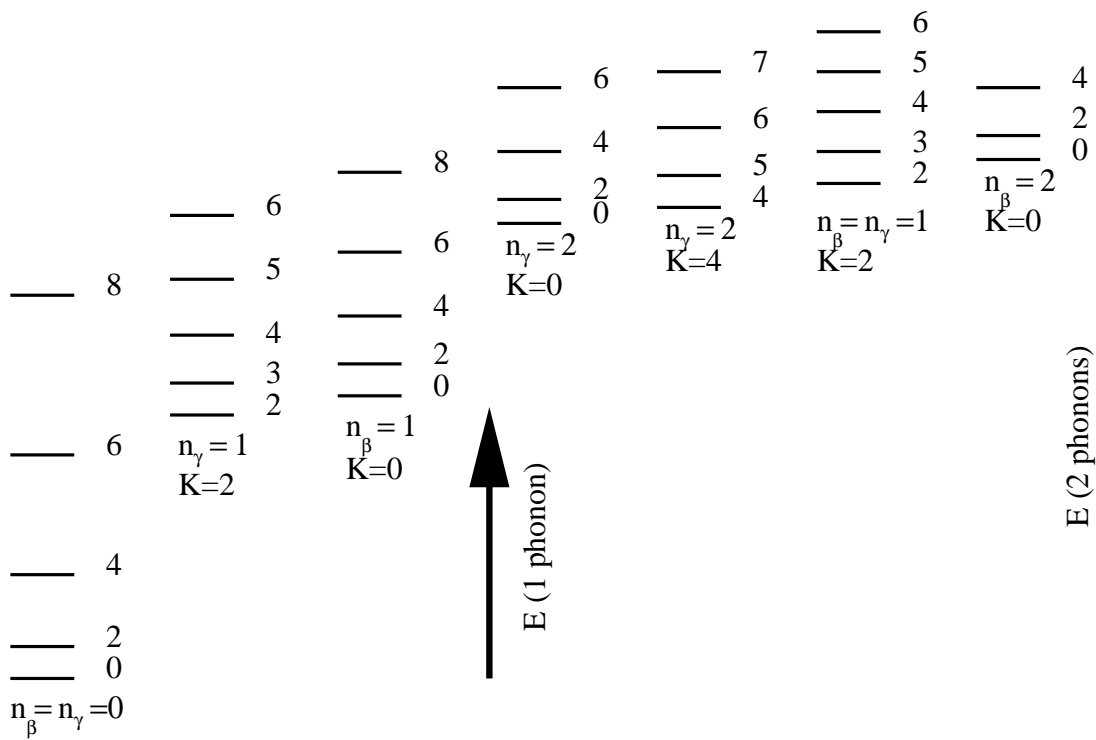
¹Institute for Theoretical Physics, Vakgroep Subatomaire en
Stralingsfysica, Proeftuinstraat 86, B-9000 **Gent, Belgium**

²Departamento de Física Atómica, Molecular y Nuclear, Universidad de
Sevilla, Apto. 1065, 41080 **Sevilla, Spain**

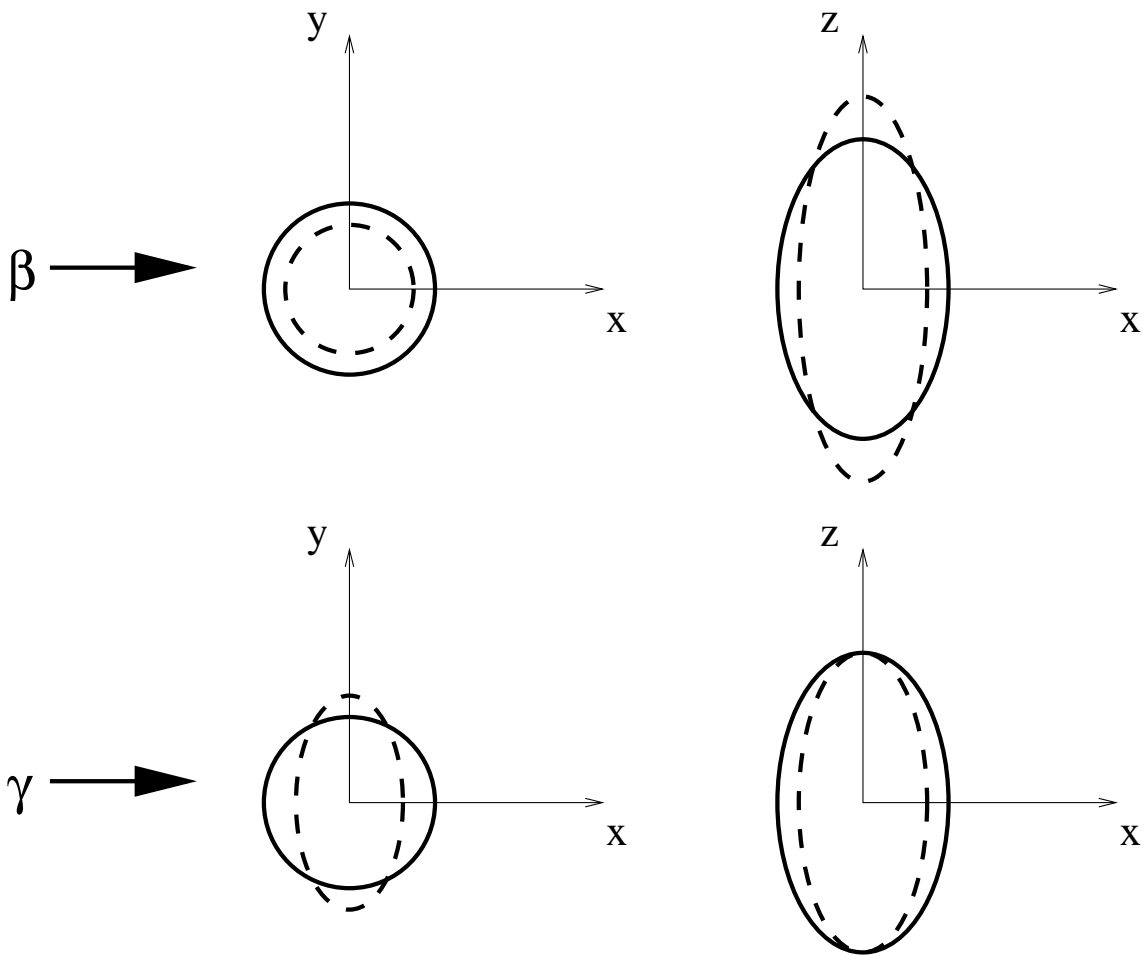
³Grand Accélérateur National d'Ions Lourds, B.P. 5027, F-14076 **Caen**
Cedex 5, **France**

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Vibro-rotational spectrum



Geometric representation of β and γ vibrations



Some experimental data

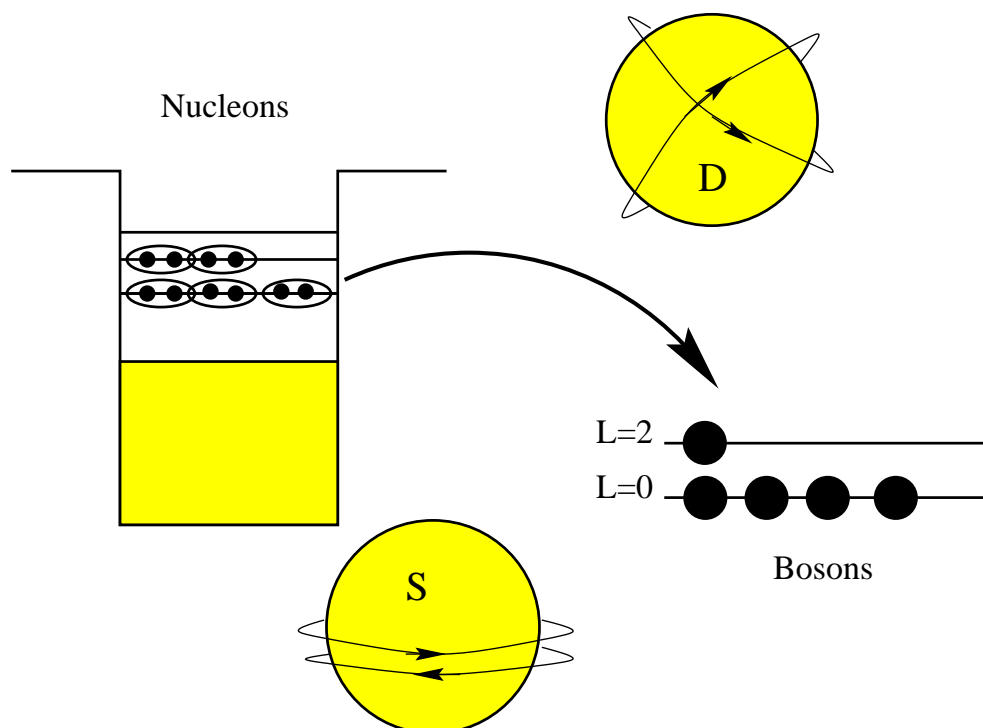
- ^{106}Mo : 1435 keV ($K^\pi = 4^+$)
- $^{154-156}\text{Gd}$: ($K^\pi = 4^+$)
- ^{164}Dy : 2173 ($K^\pi = 4^+$)
- ^{166}Er : 1943 keV ($K^\pi = 0^+$), 2028 keV ($K^\pi = 4^+$)
- ^{168}Er : 2055 keV ($K^\pi = 4^+$)
- $^{186-192}\text{Os}$, ^{194}Pt

Degree of anharmonicity

- ^{164}Dy : $\frac{E^*(4_{\gamma 2}^+)}{E^*(2_{\gamma}^+)} = \mathbf{2.95}$
- ^{166}Er : $\frac{E^*(0_{\gamma 2}^+)}{E^*(2_{\gamma}^+)} = \mathbf{2.47}$
- ^{166}Er : $\frac{E^*(4_{\gamma 2}^+)}{E^*(2_{\gamma}^+)} = \mathbf{2.58}$
- ^{168}Er : $\frac{E^*(4_{\gamma 2}^+)}{E^*(2_{\gamma}^+)} = \mathbf{2.50}$

The *I*nteracting *B*oson *M*odel

- The IBM is a model which describes the low lying collective states of medium mass and heavy nuclei.
- It can be considered as an approximation to the Shell Model. Two steps are necessary: truncation of the Shell Model space and bosonization of the nucleon pairs.



- The IBM can also be considered as the second quantization of the shape variables of the Geometric Collective Model.

Algebraic structure of the IBM

$$\begin{array}{ccc}
 s^\dagger, d_m^\dagger (m = 0, \pm 1, \pm 2) & \longrightarrow & \gamma_{lm}^\dagger, \gamma_{lm} \\
 s, d_m (m = 0, \pm 1, \pm 2) & & (l = 0, 2; \quad -l \leq m \leq l)
 \end{array}$$

$$[\gamma_{lm}, \gamma_{l'm'}^\dagger] = \delta_{ll'} \delta_{mm'}, [\gamma_{lm}^\dagger, \gamma_{l'm'}^\dagger] = 0, [\gamma_{lm}, \gamma_{l'm'}] = 0$$

- The dynamical algebra of the IBM is $U(6)$.

Generators $U(6)$: $\hat{G}_{ij} = \gamma_i^\dagger \gamma_j$, with $i, j = 1, \dots, 6$.

$$[\hat{G}_{ij}, \hat{G}_{kl}] = \hat{G}_{il} \delta_{jk} - \hat{G}_{jk} \delta_{il}$$

- Every dynamic operator can be written in terms of $U(6)$ generators.

$$\hat{H} = \sum_{ij} \varepsilon_{ij} \gamma_i^\dagger \gamma_j + \sum_{ijkl} V_{ijkl} \gamma_i^\dagger \gamma_j^\dagger \gamma_k \gamma_l$$

$$\hat{T} = \sum_{ij} t_{ij} \gamma_i^\dagger \gamma_j$$

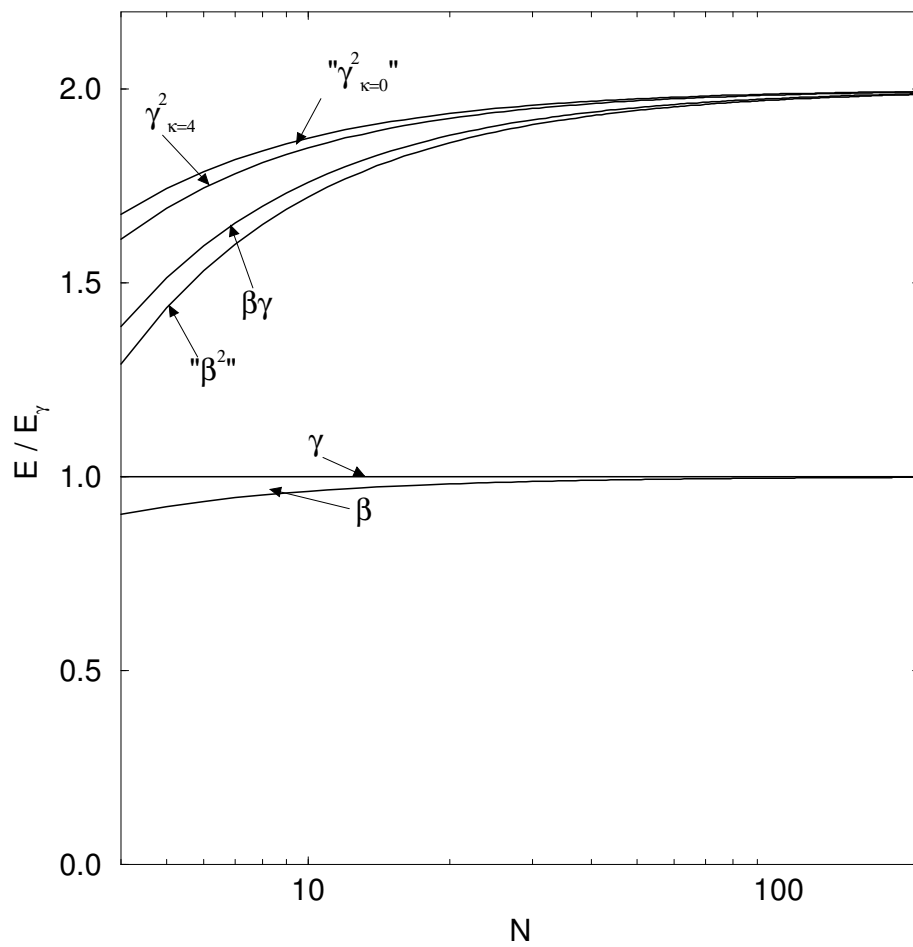
Generic Hamiltonian

$$\begin{aligned}\hat{H} &= \varepsilon_s \hat{n}_s + \varepsilon_d \hat{n}_d + \kappa_0 \hat{P}^\dagger \hat{P} + \kappa_1 \hat{L} \cdot \hat{L} \\ &+ \kappa_2 \hat{Q} \cdot \hat{Q} + \kappa_3 \hat{T}_3 \cdot \hat{T}_3 + \kappa_4 \hat{T}_4 \cdot \hat{T}_4,\end{aligned}$$

where \hat{n}_s and \hat{n}_d are the s and d boson number operators, respectively, and

$$\begin{aligned}\hat{P}^\dagger &= \frac{1}{2} d^\dagger \cdot d^\dagger - \frac{1}{2} s^\dagger \cdot s^\dagger, \\ \hat{L} &= \sqrt{10} (d^\dagger \times \tilde{d})^{(1)}, \\ \hat{Q} &= s^\dagger \tilde{d} + d^\dagger \tilde{s} + \chi (d^\dagger \times \tilde{d})^{(2)}, \\ \hat{T}_3 &= (d^\dagger \times \tilde{d})^{(3)} \\ \hat{T}_4 &= (d^\dagger \times \tilde{d})^{(4)}.\end{aligned}$$

$SU(3)$ limit



$$\hat{H} = -\hat{Q} \cdot \hat{Q}, \text{ with } \chi = \pm\sqrt{7}/2$$

Some definitions

Energy ratios.

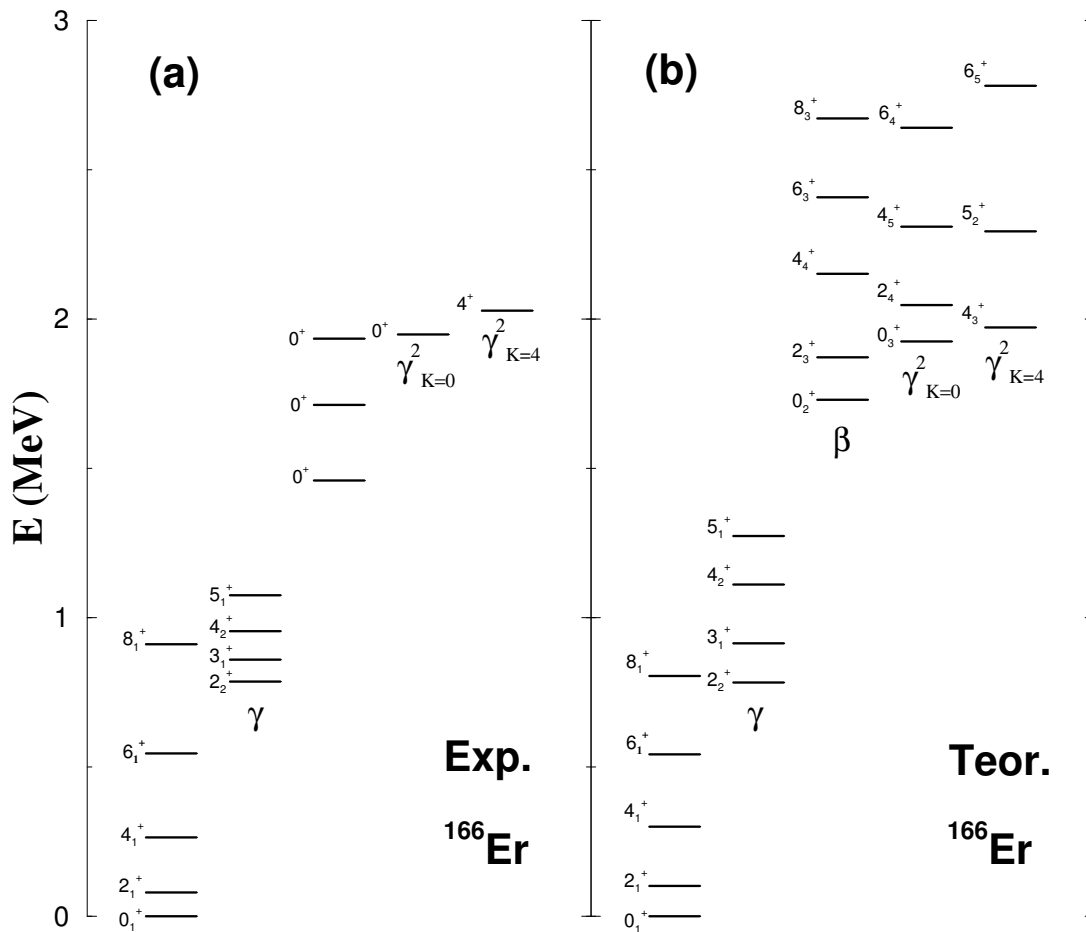
$$R_0^\gamma \equiv \frac{E_x(0_{\gamma\gamma}^+)}{E_x(2_\gamma^+) - E_x(2_1^+)}, \quad R_4^\gamma \equiv \frac{E_x(4_{\gamma\gamma}^+) - E_x(4_1^+)}{E_x(2_\gamma^+) - E_x(2_1^+)}$$

First extension: 3-body terms

Hamiltonian.

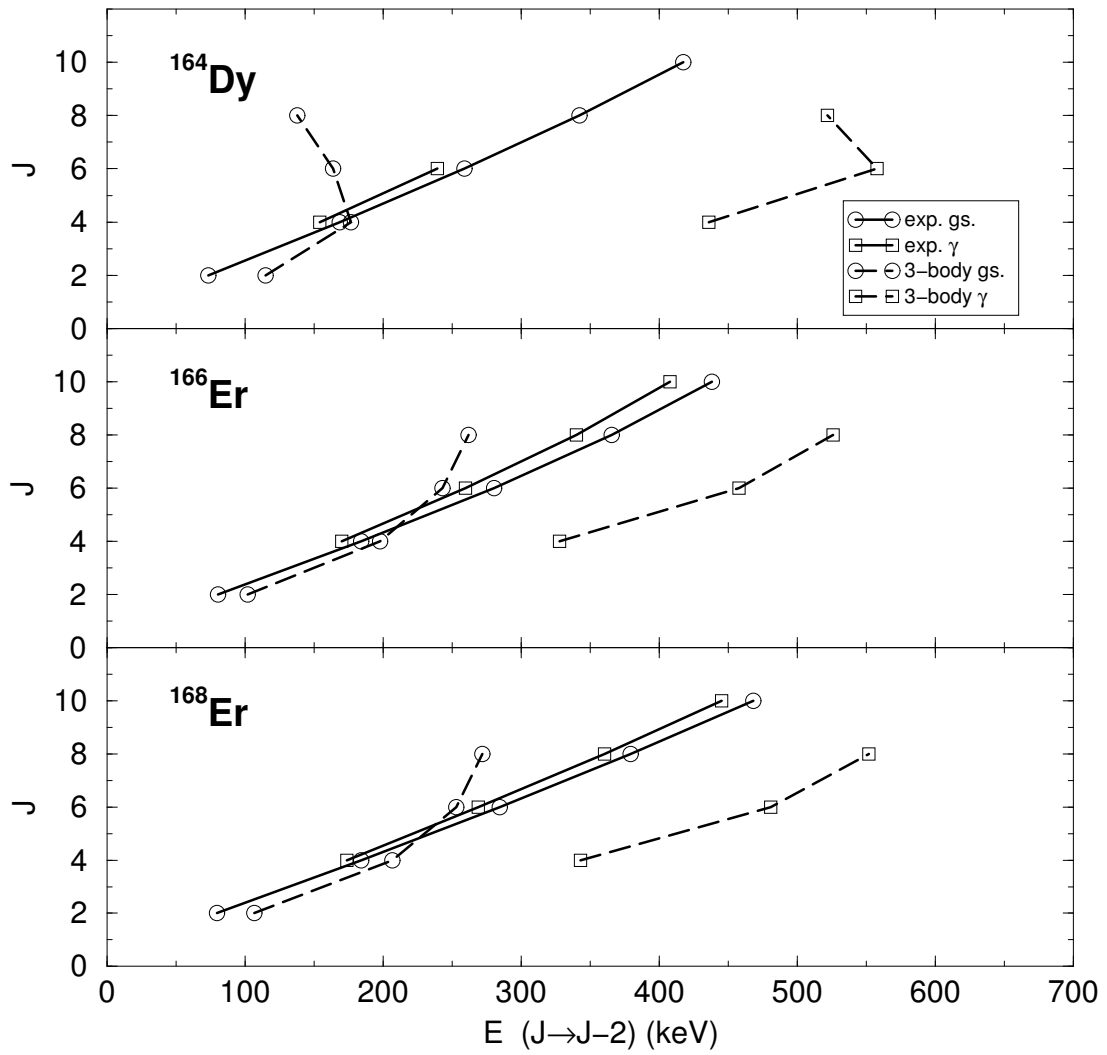
$$\begin{aligned}\hat{H} &= -\kappa\hat{Q} \cdot \hat{Q} + \kappa'\hat{L} \cdot \hat{L} \\ &+ \sum_{l=0,2,3,4,6} \theta_l \left((d^\dagger \times d^\dagger)^{(k)} \times d^\dagger \right)^{(l)} \cdot \left((\tilde{d} \times \tilde{d})^{(k)} \times \tilde{d} \right)^{(l)}\end{aligned}$$

Realistic calculation for ^{166}Er . Three-body Hamiltonian



Experimental (a) and theoretical (b) spectra for ^{166}Er . Hamiltonian parameters: $\kappa = 23.8$ keV, $\chi = -0.55$, $\kappa' = -1.9$ keV, and $\theta_4 = 31.3$ keV. $N = 15$.

Moment of inertia with a 3-body Hamiltonian



Moment of inertia: $\mathcal{I} \approx 2\hbar \frac{dJ}{dE_\gamma}$.

Beyond 3-body terms: 4-body terms?

- **An analytic approximation. The $SU(3)$ limit.**

$$\hat{H} = a \hat{C}_2[SU(3)] + b \hat{C}_3[SU(3)] + c \hat{C}_2[SU(3)]^2$$

Eigenvalues.

$$\begin{aligned} \langle (\lambda, \mu) | \hat{H} | (\lambda, \mu) \rangle &= a(\lambda^2 + \mu^2 + \lambda\mu + 3\lambda + 3\mu) \\ &+ b(\lambda - \mu)(2\lambda + \mu + 3)(\lambda + 2\mu + 3) \\ &+ c(\lambda^2 + \mu^2 + \lambda\mu + 3\lambda + 3\mu)^2 \end{aligned}$$

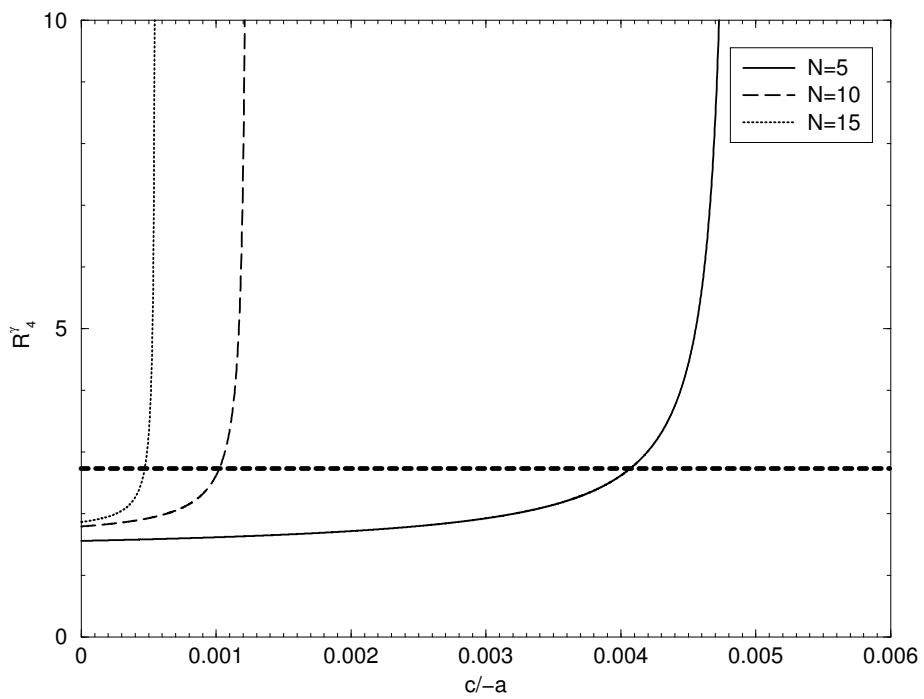
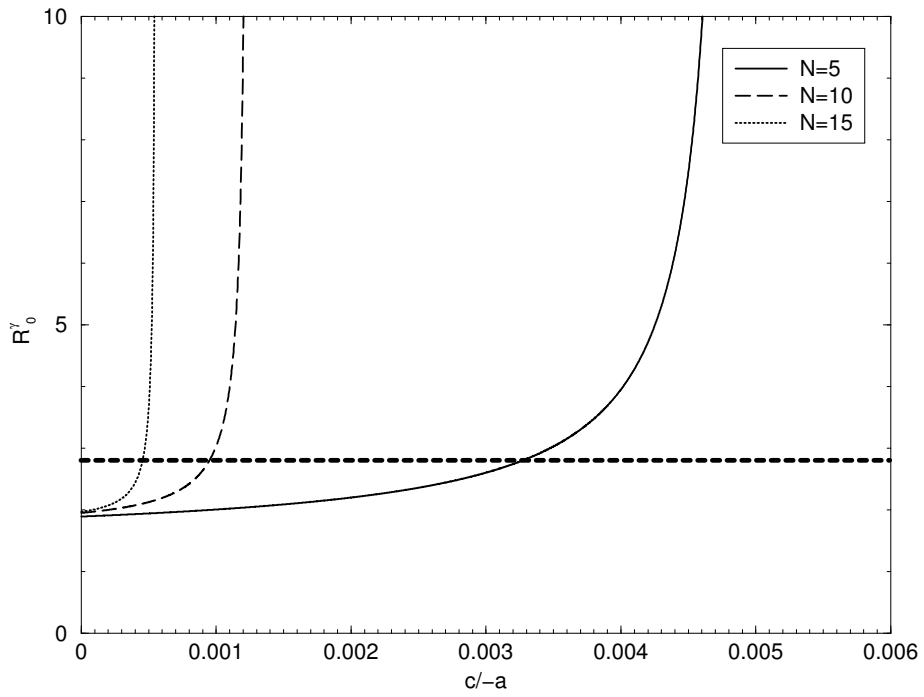
- **Q expansion.**

$$\begin{aligned} \hat{H}_Q &= \kappa' \hat{L} \cdot \hat{L} + a \hat{Q} \cdot \hat{Q} + b (\hat{Q} \times \hat{Q} \times \hat{Q})^{(0)} \\ &+ c (\hat{Q} \cdot \hat{Q})(\hat{Q} \cdot \hat{Q}) \end{aligned}$$

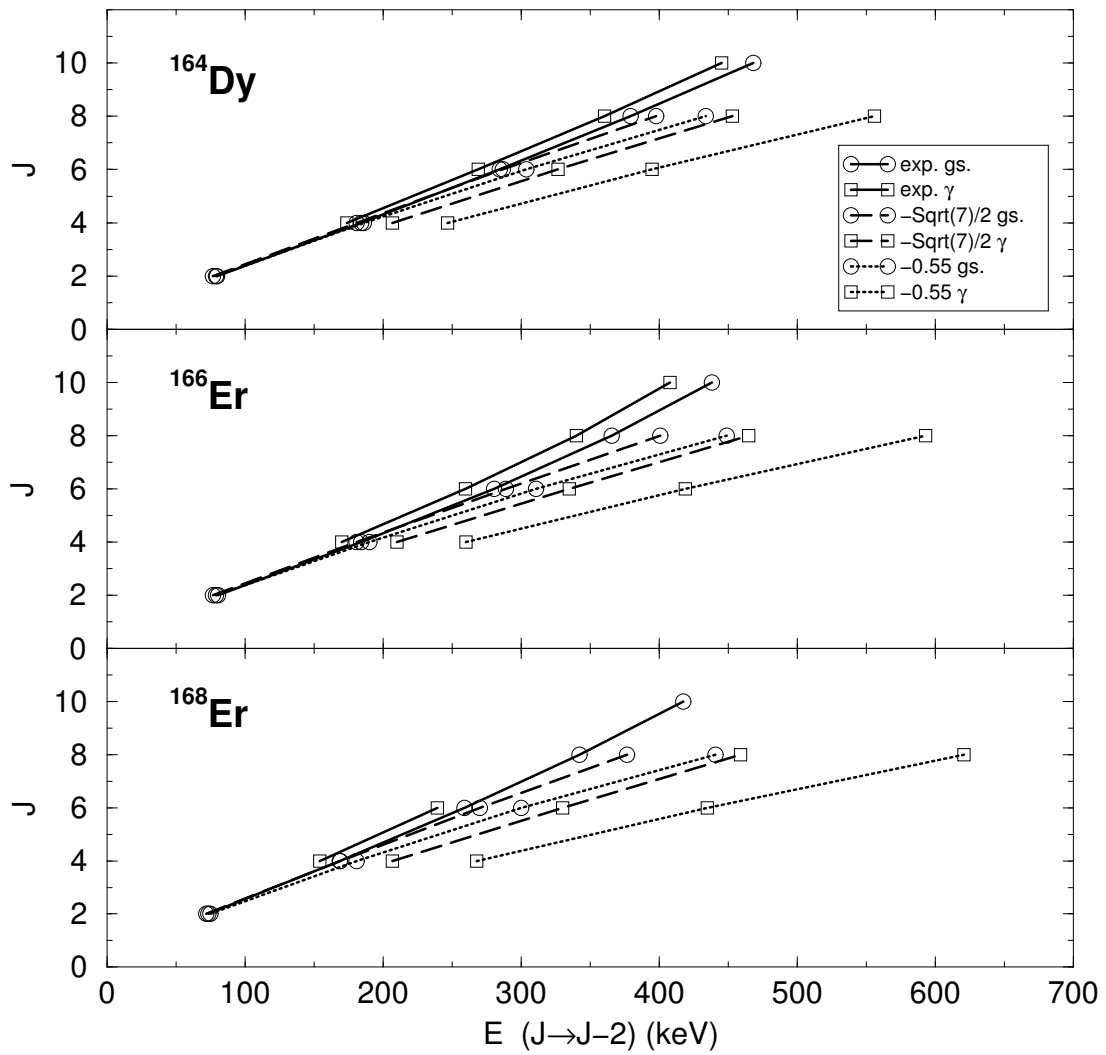
- **Pseudo-Casimir expansion.**

$$\begin{aligned} \hat{H}_{pC} &= \kappa' \hat{L} \cdot \hat{L} + a \hat{C}_2[SU(3)]_\chi + b \hat{C}_3[SU(3)]_\chi \\ &+ c \hat{C}_2[SU(3)]_\chi^2 \end{aligned}$$

Energy ratios in the $SU(3)$ limit

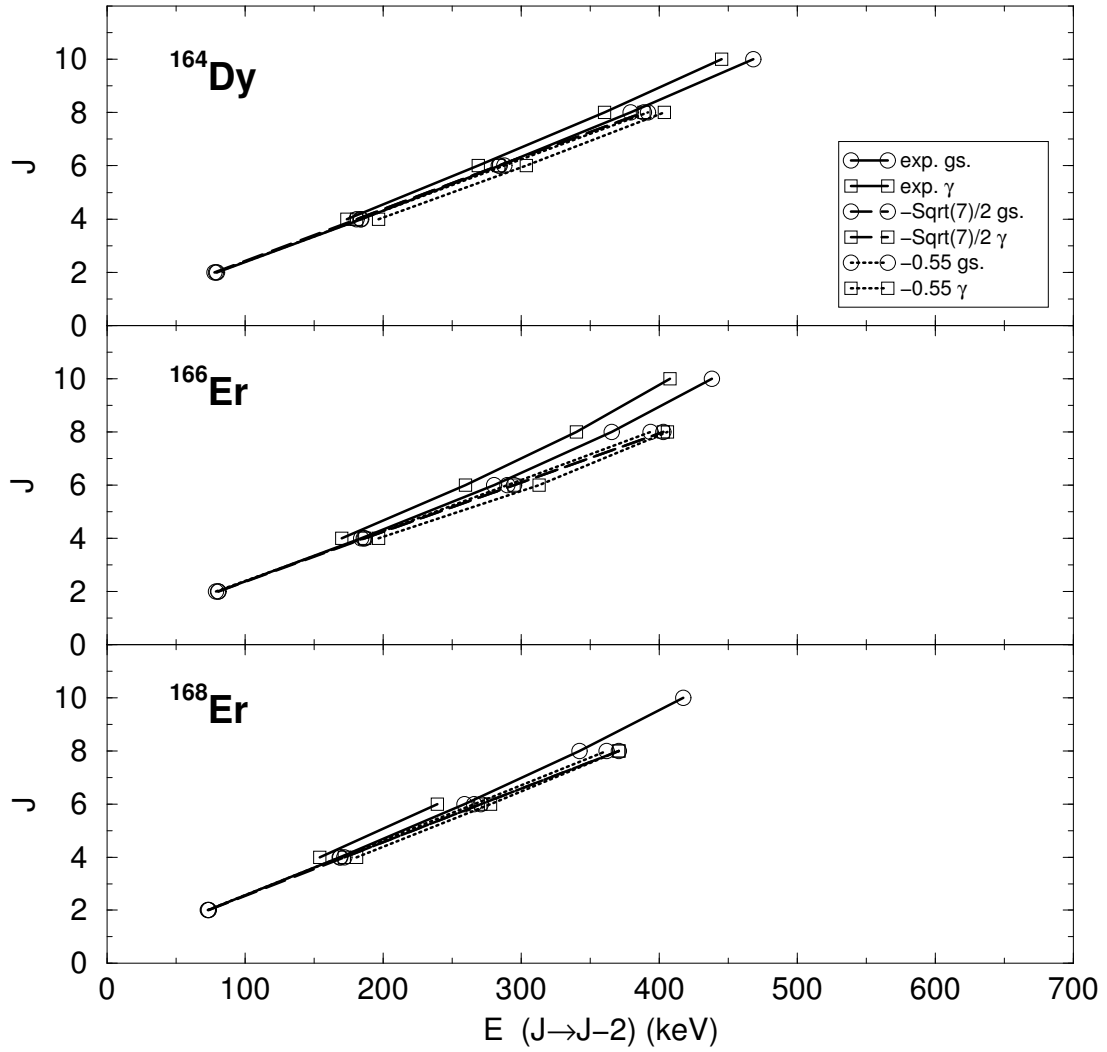


Moment of inertia with a \hat{Q} expansion



Moment of inertia: $\mathcal{I} \approx 2\hbar \frac{dJ}{dE_\gamma}$.

Moment of inertia with a pseudo-Casimir expansion

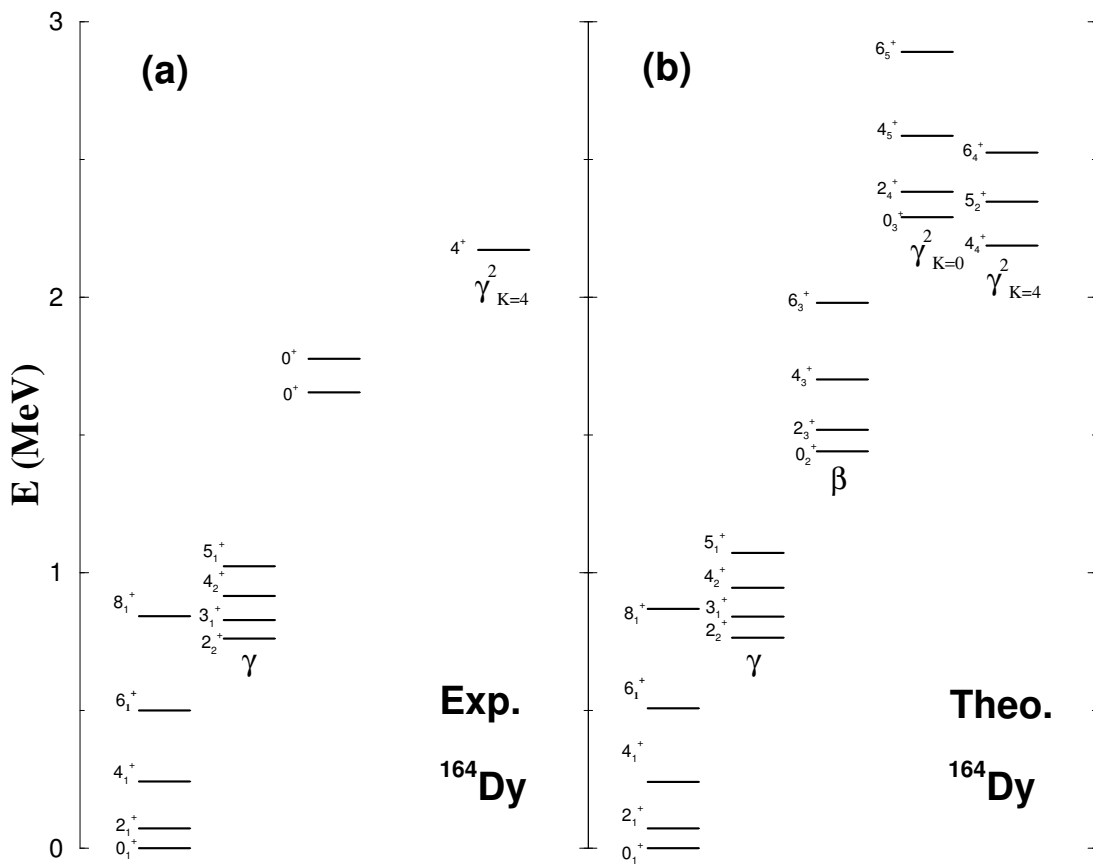


Moment of inertia: $\mathcal{I} \approx 2\hbar \frac{dJ}{dE_\gamma}$.

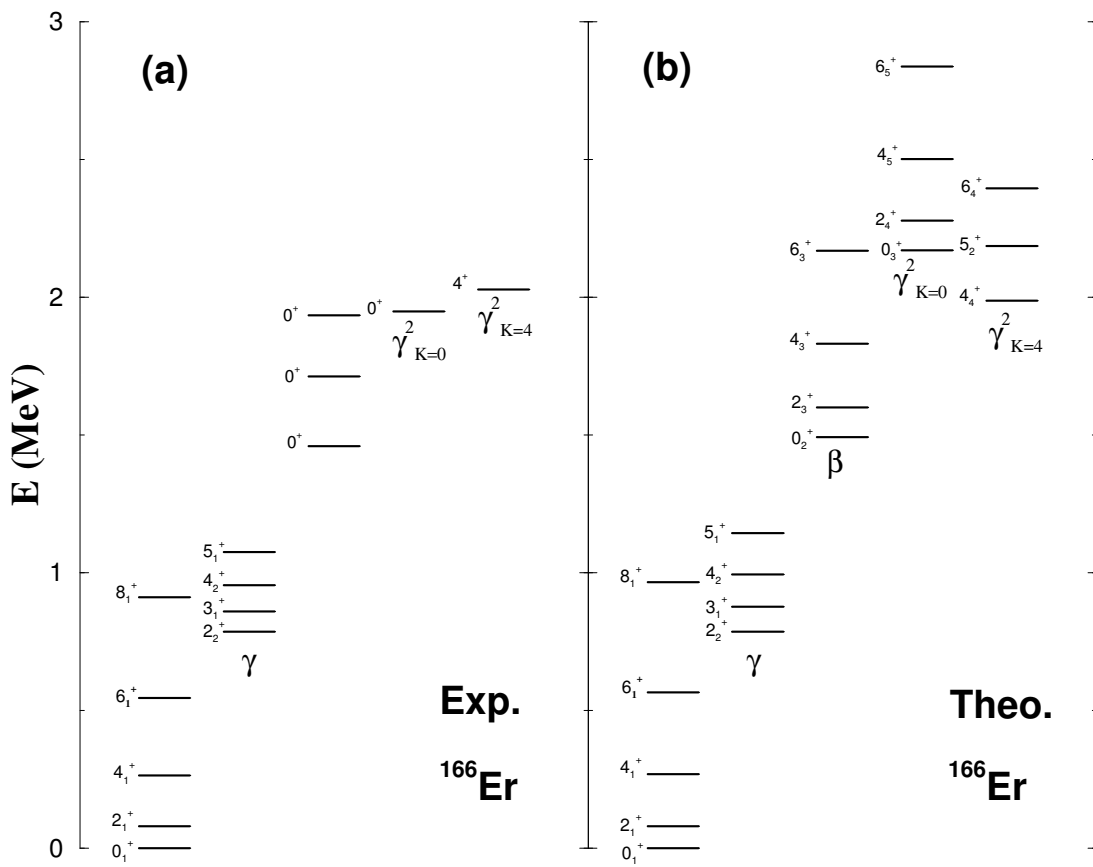
Final calculations for ^{164}Dy , ^{166}Er , and ^{168}Er

Nucleus	κ' (keV)	a (keV)	c (keV)	χ	N
^{164}Dy	12.18	-82.90	0.05150	-0.55	16
^{166}Er	13.55	-75.40	0.05286	-0.45	15
^{168}Er	13.23	-67.25	0.04080	-0.50	16

Spectrum for ^{164}Dy



Spectrum for ^{166}Er



Spectrum for ^{168}Er

